Pressure-dependent structural and electronic instabilities in LaSb₂

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Abstract

LaSb₂ exhibits a large, non-saturating, linear magnetoresistance at low temperatures, defying the expectations of Fermi liquid theory. This is thought to be caused by charge density wave order emerging below an abrupt, hysteretic anomaly in the resistivity at \sim 355 K. We find that, under hydrostatic pressure, this anomaly becomes much more pronounced, develops strong hysteresis, and shifts rapidly towards lower temperatures. The anomaly is fully suppressed by only 6 kbar. Moreover, we observe a second transition anomaly at lower temperature, which likewise disappears under pressure. These findings are discussed in the context of a structural transition recently discovered in the sister material CeSb₂ and of density functional theory calculations, which indicate that the SmSb₂-type structure adopted by LaSb₂ at ambient conditions is unstable at moderate applied pressures.

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Introduction 1

The recent discovery of unconventional superconductivity in CeSb₂ beyond a pressure-induced structural transition [1], has prompted a re-examination of the rare-earth diantimonide family. The RSb₂ materials crystallise in a series of highly-anisotropic, layered structures which host a variety of electronic and magnetic ground states [2-4]. The nature of this magnetic ordering is governed by the interplanar R-R distance, making them ideal for tuning magnetism under pressure [5].

LaSb₂, the isostructural, non-magnetic counterpart to CeSb₂, exhibits superconductivity at ambient pressure [6]. While the magnetoresistance (MR) of a conventional Fermi-liquid metal is expected to scale quadratically with increasing magnetic field, the large MR of LaSb₂ remains linear, without saturating in fields in excess of 45 T [6,7]. The observation of de



Figure 1: The temperature dependence of the resistivity of LaSb₂ is interrupted by two hysteretic anomalies; firstly, by an abrupt jump at ~ 355 K, labelled α ; and secondly, by a wider hysteresis loop at around 250 K, labelled β . (inset) LaSb₂ adopts the SmSb₂ structure (space-group Cmca) under ambient conditions.

Haas van Alphen oscillations rules out explanations stemming from impurity effects [8]. It has been proposed that, instead, this anomalous MR is associated with a charge density wave (CDW) state, driven by strong Fermi surface nesting in LaSb₂, which develops at ~ 355 K with a pronounced anomaly in the resistance, labelled α in Figure 1 [9]. The presence of CDW order in the closely related material LaAgSb₂ has been confirmed by X-ray measurements [10], but direct evidence of CDW order in LaSb₂ remains elusive so far [11, 12].

Here, we explore the nature of the 355 K resistance anomaly in a series of high pressure measurements and by *ab initio* density functional theory (DFT) calculations.

2 Methods

Single crystals of LaSb₂ were grown by self flux from La (4N) and Sb (6N) in the ratio 1:9 and excess Sb was removed with a centrifuge [13]. The resulting crystals were formed in stacks of large plates (~ 10 × 10 mm), with the *c*-axis perpendicular to the plane. AC resistivity measurements were performed in the four-point configuration with I||*ab*. Good electrical contact was ensured by spot-welding 25 μ m Au wires to each sample and applying Dupont 6838 Ag paint. Pressure was generated using a BeCu/MP35N piston-cylinder cell (PCC) [14, 15] with glycerol as the pressure-transmitting medium. The pressure inside the cell was determined from the superconducting transition temperature of a sample of Sn, while the temperature dependence of its measured resistivity was used to correct for thermal lag [16]. Measurements were recorded between 400 K and 1.76 K using the ACT module of the Quantum Design Physical Property Measurement System (PPMS).

Electronic structure calculations were performed using CASTEP [17]. Pseudopotentials were generated within the generalised gradient approximation [18] and calculations were converged on $15 \times 15 \times 5$ Monkhorst-Pack *k*-space grids [19]. These parameters reproduce the ambient pressure lattice parameters reasonably well (Table 1). The crystal structure of LaSb₂ was modelled under pressure by optimising the free energy as a function of both the lattice parameters and internal atomic positions. Within CASTEP this is done by using the Broyden–Fletcher–Goldfarb–Shanno algorithm [20], where the Hessian of atomic forces is calculated in the mixed space of internal and cell degrees of freedom. This allows simultaneous

Table 1: Ambient pressure lattice parameters of LaSb ₂ obtained by X-ray diffraction
(XRD) measurements are consistent with DFT results calculated using Perdew-Burke-
Ernzerhof (PBE) pseudopotentials.

