

ATLAS LAr Calorimeter commissioning for LHC Run-3: Energy computation in LATOME boards

Luka Selem^{1*}, on behalf of the ATLAS Liquid Argon Calorimeter Group ¹

¹ LAPP, Univ. Savoie Mont Blanc, CNRS/IN2P3, Annecy; France

* luka.selem@lapp.in2p3.fr

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Abstract

The ATLAS experiment at the Large Hadron Collider (LHC) measures proton–proton collisions at high energies. Within the Phase-I upgrade of the LHC before the start of Run-3 in 2022, the trigger system of the Liquid Argon calorimeter of ATLAS is being prepared to cope with an increased number of simultaneous proton–proton collisions. In the back-end of this new trigger system, the LATOME boards will be responsible for the computation of the energies deposited in the calorimeter. The commissioning of this computation within the LATOME is presented.

1 Introduction

The ATLAS experiment [1] is a multi-purpose detector installed at the Large Hadron Collider (LHC) [2] operated at CERN. It aims at collecting as much information as possible from the products of the high energy proton-proton collisions provided by the LHC every 25 ns. For the particles produced by these collisions, such information, collected by various sub-detectors, can be their nature, their charge, their momentum or their energy. The Liquid Argon (LAr) calorimeter is a sub-detector measuring mainly the energy of electrons, positrons and photons. It consists in alternately layers of absorbing material, mainly lead – favouring the showering of the incoming particles – and liquid Argon where low energy particles of the shower ionize producing an electrical signal. Particle positions are defined in coordinates of their pseudo-rapidity η and the azimuthal angle ϕ of the ATLAS detector seen as a barrel. The LAr calorimeters are divided in cells corresponding to specific (η, ϕ) coordinates and one of the four layers in depth. The electrical signal coming from each cell is used to recompute the energies deposited at specific position in the detector.

With collisions at a frequency of 40 MHz, it is impossible to keep every event and triggers are used to record only the physically interesting ones. The level-1 hardware trigger of the LAr calorimeter output frequency should be 100 kHz, reducing the data by a factor 400. The allocated latency for the computation of energies of particles is then of approximately 150 ns. To meet such requirements, the main idea is to reduce the number of calorimeter cells for which a deposited energy has to be computed. This is achieved with Trigger Towers, which are groups

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30 of LAr cells combining around 60 of them in depth and in a region of $\Delta\eta \times \Delta\phi = 0.1 \times 0.1$.
 31 The Trigger system groups the electrical signal coming from cells of a trigger tower and the
 32 reduced granularity allows to compute fast enough the corresponding energies.

33 Within the Phase-I upgrade of ATLAS [3] before Run-3, this trigger system is modified in
 34 order to cope with the increased pile-up in Run-3. The mean number of collision per bunch
 35 crossing is expected to go from $\langle\mu\rangle \approx 20$ to $\langle\mu\rangle \approx 80$ in Run-3. To keep the same level-1 accept
 36 rates, the trigger discriminating power has to be improved. This is achieved by increasing the
 37 granularity. For this purpose, Super-Cells are defined as a new grouping of LAr cells. Typically,
 38 there are 10 Super-Cells per Trigger Tower. However, this increased granularity puts more
 39 pressure on the trigger system as more energies have to be computed within the same time
 40 constraints. This is made possible thanks to progress in electronics and the use of FPGAs. In
 41 the end, the full front-end and back-end of the trigger system have to be changed as described
 42 in the red part of figure 1a.

