Report 1:

-Table 1 is misleading. Assign the top row the left label "Strange metal" - then the caption will match the Table.

Thank you for this suggestion. The modification has been made to the table.

-I suggest to more clearly organize the limiting factors of the method combined and in a separate subsection. I improved the reading of the manuscript.

A separate subsection related to the limiting factors was added in subsection II.D.

> The largest limiting factor of our method is the fermion sign problem that grows exponentially with the inverse temperature β and the free energy of the system [69]. This limits the maximal size of the cluster as well as the value of U. On top of the fermionic sign problem, we are also limited in temperature by the low acceptance rate of Monte-Carlo configurations at low temperatures. Because of that, it is impossible to reach T below 0.02 for the range of interactions we're interested in.
>
> Another limiting factor at high temperature comes from the method used to perform the analytical continuation of the observable. Indeed, as the temperature is increased, the interval between Matsubara frequencies increases, leading to inaccurate polynomial fits. For this reason, we limit ourselves to T lower than 0.2.
>
> For a typical value of hopping t in the cuprates of 0.3 eV, the range of temperature achievable with DCA would then be approximately between 70K and 700K. In BEDT organics the corresponding scales would be ten times smaller.
>
> Finally, DCA is a coarse-grained method. Because the observables are averaged over patches, the momentum resolution is limited.

-Similarly I suggest to make another subsection on a coherent and condensed argueing what one employs the triangular lattice Hubbard model in trying to explain linear -T behaviour on a square lattice.

As a T-linear regime with similarities to the two T-linear regimes found in this research is also found in the square lattice, we wanted to compare our results to those in the square lattice. The main motivation for this comparison is because of the fact that at $p = 25\%$, the behaviour of the scattering rate in both lattices are similar at low temperature. We added in the discussion Figure 7 for the scattering rate as a function of temperature for the triangulaire lattice ($|t'| = 1$) and for the square lattice ($t' = 0$), and concentrated the comparison to a separate subsection of the discussion. We are aware of the major difference between the two lattices, and we do not claim that the physics in the square and triangular lattice is the same.

> The results for the local scattering rate as a function of temperature on the square and the triangular lattices are presented in Fig. 7. We find that the effect of geometrical frustration does not affect the behaviour of the scattering rate at low temperature. This leads us to believe that the underlying physics responsible for the T-linear scattering rate is the same for both the square and triangular lattices. Based on this assumption, we compare our results for the triangular lattice to the characteristics of the strange metal phase found in cuprates.

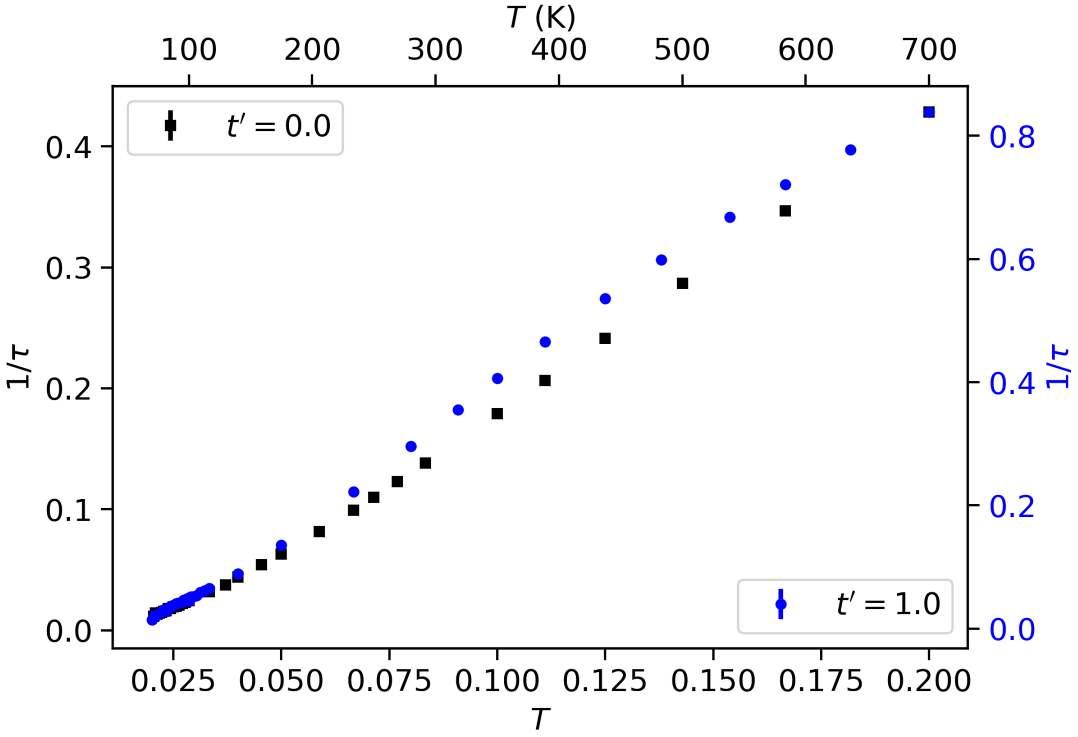

FIG. 1. Local scattering rate as a function of temperature for both the square lattice ($t' = 0$) and the triangular lattice ($t' = -t = 1$) for $p = 0.25$ and $U = 8.5$. The temperature scale is fixed by taking $t = 0.3$eV, typical of the numbers for cuprates.