Measurement of *CP* structure of Higgs-tau Yukawa coupling

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Abstract

The CMS experiment at LHC has performed the first measurement of the CP structure of the Yukawa coupling between the Higgs boson and tau leptons. The measurement is based on data collected in proton-proton collisions at $\sqrt{s} = 13$ TeV during 2016-18, corresponding to an integrated luminosity of 137 fb^{-1}. The analysis utilizes the angular correlation between the decay planes of tau leptons produced in Higgs boson decays, where dedicated analysis techniques are used to optimise the reconstruction of tau decay planes. The measured value of CP mixing angle is $4 \pm 17°$, at 68% confidence level. The pure *CP*-odd hypothesis is excluded by 3.2 standard deviations. The analysis strategies and the results of the measurement are presented.

Contents

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1 Introduction

After the discovery of a Higgs boson of mass about 125 GeV at the Large Hadron Collider(**LHC**) [1–3], many studies are being performed to ensure whether the observed particle is the standard model(SM) Higgs boson. In the standard model, the Higgs boson's coupling to fermionic and bosonic fields preserves *CP* symmetry, often referred to as *CP*-even. The **ATLAS** [4] and **CMS** [5] collaborations have already probed *CP*-violating (*CP*-odd) interaction of the Higgs boson to gauge bosons. However, the *CP* odd state can couple to gauge bosons only at NLO or higher order, while its coupling to fermions can be probed at tree level. Although these studies have excluded that the Higgs boson is a pure *CP*-odd state (pseudoscalar), a *CP* mixture state is not fully excluded.

This analysis aims to access the potential mixing between a scalar and a pseudoscalar (*CP*-odd state) in the Yukawa coupling to the τ leptons via the angle between tau decay planes, the analysis results discussed in the following are based on Ref [6]. The interaction lagrangian of Higgs boson h of arbitrary *CP* nature to τ leptons is described as [7],

$$\mathcal{L}_Y = -\frac{m_\tau}{\nu}(\kappa_\tau \bar{\tau}\tau + \bar{\kappa}_\tau \bar{\tau} i\gamma_5 \tau)h \tag{1}$$

where, m_τ is the mass of the τ lepton, and the vacuum expectation value of Higgs field ν has a value of 246 GeV. The *CP*-even and *CP*-odd Yukawa couplings κ_τ and $\bar{\kappa}_\tau$ can be expressed in terms of effective mixing angle $\phi_{\tau\tau}$ as,

$$\tan\phi_{\tau\tau} = \frac{\bar{\kappa}_\tau}{\kappa_\tau} \begin{cases} \phi_{\tau\tau} \to 0, & CP\text{-even} \\ \phi_{\tau\tau} \to \frac{\pi}{2}, & CP\text{-odd} \\ \text{else}, & CP\text{-mix} \end{cases} \tag{2}$$

We define ϕ_{CP} as the angle between the τ decay planes at Higgs rest frame. This analysis measures the mixing angle ($\phi_{\tau\tau}$) from the relationship between $\phi_{\tau\tau}$ and ϕ_{CP} in the differential cross-section [8].

$$\frac{d\Gamma}{d\phi_{CP}} \propto -\cos(\phi_{CP} - 2\phi_{\tau\tau}) \tag{3}$$

The direct access of mixing angle from the ϕ_{CP} makes this analysis model-independent. The analysis is performed using full LHC Run-2 data, recorded by the CMS detector [9], corresponding to the integrated luminosity of 137 fb^{-1}, in the final states $\tau_\mu\tau_h$ and $\tau_h\tau_h$.

2 ϕ_{CP} Reconstruction

Tau lepton, the heaviest among the leptons, has short lifetime, and hence, decays to other lighter leptons or hadrons along with associated neutrinos. The momentum of the τ-lepton is reconstructed from its decay products. However, due to the presence of neutrinos in the final state the full momentum of the tau lepton cannot be reconstructed. Therefore the decay plane is constructed from its visible decay products. The methods that are used for each decay modes are described below [8],

- **Impact parameter Method** is used for the 1-prong decays such as ($\mu^\pm, \pi^\pm$), where tau decay plane is constructed from the momentum of the charged pion or hadron and its impact parameter vector.

