

W+c-jet production at the LHC with NNLO QCD accuracy

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Abstract

In these proceedings, we highlight some aspects of the recent computation of NNLO QCD corrections for W production in association with a charm jet at the LHC. The results are presented in the form of cross sections and differential distributions and are compared to ATLAS data.



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1 Introduction

The production of W boson in association with a charm jet at the LHC has raised significant interest in the recent years. Beyond being interesting on its own as another test of perturbative QCD, the process is very sensitive to the parton distribution function (PDF) of the strange/anti-strange quark in the proton. This is simply due to the direct link between W+c-jet production and strange/anti-strange quarks in the initial state as illustrated in the left-hand side of fig. 1. It means that this process constitutes a key ingredient for the determination of the strange/anti-strange quark PDF [1, 2]. To that end, several experimental analyses by both the ATLAS [3] and CMS collaborations [4–6] have been carried out.

On the theory side, next-to-leading order (NLO) QCD corrections have been known for quite some time for both the Tevatron [7] and the LHC [8]. Very recently, these corrections have been matched to parton shower [9]. In these proceedings, we briefly review some of the results of ref. [10], where the first next-to-next-to-leading order (NNLO) QCD corrections to W production in association with a charm jet at the LHC have been computed. In the original

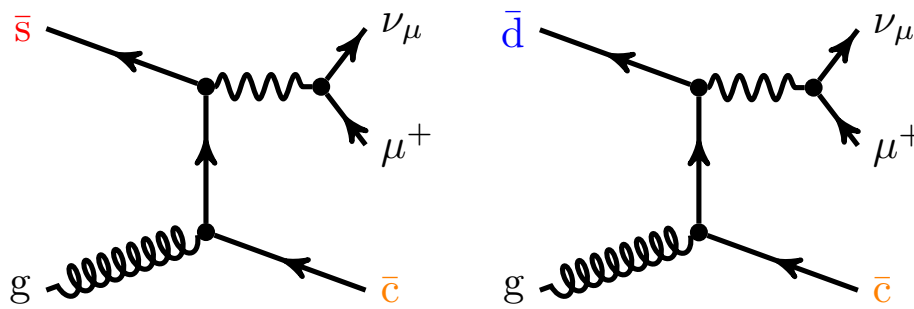


Figure 1: Feynman diagrams of $pp \rightarrow W^+j_c$ at LO with diagonal CKM matrix (left) and off-diagonal elements (right).

article, more details are provided such as predictions with different PDF sets, differential ratios between the W^+ and W^- signatures etc.

The structure of the proceedings reads as follow: in section 2, our best prediction is compared to ATLAS data [3] at 7TeV and the results are briefly discussed. In section 3, a short conclusion is provided. For more information, the interested reader is referred to the original article [10] as all the inputs utilised for the numerical simulations can be found there.

2 Results

In the introduction, the direct link between strange PDF and W +charm production at the LHC has been explained. Nonetheless, it is worth emphasising that this relation holds only at leading order (LO) when assuming a diagonal CKM matrix. Considering off-diagonal CKM elements (see the right-hand side of fig. 1) or including higher-order QCD corrections significantly complicates the situation by adding new partonic channels. This warrants therefore the precise computation of QCD corrections for this process. For the predictions presented here, the effect of $V_{cd} \neq 0$ is included at Born level only. Also, all theoretical predictions presented here have been obtained within the STRIPPER framework which is a c++ implementation of the four-dimensional formulation of the sector-improved residue subtraction scheme [11–14].

Following the ATLAS analysis [3], the phase-space definition reads

$$p_{T,\ell} > 20 \text{ GeV}, \quad |\eta_\ell| < 2.5, \quad p_{T,\text{miss}} > 25 \text{ GeV}, \quad m_T^W > 40 \text{ GeV}, \quad (1)$$

for the leptonic final states. In addition, one and only one charm jet should fulfil

$$p_{T,j_c} > 25 \text{ GeV}, \quad |\eta_{j_c}| < 2.5. \quad (2)$$

In fig. 2, the cross sections of $pp \rightarrow W^+j_c$ and $pp \rightarrow W^-j_c$ are compared to the measurements of the ATLAS collaboration [3]. In addition, the ratio of the two cross sections $R_{W^\pm j_c} = \sigma_{W^+j_c} / \sigma_{W^-j_c}$ is also provided. Based on the previous discussion, this ratio behaves approximately like $R_{W^\pm j_c} \sim (|V_{cs}|^2 \bar{s} + |V_{cd}|^2 \bar{d}) / (|V_{cs}|^2 s + |V_{cd}|^2 d)$, meaning that it also provides a sensitive probe of the strange-quark PDF.

The first interesting aspect to observe is that NLO QCD corrections are large while the NNLO ones are much more modest. This is by now a very well understood phenomenon for $V+j$ processes [15] and is also a sign of good perturbative convergence. As expected, the scale uncertainty is significantly decreasing when including higher orders. At NNLO, it becomes of the order of 2-3 per cent, meaning that it is smaller than the uncertainty due to PDF variation which is around 4%. Finally, the NNLO computations and the experimental data agree within their respective uncertainties. This holds true for both signatures as well as the ratio.

