

# MUSE: The MUon Scattering Experiment

E. Cline<sup>1</sup>, J. Bernauer<sup>1,2</sup>, E.J. Downie<sup>3</sup> and R. Gilman<sup>4\*</sup>

<sup>1</sup> Stony Brook University, Stony Brook, NY

<sup>2</sup> Riken BNL Research Center, Upton, NY

<sup>3</sup> The George Washington University, Washington, DC USA

<sup>4</sup> Rutgers, The State University of New Jersey, Piscataway, NJ

\* rgilman@physics.rutgers.edu

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## Abstract

MUSE is a high-precision muon scattering experiment aiming to determine the proton radius. Muon, electron, and pion scattering will be measured at the same time. Two-photon exchange corrections will be determined with data using both beam polarities.

## 23.1 Introduction

The charge radius is a fundamental property of the proton. It is of interest to hadronic physicists as a test of calculations of proton structure. It is of interest to atomic physicists as it affects the determination of the Rydberg constant, and so is important in precision tests of quantum electrodynamics.

The charge radius can be determined using electromagnetic interactions in two ways. In atomic physics, the proton size changes the energies of S states by

$$\Delta E = \langle \Psi_S | \delta V | \Psi_S \rangle = \frac{2}{3} \pi \alpha |\Psi_S(0)|^2 r_p^2, \quad (23.1)$$

thus allowing the radius and Rydberg constant to be determined simultaneously by measuring pairs of transition energies. In electron-proton scattering, the differential cross section depends on the square of the form factor, which is the momentum-space charge distribution. The charge radius is extracted from the slope of the electric form factor  $G_E$  at  $Q^2 = 0$ :

$$r_p^2 = -6 \frac{dG_E}{dQ^2} \Big|_{Q^2=0}. \quad (23.2)$$

As the scattering data do not extend to  $Q^2 = 0$ , the radius is extracted from fits to measured cross sections.

In 2010 the proton charge radius was determined to be  $0.84184 \pm 0.00067$  fm from a measurement of muonic hydrogen by the PSI CREMA collaboration [1]. This was quite puzzling as it was about  $5\sigma$  smaller than the nearly order-of-magnitude less precise electronic measurements [2], which used both hydrogen spectroscopy and electron-proton scattering. This proton radius puzzle was quickly confirmed with reports from two new electron scattering measurements yielding  $r_p = 0.879 \pm 0.008$  fm [3] and  $0.875 \pm 0.010$  fm [4], and a second measurement of muonic hydrogen [5] that found  $r_p = 0.84087 \pm 0.00039$  fm. New data are needed to resolve the proton radius puzzle, and a number of new experiments were developed [6–9]. Most aim to improve existing results, with new measurements of atomic hydrogen or electron-proton scattering. A new set of muonic atom measurements were also undertaken with other light nuclei.

## 38 23.2 The MUSE experiment

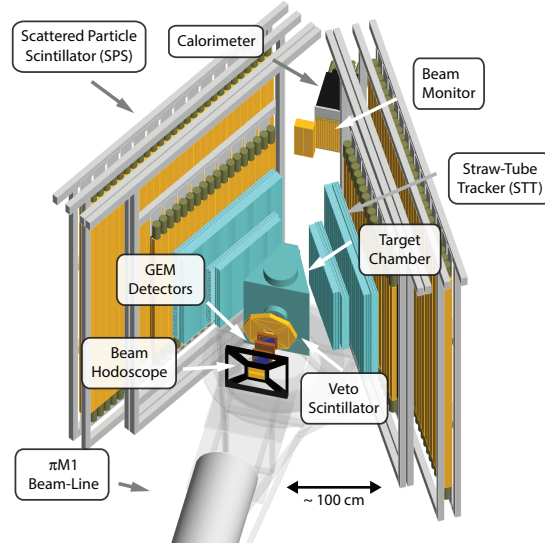


Figure 23.1: The MUSE experimental system. See text for details.

