
The Mu3e experiment

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Abstract

The *Mu3e* experiment aims for a single event sensitivity of $2 \cdot 10^{-15}$ on the charged lepton flavour violating $\mu^+ \rightarrow e^+e^+e^-$ decay. The experimental apparatus, a light-weight tracker based on custom High-Voltage Monolithic Active Pixel Sensors placed in a 1 T magnetic field is currently under construction at the Paul Scherrer Institute, where it will use the full intense $10^8 \mu^+$ /s beam available. A final sensitivity of $1 \cdot 10^{-16}$ is envisioned for a phase II experiment, driving the development of a new high-intensity DC muon source which will deliver $>10^9 \mu^+$ /s to the experiment.

20.1 Introduction

Searches for Charged Lepton Flavour Violation (CLFV) in muon decays are a remarkably sensitive method to search for new physics processes [1]. These decays are free from Standard Model backgrounds, and leave a relatively simple and clear signature in the experimental apparatus. In addition, intense muon beams are available at several facilities, where the relatively long-lived muons get transported from a production target to an experimental area.

The Paul Scherrer Institute (PSI) has been at the forefront of CLFV searches, with the current best limit on the $\mu^+ \rightarrow e^+\gamma$ decay channel of $4.2 \cdot 10^{-13}$ (90% CL) from the *MEG* experiment [2]. The *SINDRUM* experiment [3] set the best limit on the $\mu^+ \rightarrow e^+e^+e^-$ decay channel, and the *SINDRUM II* experiment [4] on muon conversion $\mu^- \rightarrow e^-$ on gold. A new generation of experiments pursuing these three *golden* channels, which probe for new physics in a complementary manner [5], is currently under construction: the *Mu2e* experiment at Fermilab, the *COMET* experiment at J-PARC, and the *MEG II* experiment at PSI. The *Mu3e* experiment aims for a 10^{-16} single-event sensitivity for the $\mu^+ \rightarrow e^+e^+e^-$ CLFV decay channel, an improvement by four orders of magnitude compared to the limit set by the *SINDRUM* experiment [3]. Such a leap in sensitivity is enabled by the availability of high-intensity muon beams, the use of silicon pixel detectors instead of multi-wire proportional chambers to track the decay products, and a modern data-acquisition system able to handle the vast amount of data produced by the detector at high beam rates. A first phase of the experiment is currently under construction at the $\pi E5$ beamline at PSI, where the intense DC surface muon beam of $10^8 \mu^+$ /s will be exploited to achieve a single event sensitivity of $2 \cdot 10^{-15}$ in 300 days of data taking [6].

The *Mu3e* detector is optimized for the $\mu^+ \rightarrow e^+e^+e^-$ decay. It is designed to track the two positrons and one electron from muons decaying at rest with a light-weight tracker placed inside a 1 T magnetic field, thereby reconstructing the decay vertex and invariant mass. The

41 momentum balance of the three reconstructed particles should consistent with a muon decay-
 42 ing at rest. Several background processes can potentially meet the same criteria as the recon-
 43 structed signal events. The dominating accidental background originates from the overlay of
 44 two ordinary muon decays where one of the positrons produces an additional electron track
 45 through Bhabha scattering in the target material. This process is sufficiently suppressed by
 46 means of a good vertex resolution of better than $300 \mu\text{m}$, a timing resolution of a few 100 ps,
 47 the requirement of an invariant mass equal to the muon mass, and a balanced momentum
 48 budget. Additional background background from $\mu^+ \rightarrow e^+ e^+ e^- \nu_e \bar{\nu}_\mu$ internal conversion de-
 49 cays can only be suppressed by means of an excellent momentum resolution of $\sigma_p < 1 \text{ MeV}$,
 50 as shown in Figure 20.1).

51 All *Mu3e* detector sub-systems, as described in Section 20.2, are currently under construc-
 52 tion. With the solenoid magnet (Figure 20.2) installed at PSI, the first engineering runs are
 53 planned for spring 2021.