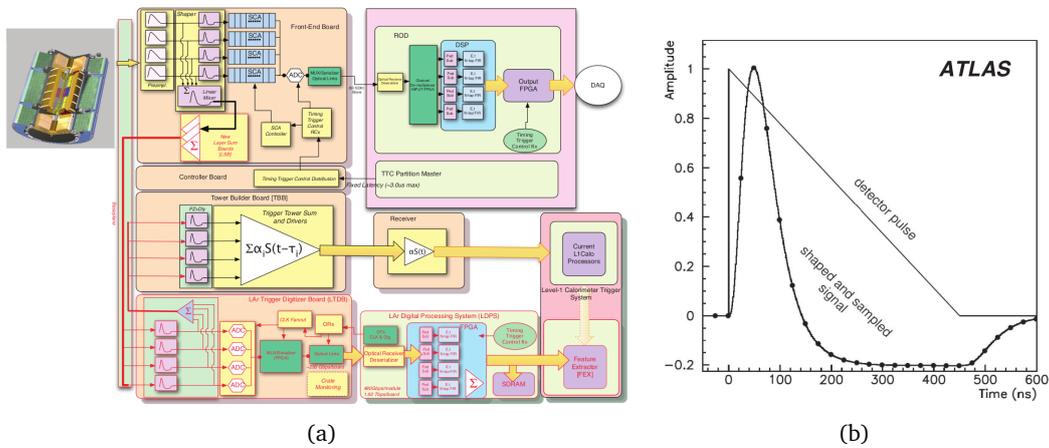


Figure 1: (a) Schematic block diagram of the LAr trigger readout architecture. Blocks concerned by the Phase-1 upgrade have a red outline. (b) Triangular signal shape from the LAr calorimeter, overlaid with the bipolar-shaped and sampled signal. Both figures are from reference [3].

43 2 From LAr pulse to Energy

44 The electrical signal coming from the LAr cells is typically a triangle as shown in figure 1b,
 45 decreasing steadily during 400 ns. However, with a bunch crossing every 25 ns, this long tail is
 46 a problem. It creates what is called of out-of-time pile-up, meaning the calorimeters still record
 47 energy deposited from previous bunch-crossing in addition to other pile-up event happening
 48 at the same bunch crossing. One has to imagine the electrical signal as the superposition of
 49 such triangles distributed randomly depending on if a particle deposited energy at a specific
 50 bunch crossing or not. With tails lasting almost 20 bunch crossings, a lot of out-of-time pile-up
 51 is created. To reduce this, the signal goes through a shaper in order to be more peaked right
 52 after the energy deposit and dive to negative values subsequently. It is then grouped in the
 53 Layer Sum Board (LSB) to create analogical Super-Cell signal. Then, it enters the LAr Trigger
 54 Digitizer Board (LTDB) where it is digitized at a rate of 40 MHz using a 12 bits analog-to-
 55 digital converter (ADC). As a result, at every bunch crossing, there is one ADC count of the
 56 original signal among the $2^{12} = 4096$ available values. The stream of ADC counts then enters
 57 the LAr Digital Processing System which contains LATOME boards grouped by four on LAr

58 carriers. The LATOME board is mainly built around an FPGA where a firmware block called
 59 `user_code` is responsible for the main physics task, most importantly computing the energy
 60 deposited in the Super-Cells.

61 The signal after going through the shaper is characterised by three important parameters:
 62 its pedestal, its amplitude and its time delay. The pedestal corresponds to an overall offset of
 63 the signal allowing for the negative part in the tail used to cancel out-of-time pile-up. The am-
 64 plitude of the signal – minus the pedestal – is proportional to the amount of energy deposited.
 65 Typically, the pedestal is around 1000 ADC counts, leaving around 3000 counts for the energy
 66 resolution. The time delay corresponds to the time between the recorded peak and the correct
 67 associated bunch crossing. This delay arises mainly from the shaping as a systematic effect of
 68 a few bunch crossings. For better precision, the phase is defined as the instant when the peak
 69 is found within the 25 ns of a bunch crossing.

70 The pedestal is extracted with calibration runs. The LATOME board is thus responsible for
 71 the computation of the energy and phase with the stream of ADC count for each Super-Cell it
 72 receives.