Method	a/Å	b/Å	c/Å	Vol. per La/Å ³
XRD (300 K) [3]	6.32	6.17	18.57	90.55
DFT	6.39	6.27	18.62	93.31

optimisation of both lattice parameters and atomic coordinates. Further, CASTEP performs optimisation at fixed external stress allowing cell optimisation to be performed with varying external pressure.

3 Results

At ambient pressure, the resistivity of LaSb₂ exhibits two clear hysteretic anomalies: anomaly α at ~ 350 K, and anomaly β at ~ 250 K (Figure 1). Anomaly α has only a small hysteresis on the order of 5 K, whereas anomaly β exhibits significant hysteresis over a range on the order of 100 K. The hysteresis indicates that both of these transition are first order. These two anomalies can be distinguished by the order of the hysteresis in the resistivity.



Figure 2: (i) The pressure dependence of the resistivity of LaSb₂ displays only a single, hysteretic anomaly above ambient pressure. This anomaly is rapidly suppressed, and disappears by ~ 6 kbar. As this occurs, the magnitude of the change in resistivity decreases, whereas the width of the hysteresis loop increases. (inset left) Anomaly α exhibits only moderate hysteresis on the order of ~ 5 K. (inset right) A small hysteresis loop is still visible at 4.69 kbar, after correcting for thermal lag. (ii) The resistivity of CeSb₂ under pressure also displays a hysteretic anomaly, which bears significant resemblance to the one seen in LaSb₂. This anomaly has been confirmed to correspond to a structural transition from the SmSb₂ ambient pressure structure to a different high pressure structure, most likely of the YbSb₂-type [1].



Figure 3: The phase diagram of LaSb₂ under pressure depicts the temperature dependence of the superconducting transition and the suppression of anomalies α and β . The superconductivity is initially enhanced, peaking at ~4 kbar, before weakening with further increasing pressure. Filled markers correspond to data measured in a PCC, while empty markers represent data recorded at ambient pressure. A circular marker indicates T_{SC} as measured by Galvis *et al.* [12]. The temperatures of anomalies α and β on warming and cooling are marked with triangles pointing in different directions.

Even at the lowest applied pressure of $\simeq 1$ kbar, a single, pronounced resistivity anomaly replaces the two original anomalies, α and β , observed at ambient pressure (Figure 2 i). As the directions of the resistivity jumps when warming and cooling are the same as in the ambient pressure α anomaly, we identify this anomaly as a continuation of α . However, the high pressure transition signature is highly hysteretic – like the β transition at ambient pressure. The transition temperature is suppressed very rapidly from its ambient pressure value and disappears altogether at approximately 6 kbar. Alongside this suppression, increasing pressure reduces the magnitude of the resistance anomaly and increases the range of its temperature hysteresis. This increase in the hysteresis width as the transition is suppressed to lower temperatures may be attributed to kinetic effects associated with a strongly first order phase transition. However, additional contributions due to non-hydrostatic pressure conditions caused by the frozen pressure medium cannot be quantified at present. Similar results were recently presented in [21].

The superconducting transition temperature, T_{SC} , also depends on pressure (Figure 3). At ambient pressure, the transition is broad with a full superconducting state only reached below 1.2 K [12]. Increasing the pressure to ~1 kbar brings the transition to above 1.7 K and significantly sharpens it. The pressure phase diagram indicates a maximum in T_{SC} of ~2.5 K at around 3-4 kbar. Further increasing the pressure then slowly suppresses T_{SC} .