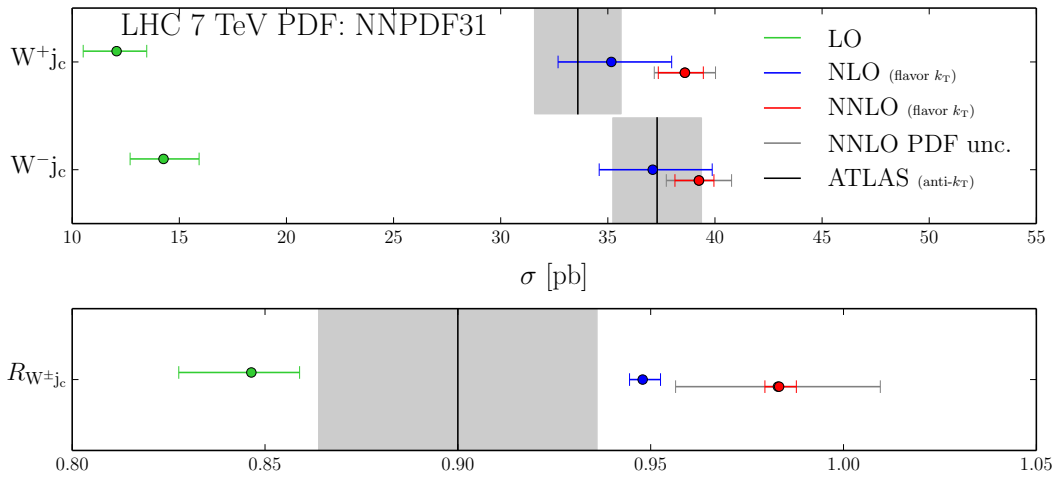


Figure 2: Cross sections for $pp \rightarrow W^+j_c$, $pp \rightarrow W^-j_c$, and the ratio $R_{W^{\pm}j_c}$ at the LHC with $\sqrt{s} = 7$ TeV. The theoretical predictions up to NNLO QCD accuracy are compared to the ATLAS data [3].

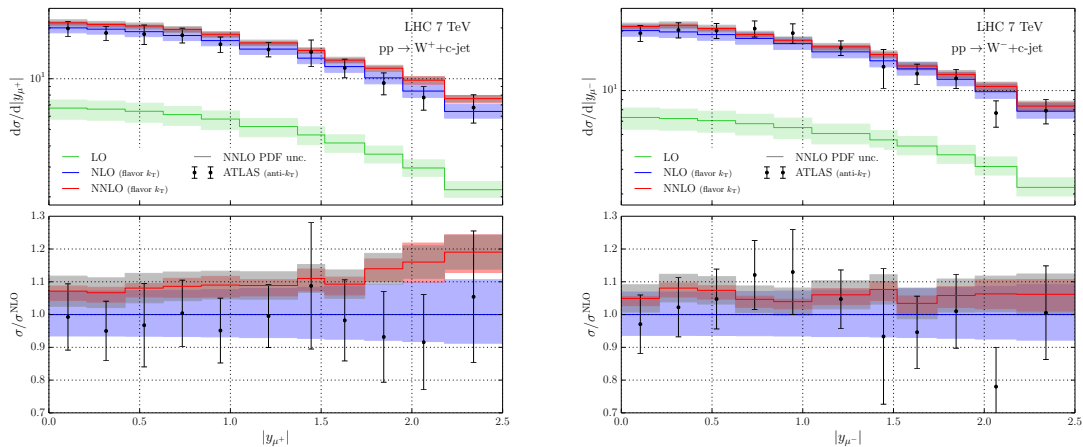


Figure 3: Differential distributions in the absolute rapidity of the muon/anti-muon in the process $pp \rightarrow W^+j_c$ (left) and of the muon in $pp \rightarrow W^-j_c$ (right) at the LHC with $\sqrt{s} = 7$ TeV. The upper panel shows the absolute predictions as well as the ATLAS data [3].

At the level of the differential distributions, the general picture is the same, with large NLO corrections and relatively moderate NNLO effects across the phase space. As an example, the differential distribution in the absolute rapidity of the muon/anti-muon is given in fig. 3. It is shown for both processes $pp \rightarrow W^+j_c$ (left) and $pp \rightarrow W^-j_c$ (right). The uncertainty due to PDF variation is again larger than the scale uncertainty obtained at NNLO accuracy. In the same way as for the fiducial cross section, the data-theory agreement is rather good over the whole kinematic range.

Nonetheless, it is worth mentioning a few limitations of this comparison. First, in our computation we have made use of the flavor- k_T algorithm [16] to ensure an infrared safe definition of the charm jets. On the other hand, the ATLAS measurement has been performed using the anti- k_T algorithm [17]. For the case of Z + b production, this mismatch has been estimated to be around 10% [18]. In the future, it is thus worth investigating such effect

for $W+c$ measurements. Second, the effects of $V_{cd} \neq 0$ has only been included at LO in our computation. At this order, it amounts to 5% to 10% depending on the signature. It is thus expected that such effects should amount to few per cent at higher orders. Finally, electroweak (EW) corrections have been here neglected. Due to Sudakov logarithms, these are usually around few per cent and grow negatively in high-energy limits. The EW corrections of order $\mathcal{O}(\alpha_s \alpha^3)$ have been found to be around -3% [19] while the subleading ones of order $\mathcal{O}(\alpha^4)$ are expected to be below a per cent [20].

3 Conclusion

In these proceedings we have reported on a recent computations of NNLO QCD effects for the $W+c$ -jet production at the LHC. In particular, we have focused on the main highlight of ref. [10] which is the comparison of our best predictions with the ATLAS measurement of $W+c$ -jet at 7 TeV [3]. Overall the agreement between theory and data is good. Nonetheless in order to perform precision comparisons, several aspects should be addressed in the future: the inclusion of off-diagonal CKM elements at higher order, the inclusion of EW corrections, and finally quantifying the effect of different jet algorithms for the charm jets identification. This work constitutes therefore only a first step toward determining with high precision the strange-quark content of the proton.

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