73 3 Energy computation in the LATOME board

74 The energy computation in the `user_code` of the LATOME is achieved thanks to the Optimal
 75 Filtering Coefficients (OFC) method [4] described in equations (1) and (2):

$$E_T(m) = \sum_{i=0}^{N-1} a_i \cdot (\text{ADC}_{m+i} - \text{ped}_{m+i}), \quad (1)$$

$$\xi(m) = \tau(m) \cdot E_T(m) = \sum_{i=0}^{N-1} b_i \cdot (\text{ADC}_{m+i} - \text{ped}_{m+i}). \quad (2)$$

76 This method uses a series of N signal value – in this case $N = 4$ – which are the signal
 77 in ADC count, labelled ADC_{m+i} , minus the pedestal, labelled ped_{m+i} . The OFC coefficients
 78 a_i (respectively b_i) are such that formula (1) (respectively (2)) gives on average the energy
 79 E_T (respectively $\xi(m)$ defined as the phase $\tau(m)$ for bunch crossing m multiplied with the
 80 energy) of the corresponding peak, with minimum variance associated.

81 Both the phase and the energy are computed at every bunch crossing. Then, a selection
 82 block in the `user_code` of the LATOME will check the phase is short enough by performing
 83 the following checks:

$$\begin{cases} -8E_T(m) < \xi(m) < 16E_T(m) & \text{for } E_T(m) > 10 \text{ GeV,} \\ -8E_T(m) < \xi(m) < 8E_T(m) & \text{for } 0 \text{ GeV} < E_T(m) \leq 10 \text{ GeV,} \\ 8E_T(m) < \xi(m) < -8E_T(m) & \text{for } -1 \text{ GeV} < E_T(m) \leq 0 \text{ GeV.} \end{cases} \quad (3)$$

84 If it is so, the energy is selected and automatically associated to its correct bunch crossing.
 85 However, a too long phase for example will indicate that the peak is best described by compu-
 86 tations at a previous bunch-crossing.

87 Additionally a saturation detection block detects and removes saturated pulses thanks to
 88 criteria on E_T and ξ , and a baseline correction blocks improves the out-of-time pile-up removal.
 89 Parameters needed for all these steps are the pedestal, the four optimal filtering coefficients
 90 a_i , again four b_i and six additional saturation detection criteria not discussed here.

91 These parameters are extracted through calibration runs. This has two consequences.
 92 Firstly, they are all Super-Cell specific. That is in total 16 parameters (the pedestal is counted
 93 twice as it appears in the two formulae) that are stored in registers of the FPGA for each Super-
 94 Cell. With 34 048 Super-Cells in total and typically around 300 Super-Cell managed by FPGA,

105 thousands of such Super-Cell specific parameters have to be stored in the correct register of
 106 the correct FPGA. Secondly these parameters are condition specific, the calibrations happening
 107 regularly. As a result, each set of parameters is stored with an interval of validity depending
 108 on run number and luminosity block of the LHC. This is achieved with the COOL condition
 109 database [5].

110 The output of the calibration runs are floats parameters stored in this format in the COOL
 111 database. However, in the firmware, the parameters are stored as integers of a specified num-
 112 ber of bits, e.g. the OFC coefficients are stored as 14 bits signed integers. Thus, a careful
 113 conversion taking into account signs and the upper limit in bits has to take place before filling
 114 the registers. Then, the outputs of the LATOME are also integers that have to be converted
 115 back to floats to be read. One important consequence is that, from floats to integers, informa-
 116 tion is lost and the result of the energy computation is degraded compared to the results with
 117 full float precision.

108 4 Commissioning

109 For commissioning, the structure detailed above was slightly modified to test separately every
 110 step. The generation of the condition database from calibration run being still under devel-
 111 opment, a dummy COOL database with smartly chosen coefficients was developed. As a first
 112 test, one of these dummy databases with all parameters set to dummy values, identical for all
 113 Super-Cells, was loaded successfully on LATOMEs.