- **Neutral-pion Method** is used when tau decay products contain at least one π^0 particle. Decay planes are constructed from the momenta of charged and neutral pions. In the case of 3-prong$\left(a_1^{3pr} \rightarrow \pi^\pm\pi^\mp\pi^\pm\right)$ decay of tau lepton $\pi^\pm$ meson that is oppositely charged to the a_1^{3pr} is considered as neutral pion vector for the purpose of constructing the decay plane.

- **Mixed Method** is used when one tau decays to one charged pion or hadron without π^0 and the other hand tau decays to charged prong along with a neutral pion. In this case the impact parameter method is used for the former and neutral pion method is used for the latter, respectively.

In all these methods the decay plane is constructed in the $\pi^+\pi^-$ zero momentum frame.

3 Analysis Strategy

We followed the same event selection strategy as used in the standard model Higgs to $\tau^+\tau^-$ analysis [10] for the $\tau_\mu\tau_h$ and $\tau_h\tau_h$ final state. However, we implemented some special methods to enhance the performance of this analysis.

3.1 Special Methods

Vertex Refitting: We exclude the tracks originating from the tau decay from the vertex fitting and apply beam spot constraints to improve primary vertex resolution, which improves the impact parameter measurement [11].
MVA decay mode identification: The analysis performance is improved by utilizing a multivariate based tau decay mode identification instead of the default HPS decay modes [12]. It enhances the assignment of 1 prong $+2\pi^0(a_1^{1pr})$ decay mode. This provides a 20% improvement in the expected sensitivity.

3.2 Background Estimation

The main background processes to consider are: Drell-Yan (Z/γ^*), $W+$jets, $t\bar{t}$, QCD multijet, electroweak W/Z, single-top and di-boson productions. All high fraction of backgrounds are estimated using data driven methods. The processes with genuine τ-leptons such as $Z/\gamma^* \rightarrow \tau\tau$ and small fraction of $t\bar{t}$ and di-boson are obtained from Embedded samples. [13]. Another major background is jets misidentified as taus ($j \rightarrow \tau_h$), which is estimated using fake factor method. [14]. The rest of the backgrounds processes like $Z/\gamma^* \rightarrow l^+l^-$ are obtained from the MC simulation.

3.3 Signal Extraction

Using multi-classification machine learning algorithm (Neural Network for $\tau_\mu\tau_h$ and BDT for $\tau_h\tau_h$) events are classified into three categories.

- **Higgs**: all signal processes (qqH,ggH and VH) combined into this category.

- **Embedded**: background processes involving two genuine τ-leptons.

- **Jet-Misidentification**: background process involving at least one misidentified jet $\rightarrow \tau$-lepton fake.

The 2D unrolled ϕ_{CP} distribution in the windows of increasing order of MVA score is used as the final discriminant. Due to the nature of the ϕ_{CP} distribution we can exploit symmetries in the background process to reduce statistical fluctuations in MC. In the final states where impact parameter method is used to reconstruct decay plane for both the tau leptons (e.g. $\mu\pi$,$\pi\pi$), the distributions of all the backgrounds are symmetrised around the central value. In other final states the background distributions are flattened. However, the jet $\rightarrow \tau$-lepton fake background distribution is symmetrised in all final states.

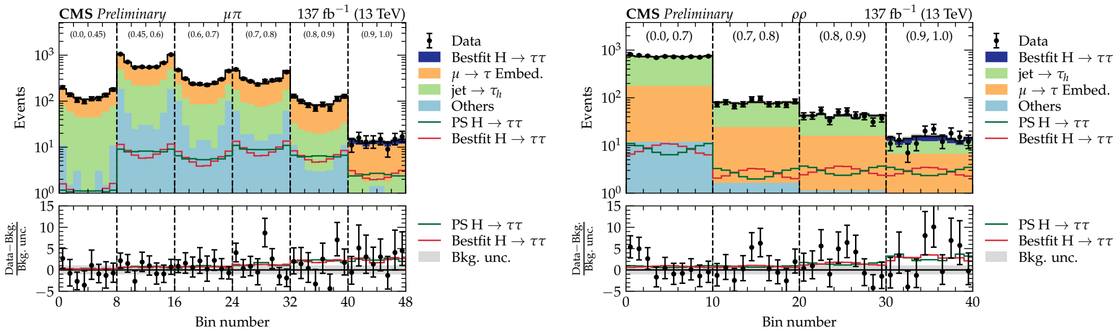

Figure 1: The unrolled ϕ_{CP} distributions for $\mu\pi$ and $\rho\rho$ are shown. The x-axis correspond to cyclic bins in ϕ_{CP} in the range of $(0, 2\pi)$. The $\mu\pi$(left) all backgrounds are symmetrised and for $\rho\rho$(right) backgrounds are flattened except jet $\rightarrow \tau$-lepton fake background [6]