4 Discussion

The shape and pressure dependence of the observed resistance anomaly are strongly reminiscent of the case of $CeSb_2$ (Figure 2 ii). In $CeSb_2$, high pressure x-ray diffraction (XRD) measurements confirm a pressure-induced structural transition from the $SmSb_2$ -type structure to $YbSb_2$ -type structure [1]. This prompts the question whether the observed resistivity anomalies in $LaSb_2$ might, likewise, be attributed to a structural instability rather than CDW order.

There has not yet been any direct measurement of CDW order in LaSb₂. The identification of the α anomaly as a signature of a CDW transition [21] primarily rests on the large linear MR observed in LaSb₂ at low temperature. However, CDW order is neither necessary nor sufficient to explain this phenomenon, as similar linear MR also arises in some Weyl and Dirac semi-metals [22–24] characterised by small Fermi surface pockets. Moreover, linear MR is observed at higher pressures even after anomaly α is fully suppressed [21]. CDW order has been reported in the (La/Ce)Sb₂ substitution series [9], but the study lacks direct observation of the CDW ordering wavevector by XRD, whereas the reported changes in bond lengths and lattice parameters might also be consistent with a structural transition. Furthermore, no evidence of a CDW ordering wavevector was found when directly probing LaSb₂ at low temperatures via scanning tunnelling microscopy [12].

It is notable that crystals of $LaSb_2$ grown by flux are predisposed towards twinning [6], which is often an indication that the crystal has been cooled through a structural transition, like in the case in the famous Dauphiné twins of quartz [25]. This is consistent with reports of a distinct, high-temperature form of $LaSb_2$ [26]. It has also been suggested that $LaSb_2$ is susceptible to a pressure-induced structural transition, in which the structure becomes more symmetric in the *ab* plane, and the *c*-axis is compressed. This is indeed the case in CeSb₂ [6].

The crystal structures of the rare-earth diantimonides which can be grown at ambient pressure (R =La-Tb, Yb) consist of R-Sb bilayers, separated by a distorted square net of Sb [2, 5]. The majority of the rare-earth diantimonides, including LaSb₂, crystallise in the orthorhombic SmSb₂ structure (space group Cmca, Figure 4 left). EuSb₂ and YbSb₂, however, both adopt higher symmetry variations of this structure, in which the neighbouring R-Sb layers are offset and pulled together by zig-zag chains of Sb atoms. The former adopts the P2₁/m space-group,



Figure 4: *Ab initio* calculations (see text) imply that $LaSb_2$, which adopts the $SmSb_2$ structure under ambient conditions (left), becomes unstable to the more symmetric YbSb₂ structure (right) under pressure. Structures generated using VESTA software package [27].

due to a distortion driven by the Eu^{2+} valence, whereas in the latter, a structure with the Cmcm space-group is stabilised by the contribution of trivalent Yb on the predominantly Yb²⁺ site [5]. As La does not support a divalent state, we expect it to undergo a transition to the YbSb₂-type structure under pressure, where the compressed layers and smaller volume reduce the free energy relative to the SmSb₂ structure [1].

In the transition from the $SmSb_2$ -type structure to the $YbSb_2$ -type structure, the planes of La and Sb atoms undergo a relative shear which allows the individual planes to come closer to each other. Consequently, the *c*-axis shortens, similar to the uncollapsed to collapsed tetragonal transition in some iron pnictides such as CaFe₂As₂ [28]. This behaviour is also seen in XRD reconstructed lattice parameters for the SmSb₂-to-YbSb₂ transition in CeSb₂ [1].

We have conducted extensive DFT calculations to investigate the possibility of a pressureinduced structural transition in LaSb₂ that matches the one observed in CeSb₂. These calculations suggest that at a surprisingly low pressure $\simeq 5$ kbar, the SmSb₂-type structure, adopted by LaSb₂ under ambient conditions, could become unstable to the YbSb₂-type structure. At ambient pressure, the SmSb₂-type structure is favoured as the ground state relative to the YbSb₂-type structure by 2.5×10^{-3} eV/La. However, under pressure the larger volume of the SmSb₂-type structure relative to the YbSb₂-type structure means that it is rapidly destabilised. By 5 kbar the YbSb₂-type structure is energetically more favourable than the SmSb₂type structure and remains the relative ground state structure up to at least 10 kbar. This is in good agreement with our high pressure experiments and supports the existence of a pressureinduced structural transition in LaSb₂.