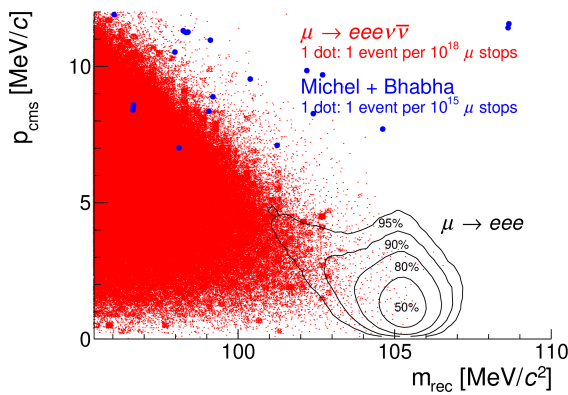


Figure 20.1: The simulated reconstructed mass versus the momentum balance of two positrons and one electron from a common vertex [6]. The accidental background is shown in blue, the dominating background from internal conversion is shown in red.

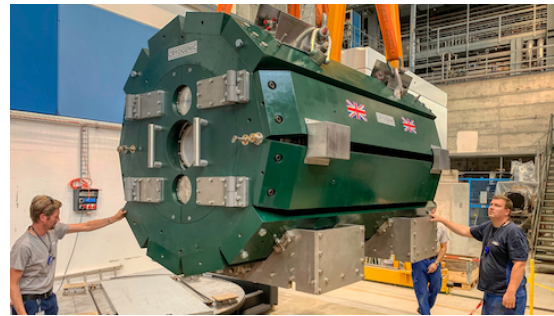


Figure 20.2: The 30 ton *Mu3e* magnet arriving at PSI. The magnet is currently installed and commissioned in the πE5 experimental area, providing a magnetic field of up to 2.6 Tesla with a $\frac{\Delta B}{B}$ uniformity and stability of $\mathcal{O}(10^{-4})$.

54 20.2 The *Mu3e* detector

55 The *Mu3e* detector is located at the Compact Muon Beam Line at the πE5 channel. After the
 56 positron contamination from the beam is removed by a Wien filter, the surface μ^+ beam of
 57 up to $10^8 \mu^+/\text{s}$ is transported to the center of the *Mu3e* solenoid magnet, and stopped on
 58 a hollow double-cone target, which spreads out the decay vertices in z and minimises the
 59 amount of target material traversed by the decay particles. The target is surrounded by the
 60 cylindrical central tracker, consisting of the inner silicon pixel detector, a scintillating fibre
 61 tracker for time measurements, and the outer silicon pixel detector. A momentum resolution of
 62 better than 1 MeV/c is achieved by letting the positrons(electrons) recurl in the magnetic field,
 63 either crossing the central tracker again, or hitting the outer tracking stations surrounding the
 64 upstream and downstream beam pipe. These stations consist of a silicon pixel tracker, and
 65 a scintillating tile detector mounted on the inside of the pixel tracker. The 5 mm thick tiles
 66 enable a time resolution for the tracks reaching these outer stations of better than 100 ps. The
 67 active part of the *Mu3e* detector is depicted in Figure 20.3.

68 As multiple Coulomb scattering is the dominating factor affecting the momentum resolu-
 69 tion, it is crucial to minimize the material budget in the tracking detectors. For this purpose,