4 Estimation of $\phi_{\tau\tau}$

The estimation of the $\phi_{\tau\tau}$ is obtained by the maximum likelihood fit using enrolled ϕ_{CP} distribution. The likelihood function $L(\vec{\mu}, \mu^{\tau\tau}, \vec{\theta})$ depends on the SM Higgs boson production signal strength ($\vec{\mu} = \mu_{ggH}, \mu_{qqH}, \mu_{VH}$), the $H \rightarrow \tau\tau$ decay branching fraction, *CP*-mixing angle, and the nuisance parameter ($\vec{\theta}$) accounted for the systematic uncertainties. The negative log-likelihood scan for the combination (NLL) of the $\tau_\mu\tau_h$ and $\tau_h\tau_h$ channel shown in Figure 2,where the negative likelihood is defined as:

$$-2\Delta \ln L = -2\big(\ln(L\phi_{\tau\tau}) - \ln(L\phi_{\tau\tau}^{\text{best fit}})\big)$$

We find the 68.3, 95.5, and 99.7% confidence intervals when $-2\Delta \ln L = 1.00, 4.02$ and 8.81 respectively [15]. The fit favours a scalar over the pseudoscalar $H\tau\tau$ coupling hypothesis at an observed(expected) sensitivity of 3.2(2.3) standard deviations. The measured value of the $\phi_{\tau\tau}$ with the decomposed uncertainty [6] is

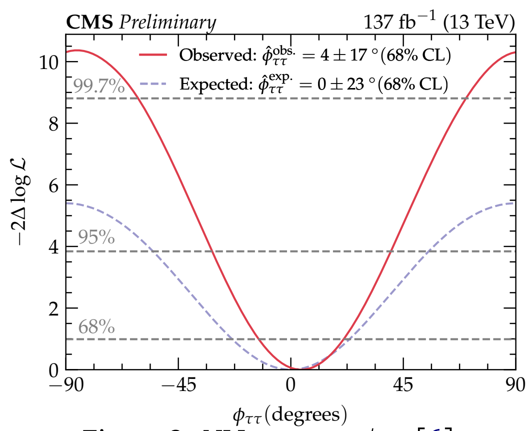

Figure 2: NLL scan on $\phi_{\tau\tau}$ [6]

$$\phi_{\tau\tau} = (4 \pm 17(stat) \pm 2(\text{bin-by-bin}) \pm 1(\text{syst}) \pm 1(\text{theory}))^\circ$$

Furthermore, we performed 2D fit of the branching fraction modifier concerning the SM value $\mu^{\tau\tau}$ versus $\phi_{\tau\tau}$, where we observe that there is no strong correlation. Also, the 2D scan for scalar and pseudoscalar Yukawa coupling fit shows that the best fit value is closer to the SM prediction.

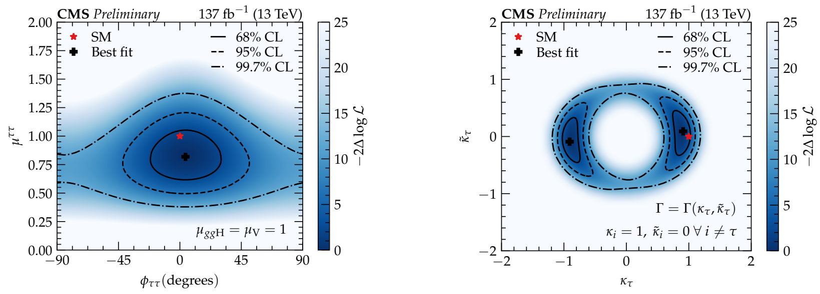

Figure 3: 2D scan of the branching fraction modifier with respect to the SM value of $\mu^{\tau\tau}$ versus $\phi_{\tau\tau}$(left). the 2D scan for scalar(κ) and pseudoscalar($\tilde{\kappa}$) τ Yukawa coupling (right) [6]

5 Conclusion

A measurement is performed of the CP mixing angle $\phi_{\tau\tau}$ in the Higgs to $\tau\tau$ coupling using 137 fb^{-1} of data recorded by the **CMS** experiment at centre-of-mass energy of 13 TeV. The best fit value of $\phi_{\tau\tau}$ is found to be $4 \pm 17°$. The analysis excludes a pure *CP*-odd scalar at a significance of 3.2 standard deviations. The results are consistent with the Standard Model prediction.