Figure 5: (i) The absolute energies of the SmSb_2 -type structure and YbSb_2 -type structure vary as function of pressure from 0 to 10 kbar. (ii) The absolute difference in energies between the SmSb_2 -type structure and the YbSb_2 -type structure indicates that there is a theoretical transition at 5 kbar. (iii) The volume of both cells per La atom demonstrates that the YbSb_2 has a smaller volume making it more favourable at high pressures. (iv) The *c*-axis of both structures shows how the transition from the SmSb_2 -type structure to the YbSb_2 -type structure corresponds to a collapse of the *c*-axis and shift in atomic planes. (v) The behaviour of the *a* and *b* lattice parameters in the SmSb_2 -type structure and effective lattice parameter a' in the YbSb_2 -type structure indicate that the dominant compression is in the c-direction. (vi) The La-La separation in both structures is a further indicator of how the SmSb_2 -to- YbSb_2 transition corresponds to a volume collapse.

5 Conclusion

The two high-temperature resistivity anomalies and the low temperature superconducting transition in ambient pressure LaSb₂ have been tracked under applied pressure. Anomaly α is rapidly suppressed to 0 K by ~ 6 kbar, whereas anomaly β is not visible above ambient pressure. At present, these anomalies cannot be definitively assigned to either CDW ordering or a structural transition. Further work such as high pressure XRD will be required to determine whether the high pressure resistivity anomaly observed in LaSb₂ may be attributed to a structural transition similar to that observed in CeSb₂.

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References

- [1] O. P. Squire et al., *Superconductivity beyond the Pauli limit in high-pressure CeSb*₂, (arXiv preprint) doi:10.48550/arXiv.2211.00975.
- [2] R. Wang and H. Steinfink, The crystal chemistry of selected AB2 rare earth compounds with selenium, tellurium, and antimony, Inorg. Chem. 6, 1685 (1967), doi:10.1021/ic50055a017.
- G. Borzone, A. Borsese, A. Saccone and R. Ferro, *Heats of formation of lanthanum-antimony alloys*, J. Less Common Met. **65**, 253 (1979), doi:10.1016/0022-5088(79)90115-2.
- [4] S. L. Bud'ko, P. C. Canfield, C. H. Mielke and A. H. Lacerda, Anisotropic magnetic properties of light rare-earth diantimonides, Phys. Rev. B 57, 13624 (1998), doi:10.1103/PhysRevB.57.13624.
- [5] W. Tremel and R. Hoffmann, *Square nets of main-group elements in solid-state materials*, J. Am. Chem. Soc. **109**, 124 (1987), doi:10.1021/ja00235a021.
- [6] K. F. F. Fischer, N. Roth and B. B. Iversen, *Transport properties and crystal structure of layered LaSb*₂, J. Appl. Phys. **125**, 045110 (2019), doi:10.1063/1.5081819.
- [7] D. P. Young et al., *High magnetic field sensor using LaSb*₂, Appl. Phys. Lett. 82, 3713 (2003), doi:10.1063/1.1577390.
- [8] R. G. Goodrich et al., *de Haas-van Alphen measurements of the electronic structure of LaSb*₂, Phys. Rev. B **69**, 125114 (2004), doi:10.1103/PhysRevB.69.125114.