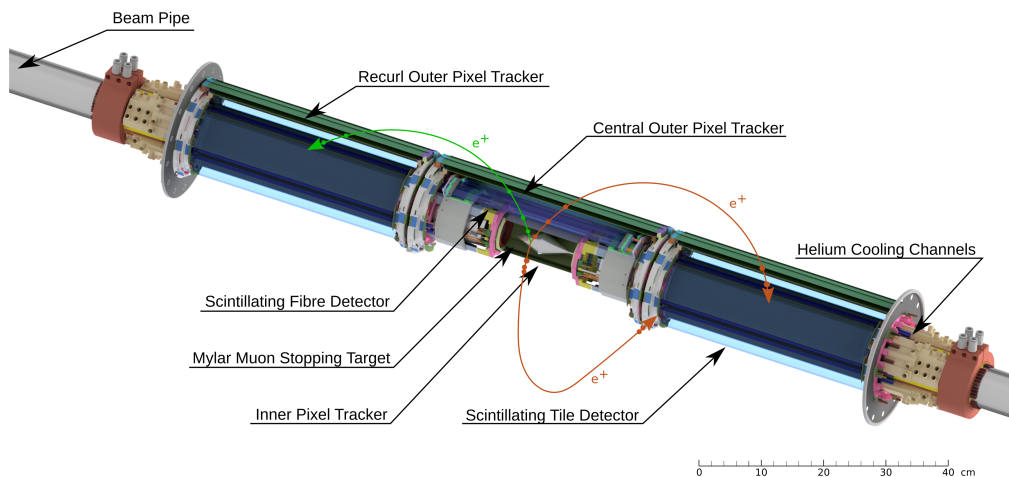


Figure 20.3: The active part of the *Mu3e* detector, with a central tracker surrounding the target, and upstream and downstream outer pixel tracking stations. The large lever arm created by the recurling tracks enables the high momentum resolution required.

70 the collaboration has developed a custom High-Voltage Monolithic Active Pixel Sensor [7]
 71 (HV-MAPS) based on a commercial 180 nm HV-CMOS process. After a series of prototypes
 72 showed good efficiency ($>99\%$) and time resolution ($\mathcal{O}(10\text{ ns})$) [8] [9], the current version,
 73 , the final *MuPix* sensor is a $2 \times 2\text{ cm}^2$ sensor with $80 \times 80\ \mu\text{m}^2$ active pixels, thinned to $50\ \mu\text{m}$
 74 (Figure 20.4). The digital periphery provides up to three 1.25 Gbit/s Low-Voltage Differential
 75 Signaling (LVDS) continuous data connections to the front-end electronics. The sensors are
 76 bonded to a thin aluminum/polyimide flex print carrying all electrical signals. Together with
 77 a polyimide support structure, the entire silicon tracking module has a thickness of ca. 0.0012
 78 radiation lengths. The pixel sensors generate about 250 mW/cm^2 of heat. To remove this
 79 heat whilst keeping the material budget of the tracker sufficiently low, a gaseous He cooling
 80 system [10] is deployed providing well controlled He flows at atmospheric pressure between
 81 and outside the pixel layers.

82 A time resolution of about 10 ns is insufficient to determine the direction and thus the
 83 charge of the decay particles. A scintillating fibre detector is therefore placed between the
 84 inner and outer layer of the central silicon-pixel tracker, consisting of a dozen 30 cm long rib-
 85 bions made from three staggered layers of $250\ \mu\text{m}$ diameter multiclad round fibers, read out by
 86 Silicon Photomultipliers (SiPM) arrays on both sides [11]. Located at the very end of the re-
 87 curving particle trajectories hitting the upstream or downstream tracker, where the constraints
 88 on the material budget are less stringent, the tile detector provides the needed precise timing
 89 information of the particle tracks, in conjunction with the fibre detector significantly reducing
 90 the accidental background associated with the intense rate of $10^8\ \mu^+/\text{s}$. Each of the 5824
 91 individually wrapped tiles is read out by a single SiPM. Both the fibre and tile SiPM signals are
 92 processed by a custom Application-Specific Integrated Circuit (ASIC), the 32 channel *MuTrig*
 93 chip [12], which applies 2 thresholds to the analogue signal for time and energy information.
 94 The *MuTrig* chip has a 1.25 Gbit/s LVDS data connection, similar to the *MuPix* chip readout.
 95 For tile and fibre detector a respective time resolution of $<50\text{ ps}$ and $<400\text{ ps}$ is achieved.

96 The entire *Mu3e* detector is mounted in the bore of a superconducting magnet. Figure 20.2
 97 shows the 3 m long solenoid magnet with the iron return yoke. It has a 1 m wide bore housing
 98 the active detector, in addition to the support structures and services such as the front-end
 99 readout electronics and DC-DC power converters for the detector ASICs. The two flanges

100 below and above the beam pipe provide access for the water and gaseous helium cooling
101 pipes, the power cables, and the optical data connections.