References

[1] S. Chatrchyan *et al.*, *Observation of a New Boson at a Mass of 125 GeV with the CMS Experiment at the LHC*, Phys. Lett. B **716**, 30 (2012), doi:10.1016/j.physletb.2012.08.021, 1207.7235.

[2] G. Aad *et al.*, *Observation of a new particle in the search for the Standard Model Higgs boson with the ATLAS detector at the LHC*, Phys. Lett. B **716**, 1 (2012), doi:10.1016/j.physletb.2012.08.020, 1207.7214.

[3] S. Chatrchyan *et al.*, *Observation of a New Boson with Mass Near 125 GeV in pp Collisions at $\sqrt{s} = 7$ and 8 TeV*, JHEP **06**, 081 (2013), doi:10.1007/JHEP06(2013)081, 1303.4571.

[4] M. Aaboud *et al.*, *Measurement of the Higgs boson coupling properties in the $H \to ZZ^* \to 4\ell$ decay channel at $\sqrt{s}$ = 13 TeV with the ATLAS detector*, JHEP **03**, 095 (2018), doi:10.1007/JHEP03(2018)095, 1712.02304.

[5] S. Chatrchyan *et al.*, *Study of the Mass and Spin-Parity of the Higgs Boson Candidate Via Its Decays to Z Boson Pairs*, Phys. Rev. Lett. **110**(8), 081803 (2013), doi:10.1103/PhysRevLett.110.081803, 1212.6639.

[6] *Analysis of the CP structure of the Yukawa coupling between the Higgs boson and τ leptons in proton-proton collisions at $\sqrt{s}$ =* 13 TeV, Tech. rep., CMS-PAS-HIG-20-006, Geneva, http://cds.cern.ch/record/2725571 (2020).

[7] A. V. Gritsan, R. Röntsch, M. Schulze and M. Xiao, *Constraining anomalous Higgs boson couplings to the heavy flavor fermions using matrix element techniques*, Phys. Rev. D **94**(5), 055023 (2016), doi:10.1103/PhysRevD.94.055023, 1606.03107.

[8] S. Berge, W. Bernreuther and S. Kirchner, *Determination of the Higgs CP-mixing angle in the tau decay channels at the LHC including the Drell–Yan background*, Eur. Phys. J. C **74**(11), 3164 (2014), doi:10.1140/epjc/s10052-014-3164-0, 1408.0798.

[9] S. Chatrchyan *et al.*, *The CMS Experiment at the CERN LHC*, JINST **3**, S08004 (2008), doi:10.1088/1748-0221/3/08/S08004.

[10] *Measurement of Higgs boson production and decay to the $\tau\tau$ final state*, Tech. rep., CMS-PAS-HIG-18-032, Geneva, https://cds.cern.ch/record/2668685 (2019).

[11] A. Cardini, *Tau identification exploiting deep learning techniques*, PoS **ICHEP2020**, 723 (2021), doi:10.22323/1.390.0723.

[12] *Identification of hadronic tau decay channels using multivariate analysis (MVA decay mode)* (2020), https://cds.cern.ch/record/2727092.

[13] *An embedding technique to determine backgrounds in proton-proton collision data*, Journal of Instrumentation **14**(06), P06032 (2019), doi:10.1088/1748-0221/14/06/p06032.

[14] *Measurement of the* $Z\gamma^* \rightarrow \tau\tau$ *cross section in pp collisions at* $\sqrt{s} = 13$ *TeV and validation of* τ *lepton analysis techniques*, Eur. Phys. J. C **78**(9), 708 (2018), doi:10.1140/epjc/s10052-018-6146-9, 1801.03535.

[15] *Precise determination of the mass of the Higgs boson and tests of compatibility of its couplings with the standard model predictions using proton collisions at 7 and 8 TeV*, The European Physical Journal C **75**(5) (2015), doi:10.1140/epjc/s10052-015-3351-7.