

Measurement of intermittency for charged particles in Au + Au collisions at $\sqrt{s_{NN}} = 7.7\text{-}200$ GeV from STAR

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Local density fluctuations near the QCD critical point can be probed by intermittency analysis of scaled factorial moments in relativistic heavy-ion collisions. We report the first measurement of intermittency for charged particles in Au + Au collisions at $\sqrt{s_{\rm NN}}=7.7\text{-}200$ GeV from the STAR experiment at RHIC. We observe scaling behaviors in central Au + Au collisions, with the extracted scaling exponent decreasing from mid-central to the most central Au + Au collisions. Furthermore, the scaling exponent exhibits a non-monotonic energy dependence with a minimum around $\sqrt{s_{\rm NN}}=20\text{-}30$ GeV in central Au + Au collisions.

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1 Introduction

The major goal of the Beam Energy Scan (BES) at the Relativistic Heavy Ion Collider (RHIC) is to explore the phase diagram of quantum chromodynamics (QCD) [1, 2]. An important landmark of the QCD phase structure is the critical point (CP), which is the end point of first-order phase boundary between quark-gluon and hadronic phases [3]. In the thermodynamic limit, the correlation length diverges at the CP and the system becomes scale invariant and fractal [4]. It is shown that the density fluctuations near the QCD critical point form a distinct pattern of power-law or intermittent behavior in the matter produced in high energy heavy-ion collisions [5].

In analogy to the critical opalescence observed in conventional matter near the critical point, the related fractal and self-similar geometry of QCD matter will lead to local density fluctuations that obey intermittent behavior [5]. Based on the effective action belonging to three-dimensional Ising universality class, the intermittency of QCD matter is revealed in transverse momentum spectra as a power-law (scaling) behavior of scaled factorial moment (SFM) in heavy-ion collisions [5]. An intermittent behavior has observed in Si + Si collisions at 158A



GeV from the NA49 experiment [6]. Meanwhile, studies based on a critical Monte Carlo with self-similar property [7] and transport model with hadronic potentials [8] demonstrate that the intermittency could be visible in Au + Au collisions at RHIC energies.

2 Analysis Details

In high-energy experiments, local power-law fluctuations can be detectable through the measurements of scaled factorial moment (SFM) which is defined as:

$$F_{q}(M) = \frac{\langle \frac{1}{M^{D}} \sum_{i=1}^{M^{D}} n_{i}(n_{i}-1) \cdots (n_{i}-q+1) \rangle}{\langle \frac{1}{M^{D}} \sum_{i=1}^{M^{D}} n_{i} \rangle^{q}},$$
(1)

where M^D is the number of cells in D-dimensional momentum space, n_i is the measured multiplicity in the i-th cell, and q is the order of moment.

Another expected power-law behavior that describes relationship between $F_q(M)$ and $F_2(M)$ is defined as [9,10]:

$$F_q(M) \propto F_2(M)^{\beta q}.$$
 (2)

Moreover, the scaling exponent ν quantitatively describes the values of β_q :

$$\beta_q \propto (q-1)^{\gamma}. \tag{3}$$

Here ν specifies scaling (power-law) behavior of $F_q(M)$. According to Ginzburg-Landau (GL) theory, the critical ν is equal to 1.304 in entire space phase [9], while it is equal to 1.0 from the two-dimensional Ising model [10].

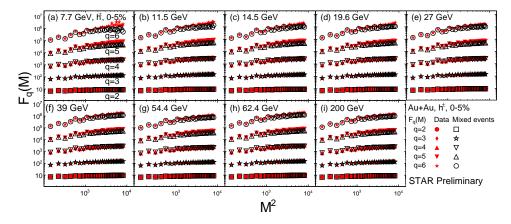


Figure 1: $F_q(M)$ (up to sixth order) of charged particles in transverse momentum space for the most central (0-5%) Au + Au collisions at $\sqrt{s_{\rm NN}}=7.7$ -200 GeV in double-logarithmic scale.

The data reported here were obtained from Au + Au collisions at $\sqrt{s_{\rm NN}}=7.7,\,11.5,\,14.5,\,19.6,\,27,\,39,\,54.4,\,62.4$ and 200 GeV, which were recorded by the STAR experiment at RHIC from 2010 to 2017. Protons (p), antiprotons (\bar{p}) , kaons (K^\pm) and pions (π^\pm) are analyzed as charged particles, and their identifications are carried out using the Time Projection Chamber (TPC) and the Time-of-Flight (TOF) detectors. To avoid the self-correlation, the centrality was determined from uncorrected charged particles within a pseudo-rapidity window of $0.5 < |\eta| < 1$, which was chosen to be beyond the analysis window of $|\eta| < 0.5$.