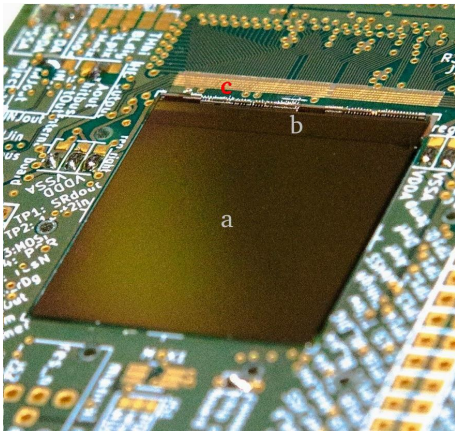


Figure 20.4: The full sized MuPix sensor, with a) a $2 \times 2 \text{ cm}^2$ sized active area, and b) a periphery with the pixel hit digitization and read-out state machine. This chip is c) wire bonded to a PCB for testing purposes.

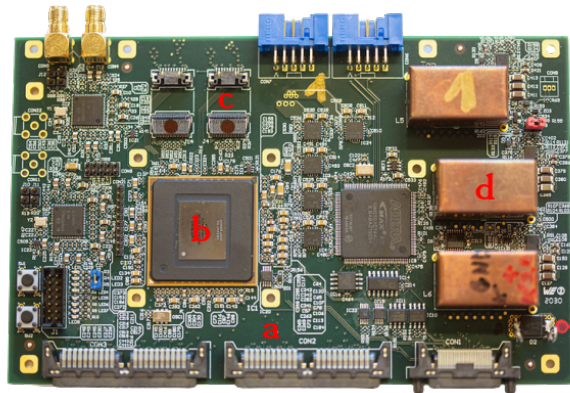


Figure 20.5: The front-end readout board, combining and time sorting the a) data from up to 36 detector ASICs on b) an Arria V FPGA, before sending the data via the c) optical Samtec FireFly transceivers. d) Custom DC-DC converters with air coils regulate the power on the board.

102 20.3 Readout and online event selection

103 With three lepton tracks going in different (opposite) directions, the topology of a $\mu^+ \rightarrow e^+e^+e^-$
104 event is such, that a global picture of the detector is needed before candidate events can be
105 selected. This leads to a trigger-less readout scheme as shown in Figure 20.6, where all pixel,
106 fibre and tile hits are continuously being digitized and merged into a data stream of up to 100
107 Gbit/s. A series of PC's housing powerful Graphics Processing Units (GPU) perform an online
108 event-selection, reducing the data rate to a manageable 50-100 MByte/s which is stored for
109 further offline processing.

110 Each detector ASIC, a *MuTrig* or *MuPix* chip, assigns a timestamp and address to each hit,
111 and sends the serialized data through a series of flex-prints and twisted pair cables to a front-
112 end board (Figure 20.5). Each of these readout boards is located inside the magnet bore and
113 accepts up to 45 electric LVDS links. The data streams are merged and time-sorted on an Arria V
114 Field-Programmable Gate Array (FPGA). Two optical transceivers provide eight 6 GBit/s links
115 to the outside, sending off the merged and sorted hit information combined with the slow-
116 control data. In addition, the front-end FPGA also configures the detector ASICs, including
117 tuning the very large number of individual *MuPix* pixels, and distributes the clock and reset
118 signals.

119 All incoming and outgoing data connections to and from the detector volume travel via
120 optical fibres to the counting house. The data links from the 112 front-end boards are con-
121 nected to the *Switching boards*, where the data from different detector modules are merged
122 into 64 ns time slices containing the full detector hit information. This custom *PCIe40* board
123 housing a large Arria 10 FPGA and 48 fast optical receivers and 48 fast optical transmitters
124 was developed for the LHCb and ALICE upgrades [13].