- [9] R. F. Luccas et al., *Charge density wave in layered La*_{1-x}*Ce*_x*Sb*₂, Phys. Rev. B **92**, 235153 (2015), doi:10.1103/PhysRevB.92.235153.
- [10] C. Song et al., *Charge-density-wave orderings in LaAgSb*₂: An x-ray scattering study, Phys. Rev. B 68, 035113 (2003), doi:10.1103/PhysRevB.68.035113.
- [11] J. F. DiTusa et al., Optical conductivity and superconductivity in LaSb₂, J. Phys. Conf. Ser. 273, 012151 (2011), doi:10.1088/1742-6596/273/1/012151.
- [12] J. A. Galvis, H. Suderow, S. Vieira, S. L. Bud'ko and P. C. Canfield, Scanning tunneling microscopy in the superconductor LaSb₂, Phys. Rev. B 87, 214504 (2013), doi:10.1103/PhysRevB.87.214504.
- [13] P. C. Canfield and Z. Fisk, Growth of single crystals from metallic fluxes, Philos. Mag. B 65, 1117 (1992), doi:10.1080/13642819208215073.
- K. Semeniuk, Correlated low temperature states of YFe₂Ge₂ and pressure metallised NiS₂, PhD thesis, University of Cambridge, UK (2018), doi:10.17863/CAM.21468.
- [15] M. I. Eremets, *High pressure experimental methods*, Oxford Science Publications, Oxford, UK, ISBN 9780198562696 (1996).
- [16] A. Eiling and J. S. Schilling, Pressure and temperature dependence of electrical resistivity of Pb and Sn from 1-300K and 0-10 GPa-use as continuous resistive pressure monitor accurate over wide temperature range; superconductivity under pressure in Pb, Sn and In, J. Phys. F Met. Phys. 11, 623 (1981), doi:10.1088/0305-4608/11/3/010.
- [17] S. J. Clark et al., First principles methods using CASTEP, Z. Krist. Cryst. Mater. 220, 567 (2005), doi:10.1524/zkri.220.5.567.65075.
- [18] J. P. Perdew, K. Burke and Y. Wang, Generalized gradient approximation for the exchange-correlation hole of a many-electron system, Phys. Rev. B 54, 16533 (1996), doi:10.1103/PhysRevB.54.16533.
- [19] H. J. Monkhorst and J. D. Pack, Special points for Brillouin-zone integrations, Phys. Rev. B 13, 5188 (1976), doi:10.1103/PhysRevB.13.5188.
- [20] B. G. Pfrommer, M. Côté, S. G. Louie and M. L. Cohen, *Relaxation of crystals with the quasi-Newton method*, J. Comput. Phys. 131, 233 (1997), doi:10.1006/jcph.1996.5612.
- [21] S. L. Bud'ko, S. Huyan, P. Herrera-Siklody and P. C. Canfield, Rapid suppression of charge density wave transition in LaSb₂ under pressure, Philos. Mag. 103, 561 (2022), doi:10.1080/14786435.2022.2159561.
- [22] M. N. Ali et al., Large, non-saturating magnetoresistance in WTe₂, Nature 514, 205 (2014), doi:10.1038/nature13763.
- [23] J. Feng et al., Large linear magnetoresistance in Dirac semimetal CdAs₂ with Fermi surfaces close to the Dirac points, Phys. Rev. B 92, 081306 (2015), doi:10.1103/PhysRevB.92.081306.
- [24] C. Shekhar et al., Extremely large magnetoresistance and ultrahigh mobility in the topological Weyl semimetal candidate NbP, Nat. Phys. 11, 645 (2015), doi:10.1038/nphys3372.
- [25] J. Van Landuyt, G. Van Tendeloo, S. Amelinckx and M. B. Walker, Interpretation of Dauphiné-twin-domain configurations resulting from the α-β phase transition in quartz and aluminum phosphate, Phys. Rev. B 31, 2986 (1985), doi:10.1103/PhysRevB.31.2986.

- [26] J. J. Murray and J. B. Taylor, Halide vapor transport of binary rare-earth arsenides, antimonides and tellurides, J. Less Common Met. 21, 159 (1970), doi:10.1016/0022-5088(70)90114-1.
- [27] K. Momma and F. Izumi, VESTA3 for three-dimensional visualization of crystal, volumetric and morphology data, J. Appl. Cryst. 44, 1272 (2011), doi:10.1107/S0021889811038970.
- [28] A. Kreyssig et al., Pressure-induced volume-collapsed tetragonal phase of CaFe₂As₂ as seen via neutron scattering, Phys. Rev. B 78, 184517 (2008), doi:10.1103/PhysRevB.78.184517.