To subtract the background at the level of SFM, a correlator $\Delta F_q(M)$ is defined in terms of original and mixed events, i.e., $\Delta F_q(M) = F_q(M)^{data} - F_q(M)^{mix}$ [6]. In addition, a cell-by-cell method is proposed for efficiency correction on SFM [11]. The statistical uncertainties are estimated by Bootstrap method, and the systematic uncertainties are estimated by varying the experimental requirements for tracks in the TPC and TOF.

3 Results and Discussion

Figure 1 shows $F_q(M)^{data}$ and $F_q(M)^{mix}$, from the second order to the sixth order in the most central (0-5%) collisions for various $\sqrt{s_{\rm NN}}$. Based on the statistics of BES-I data, $F_q(M)$ can be calculated in the range of M^2 from 1 to 100^2 and up to the sixth order (q=6). It is observed that $F_q(M)^{data}$ is larger than $F_q(M)^{mix}$ at large M^2 region for various $\sqrt{s_{\rm NN}}$, thus a deviation of $\Delta F_q(M)$ from zero is present in central Au + Au collisions.

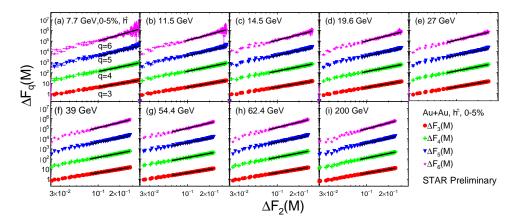


Figure 2: $\Delta F_q(M)$ (q=3-6) as a function of $\Delta F_2(M)$ in the most central (0-5%) Au + Au collisions at $\sqrt{s_{\rm NN}} = 7.7\text{-}200$ GeV in double-logarithmic scale.

Figure 2 shows $\Delta F_q(M)$ (q=3-6), as a function of $\Delta F_2(M)$ in the most central (0-5%) collisions for various $\sqrt{s_{\rm NN}}$. We clearly observe that the correlators $\Delta F_q(M)$ (q=3-6) exhibit scaling behavior with $\Delta F_2(M)$.

The value of β_q is obtained through a power-law fit of Eq. (2) as shown in Figure 2, and its statistical error is determined by the fit. Figure 3(a) shows β_q as a function of q-1 in the most central Au + Au collisions for $\sqrt{s_{\rm NN}}=7.7\text{-}200$ GeV. Consistent with theoretical expectation, β_q also obeys a good scaling behavior with q, thus ν can be obtained through a power-law fit of Eq. (3). Figure 3(b) shows the extracted ν as a function of $\langle N_{part} \rangle$ in central Au + Au collisions at various $\sqrt{s_{\rm NN}}$. We find that ν decreases from mid-central (30-40%) to the most central (0-5%) Au + Au collisions.

Figure 4 shows the energy dependence of ν of charged particles in central Au + Au collisions at $\sqrt{s_{\rm NN}} = 7.7\text{-}200$ GeV. It is observed that the ν exhibits a non-monotonic behavior on collision energy and seems to reach a minimum around $\sqrt{s_{\rm NN}} = 20\text{-}30$ GeV. Higher statistics data from BES-II will help to confirm the trend of energy dependence of ν .

4 Summary

In summary, we report the first measurements of intermittency for charged particles in Au + Au collisions at $\sqrt{s_{\rm NN}} = 7.7\text{-}200$ GeV from the STAR experiment. Scaled factorial moments



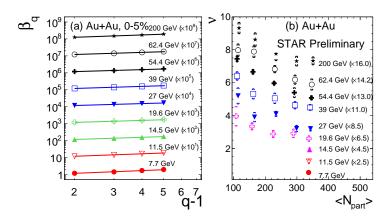


Figure 3: (a) β_q (q=3-6) as a function of q-1 in most central Au + Au collisions at $\sqrt{s_{\rm NN}} = 7.7\text{-}200$ GeV. (b) ν as a function of $\langle N_{part} \rangle$ in central Au + Au collisions.

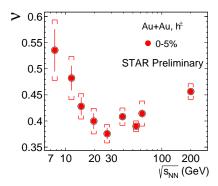


Figure 4: Energy dependence of ν for charged particles in Au + Au collisions at $\sqrt{s_{\rm NN}}$ = 7.7- 200 GeV. The statistical and systematic errors are shown in bars and brackets, respectively.

(up to the sixth order) for p, \bar{p} , K^{\pm} and π^{\pm} within $|\eta| < 0.5$, have been measured in available transverse momentum space. Scaling behavior is clearly visible in Au + Au collisions which is consistent with theoretical predictions. The scaling exponent is related to the critical component, and we observe that it shows a non-monotonic behavior on $\sqrt{s_{\rm NN}}$ with a dip around 20-30 GeV in the most central (0-5%) Au + Au collisions. This non-monotonic behavior needs to be understood with more theoretical inputs. With significantly improved statistics, the RHIC BES Phase-II program will allow for a more precise measurement of intermittency in heavy-ion collisions.

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