39 The MUon Scattering Experiment (MUSE) addresses the radius puzzle in a unique way. The  
 40 intent is to extract the first precise proton radius measurement from muon-proton scattering.  
 41 The experiment uses the PSI HIPA PiM1 channel [10, 11], which provides a secondary beam of  
 42 pions, muons, and electrons. This enables simultaneous measurements of both electron and  
 43 muon scattering, so that the extracted proton radii and the cross sections for the two reactions  
 44 can be directly compared. The PiM1 channel can produce beams with similar beam properties  
 45 for both polarities. A difference between the scattering probability for the two beam polarities  
 46 would result from two-photon exchange, a higher-order correction to the interaction. This  
 47 correction is expected to be small,  $O(0.1 - 1\%)$ , depending on kinematics, but it is difficult to  
 48 calculate accurately. It might affect the determination of the radius.

49 Figure 23.1 shows the experimental apparatus, taken from the MUSE Geant4 simulation.  
 50 Beam particles exiting the channel first pass through a beam hodoscope, which measures par-  
 51 ticle times. In conjunction with the accelerator RF signal, these times can be used to de-  
 52 termine particle species. The beam next passes through GEM chambers, which measure the  
 53 beam-particle trajectories. A veto scintillator is used to suppress background events such as  
 54 upstream beam particle decays in flight or scattering from the detectors, leading to particles  
 55 passing through the vacuum chamber wall. The target system inside the vacuum chamber  
 56 includes a liquid hydrogen cell, an empty cell, solid targets, and a beam focus monitor. The  
 57 unscattered beam exits through a thin window, and reaches the downstream beam monitor and  
 58 a calorimeter, which are used to study radiative corrections. Scattered particles exit through  
 59 thin side windows, are tracked by the straw tube tracker, and their times measured with the  
 60 scattered particle scintillators.

61 The PiM1 channel has been used previously for precise pion scattering measurements. This  
 62 is feasible as pions are often the dominant species in the beam, and hadronic scattering cross  
 63 sections can be orders-of-magnitude larger than electromagnetic cross sections. A primary  
 64 challenge of MUSE is to measure precise cross sections for the smaller muonic component of

65 the beam. The first aspect of the challenge is that previous determinations of beam proper-  
 66 ties concentrated on the pionic component of the beam, so the properties of the muonic and  
 67 electronic components are not as well known. The second aspect is that the experimental sys-  
 68 tem has to largely prescale away pion scattering to be able to efficiently measuring muon and  
 69 electron scattering.

70 To address the challenge of beam properties, MUSE has undertaken a program of simula-  
 71 tions and measurements. The first step is to simulate the particle production mechanisms at  
 72 the M target. Charged pions are produced at the M target through  $pC \rightarrow \pi^\pm X$  reactions. From  
 73 the perspective of the PiM1 channel, the proton beam crosses the M target generating pions  
 74 with an effective millimeter-sized source. Muons are produced by the decays in flight of those  
 75 pions. Simulations show that the majority of the muons that will pass through the PiM1 chan-  
 76 nel are generated by pions that decay in the first few centimeters of flight, at an angle of nearly  
 77  $90^\circ$  in the pion rest frame. The effective muon source size is larger than the pion source size,  
 78 but still only a few millimeters. Electrons and positrons are produced mainly by a sequence  
 79 of reactions, with  $pC \rightarrow \pi^0 X$  producing neutral pions, followed by the decay  $\pi^0 \rightarrow \gamma\gamma$ , and  
 80 subsequently pair production in the M target via  $\gamma C \rightarrow e^\pm X$ . Geant4 simulations show that  
 81 higher momentum electrons and positrons are only produced when all these processes are in  
 82 the direction of the PiM1 channel. As a result, the effective source size remains very close to  
 83 that for pions.

84 The source simulations generate charged particles that are input to the TURTLE [12] and  
 85 G4 beamline [13] magnetic transport codes. These codes include the channel quadrupoles and  
 86 dipoles, as well as apertures from beam pipes and jaws. The simulation describes well several  
 87 measured properties of the beam, including the beam distributions in position and angle at  
 88 the channel intermediate focal plane and at the scattering target position, and the variation of  
 89 particle times at the scattering target with respect to accelerator RF as a function of momen-  
 90 tum: the pion time distribution is wider than that for electrons or muons due to the interplay  
 91 of faster speed vs longer flight path for higher-momentum particles within the channel. While  
 92 the measured time distributions of all particles are quite similar, the muon distribution is pre-  
 93 dicted to be somewhat larger than the pion and electron distributions, indicating that extreme  
 94 rays are more constrained in reality than in the simulation.