114 The next step was to verify that the OFCs were correctly filled in the FPGA. First, the OFCs
 115 filled in the dummy database were chosen to be $a_i = b_i = (1, 0, 0, 0)$. The bunch crossings
 116 all have an index called BCID. For this test, the ADC counts being sent were exactly the BCID
 117 at which they were sent. To test LATOME boards, different output configuration are devised
 118 under the name *monitoring recipe*. Here, with the chosen monitoring recipe, the energy com-
 119 puted at one specific bunch crossing was read from the LATOME. The ADC counts used in
 120 the formula (1) were thus the BCID and the three subsequent ones. With this simple OFC
 121 configuration and all pedestals set to zero, the energy had to be the BCID at which the energy
 122 was computed. This configuration allowed to check specifically the first OFC, but other similar
 123 configurations were tried. One such was $a_i = b_i = (1, 2, 3, 4)$, validating all OFCs at once.

124 The general energy computation formula being validated, the next step was to check with
 125 real OFCs coming from calibration. They would thus all be different for each Super-Cell.
 126 Comparing the LATOME output to an offline computation allowed to check that the coeffi-
 127 cients were filled in the correct register in the correct format. As a result, in addition to the
 128 correct matching of OFCs with their Super-Cell, the conversion from floats to 14 bits signed
 129 integers was also verified in this test. In the monitoring recipe used, eleven ADC counts and
 130 the corresponding computed energy was read from the LATOME. The test run was sending to
 131 all channels of the LATOME ADC counts corresponding to a pulse. The formula (1) applies at
 132 every bunch crossing if at least four have been already received, so an energy output can be
 133 recomputed offline for the eight earliest ADC counts of an event, for each channel, and then
 134 compared to the LATOME energy output. This test was conducted with two LATOMEs, for a
 135 total of 580 Super-Cells, and for 1000 events. This amounts to $8 \times 580 \times 1000 = 4\,640\,000$
 136 energies recomputed in different conditions. The comparison of these to the corresponding
 137 LATOME output show strictly no difference. This validates the energy computation part of the
 138 `user_code` of the LATOME board.

139 However, the use of only half of the OFC coefficients can be verified by this method. The
 140 computation equivalent to formula (2) needs to be tested. This is first achieved by using the
 141 previous method and by replacing the OFCa by OFCb and comparing the offline computation

142 of ξ to the result of the LATOME output, showing again no difference. These tests also allowed
143 to check that the number of bits in the registers cover the a_i and b_i range. In this case, it had
144 to be adapted accordingly. A more complete test was performed by activating the selection
145 block in the user_code and implementing it also in the offline computation. On failing the
146 selection tests described in formula (3), the LATOME finally outputs an energy of 0 GeV. The
147 comparison of LATOME output and recomputed energy show again no difference as previously.

148 A similar offline computation can be done using this time the full float precision of the
149 OFC coefficients. The direct comparison of the full precision energy to the LATOME output
150 gives the energy resolution. Representing the difference between both results, a gaussian
151 distribution is obtained, yielding an energy resolution of around 50 MeV. This is approximately
152 the precision of the third least significant bit of the LATOME output – the least significant bit
153 representing 12.5 MeV. This precision choice is thus very conservative. Still, this resolution is
154 small compared to the typical values of the energy computed. The relative difference between
155 both energies is found to be very small, of the per mil level.

156 5 Conclusion

157 The computation of energy and time delay of LAr pulses is made possible by the Optimal
158 Filtering Coefficients method. The method is implemented in the firmware of the LATOME
159 board responsible for the computation of energy from Super-Cells in the new trigger chain
160 of the LAr calorimeter. As presented here, the commissioning of the computation of energy
161 and selection block is a success. Additionally, the energy output of the LATOME has a good
162 resolution, of the order of 50 MeV, corresponding to an accuracy at the per mil level. Yet, this
163 was only validated on a few LATOME and the full commissioning will have to happen on all
164 LATOME boards to check for possible outliers.

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