125 The online event selection must decide which of these 64 ns *snapshots* of the detector to
126 store for later (offline) processing, in the process keeping less than 1% of the data. A simple
127 time coincidence between 3 tracks is insufficient to achieve this. Instead an online filter farm

128 reconstructs all tracks in software, and performs the selection by requiring 3 tracks having a
 129 common vertex and the kinematics of a possible $\mu^+ \rightarrow e^+e^+e^-$ event. The filter farm consists of
 130 12 PC's housing a FPGA board receiving the data and a powerful commercial GPU performing
 131 the event selection. With simple geometric cuts, candidate tracks are first selected on the FPGA
 132 from hits in the central pixel tracker. The track fitting [14] is performed on the GPU, where
 133 $1 \cdot 10^9$ fits per second have been achieved on a NVIDIA GTX 980 GPU, sufficient to be able to
 134 process the expected 10^8 muon decays/s. A newer more powerful GPU will be selected when
 135 equipping the farm PCs.

136 The MIDAS¹-based data-acquisition system sends the filtered data to on-site and off-site
 137 storage for later processing. This integrated DAQ also takes care of the configuration, mon-
 138 itoring, and logging of all parameters of the detector and its services such as the water and
 139 helium cooling system and power distribution.

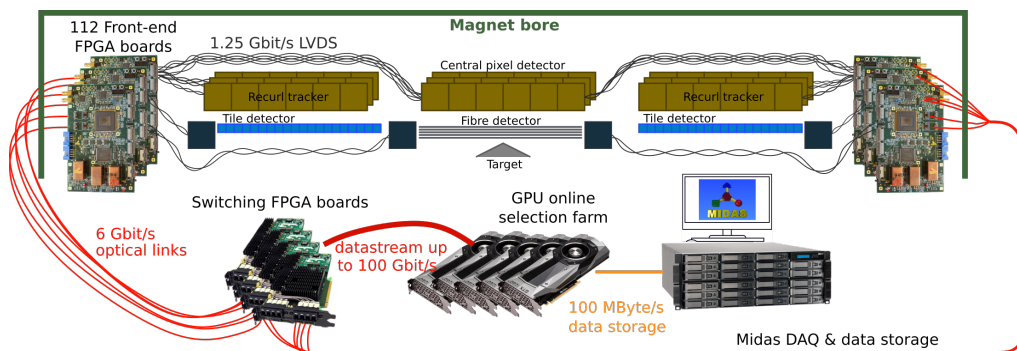


Figure 20.6: A sketch of the *Mu3e* triggerless readout scheme [15], where all detector hits are piped to the online filter farm. A selection algorithm based on massive parallelised track fitting sends off a subset of the data for further offline processing.

140 20.4 Conclusions and outlook

141 With the magnet installed at the Paul Scherrer Institute, the *Mu3e* experiment is entering
 142 its construction phase. All sub-detector demonstrators have met the required specification,
 143 and are currently being integrated to a single lightweight electron/positron tracker. This also
 144 includes a novel read-out system of the apparatus, which pipes the full detector information to
 145 an online filter farm. Aside from being a necessary requirement set by the CLFV decay event
 146 topology, this readout scheme where the full and global detector information is available for
 147 online analysis, also allows other new-physics searches such as CLFV two-body decays and
 148 Dark Photon searches [16].

149 The *Mu3e* phase II experiment envisions a branching ratio sensitivity of $1 \cdot 10^{-16}$. Many
 150 detector sub-systems are already designed with this goal in mind, but significant research
 151 and development on the detector side still has to be done. An order of magnitude increase in
 152 sensitivity also requires a more intense, and currently unavailable muon flux of μ^+ /s of $\mathcal{O}(10^9)$.
 153 For this purpose, a new High-Intensity Muon Beamline [17] to be installed at the target M is
 154 currently under development at the Paul Scherrer Institute, replacing the conventional muon
 155 extraction beamline elements with solenoids. The timeline of this project coincides with the
 156 envisioned start of the *Mu3e* Phase II construction at the end of this decade.

¹<https://midas.triumf.ca>

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