95 In addition to the particle trajectories, it is important to know the beam momentum at the  
 96 0.2% (0.3%) level for muons (electrons). The channel momentum resolution is better than  
 97 this. The absolute momentum of the beam selected by the PiM1 channel is determined in  
 98 3 ways. First, dedicated time-of-flight measurements with changes of the beam hodoscope  
 99 and beam monitor positions determine the pion and muon momenta to the 0.2 – 0.3% level.  
 100 Second, the timing of particles in the beam hodoscope relative to the accelerator RF provides  
 101 an independent momentum measurement at the same level.<sup>1</sup> Third, the dispersion of the  
 102 channel at the intermediate focal point, of 7 cm/%, combined with the dispersion of the beam  
 103 from the intermediate focus to the scattering target of  $\approx 9.5$  cm/%, provides a check of any  
 104 momentum difference between the different particle species at the  $\approx 0.1\%$  level, through the  
 105 similarity of the measured beam spot positions.

106 The challenge of suppressing pion scattering while efficiently measuring muon and electron  
 107 scattering is addressed by the MUSE trigger system. A first-level trigger FPGA identifies all  
 108 particle species in the 3.5-MHz beam using the time difference between the beam-hodoscope  
 109 signal and the accelerator RF signal. Other first-level triggers identify scattered particles and  
 110 hits in the veto detector. The combination of these first-level triggers allows muon and electron  
 111 scattering to be read out efficiently while suppressing pion scattering.

112 One important feature of MUSE will be the implementation of a blinded analysis in the  
 113 cross section measurement. A Monte Carlo simulation is needed to determine precise cross sec-

<sup>1</sup>This timing measurement also checks the beam momentum stability at the  $\approx 0.1\% - 0.2\%$  level.

114 tions, and from them the proton radius. The blinding will be accomplished primarily through  
 115 modifying the simulation-derived weight factor, while encrypting the actual weights. Addi-  
 116 tionally, some small fraction of the tracks for different particle species will be thrown away  
 117 as a function of angle, to prevent accidental unblinding by direct comparison of charge and /  
 118 or particle species. This will be programmed to be reversed by the application of two encryp-  
 119 tion keys. Once the analysis is complete, the actual weights can be extracted and the physics  
 120 analysis rerun.

121 A more detailed description of the MUSE system is available in [14]. Detailed publications  
 122 are also available for the target [15] and the SiPM detectors [16].

### 123 23.3 Anticipated results

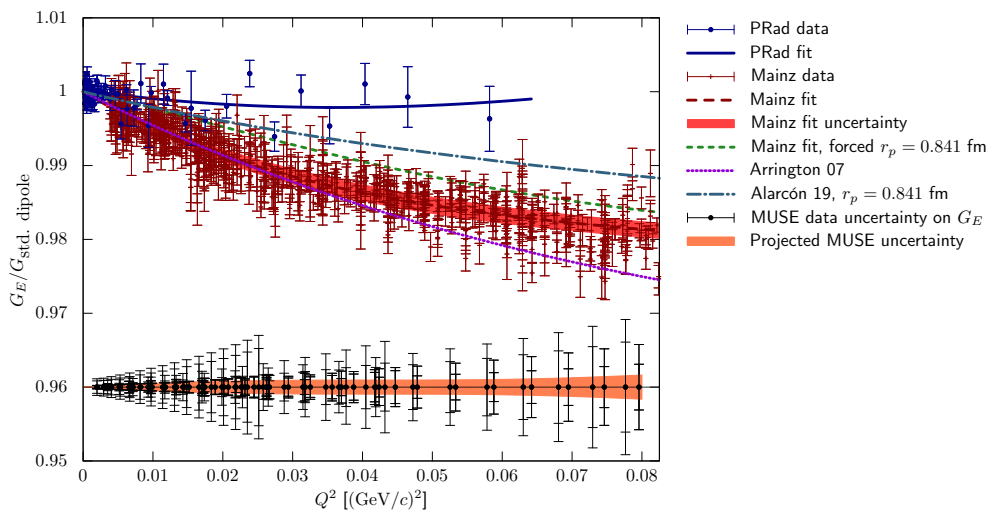


Figure 23.2: Anticipated data for  $G_E$  from MUSE, arbitrarily placed at 0.96, compared to recent electron scattering experiments, and fits to these data, and to two world data fits. The MUSE data include both electron and muon points. The doubled uncertainty bars represent the uncertainties for + (inner bar) and - (outer bar) beam polarity. The muon and electron points are slightly offset due to the mass difference of muons and electrons. See text for further details.

124 With the planned 12 months of beam time,  $4 \times 10^7 \mu^+$  ( $2 \times 10^7 \mu^-$ ) scattering events are  
 125 expected for MUSE. This should give better than 1% statistical precision for the cross section  
 126 in almost all of the 16 planned angle bins at each of 3 beam momenta and two beam polarities.  
 127 Figure 23.2 shows the expected uncertainties for the determination of the electric form factor,  
 128  $G_E$ , from MUSE, together with the results from Mainz [3] and from PRad [17], along with two  
 129 selected fits [18, 19]. The Arrington07 fit [18] is to older world data that are not shown, and  
 130 has a large radius. The Alarcon19 curve [19] is a dispersively improved effective field theory  
 131 calculation which has one free parameter, the radius, which can be fit, but here is chosen to  
 132 be the muonic spectroscopy value. The green dashed “Mainz-fit” line is a fit to the Mainz data,  
 133 but with the radius term set to the muonic spectroscopy value.

134 The experiments each measure in different kinematic regions, with MUSE at the lowest  
 135 beam momentum and largest angles, and PRad at the highest beam momentum and smallest  
 136 angles. The experiments also use different techniques. The more recent PRad measurement  
 137 used a forward angle calorimeter to measure cross sections for 1.1 and 2.2 GeV beam energies  
 138 at angles up to  $\approx 7.5^\circ$ . The earlier Mainz measurements used magnetic spectrometers at larger

139 scattering angles, with beam energies from 180 – 855 MeV. The Mainz and PRad data can be  
140 seen to diverge from each other, which probably indicates problems either with the experi-  
141 ments or with the radiative corrections. While the Mainz data are in good agreement with the  
142 Arrington fit to earlier data, neither the PRad nor the Mainz data agree with the prediction  
143 by Alarcon using the muonic radius. The expected MUSE uncertainties are competitive with  
144 those of the existing experiments. Muon scattering has much smaller single-photon radiative  
145 corrections, due to the larger muon mass, so any differences between muons and electrons  
146 might point to issues of radiative corrections or new physics.

147 The comparison of the cross sections for + and – polarities will yield a measurement of  
148 the two-photon exchange contribution, expected to be of similar size to the experimental un-  
149 certainties shown in Figure 23.2. The proton radius should be determined with an uncertainty  
150 of 0.006 – 0.010 fm, based on a sample of fits. The electron scattering data will have supe-  
151 rior statistical precision, but larger systematic uncertainties due to radiative corrections. This  
152 should result in slightly better measurements for both the radius and the two-photon exchange  
153 contribution.

154 In addition to the electromagnetic scattering, pion cross sections need to be measured dur-  
155 ing MUSE to sufficiently characterize experimental backgrounds. The pion cross sections are  
156 interesting by themselves as a test of the application of chiral perturbation theory, to improve  
157 the existing  $\pi N$  scattering database, and as a constraint on occasional speculations about  
158 undiscovered resonances in the  $\pi N$  system. Because MUSE operates with a mixed beam, pion  
159 scattering will be measured in all MUSE kinematics at the same time as the electron and muon  
160 scattering. The experimental trigger includes beam particle information, which allows the  
161 pion scattering events to be pre-scaled to become a small fraction of the data set, while still  
162 recording on the order of  $10^7$  events.

## 163 23.4 Outlook

164 A test of the full MUSE system in December 2019 led to several planned upgrades to make  
165 the system more robust. Due to the ongoing international public health crisis and its resulting  
166 impact on international travel, we were only able to partially complete the upgrades during  
167 2020. We plan to complete the upgrades and start MUSE production data taking in 2021.  
168 With 12 months of data taking and analysis to be performed, we anticipate publication of  
169 first results in 2023/24. MUSE will be the first experiment to measure elastic muon-proton  
170 scattering in an appropriate kinematic region, with a precision sufficient to address the proton  
171 radius puzzle. The corresponding results for the simultaneously-measured electron scattering,  
172 will put a strong constraint on potential systematic uncertainties, and may help settle the  
173 discrepancies between the Mainz and PRad results. MUSE will be the only experiment that  
174 can directly measure with its own data the difference between electron and muon extractions  
175 of the radius, making it highly compelling.

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