Studying hadronization at LHCb

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August 15, 2021

*Proceedings for the XXVIII International Workshop
on Deep-Inelastic Scattering and Related Subjects,
Stony Brook University, New York, USA, 12-16 April 2021*

Abstract

LHCb is a fully instrumented forward spectrometer with particle identification and muon reconstruction covering the pseudorapidity (η) range from 2 to 5. Its full jet reconstruction capability makes the LHCb experiment a suitable venue to explore jet substructure observables, particularly formation of hadrons and heavy quarkonia resonances inside jets. This contribution presents a brief overview of ongoing research program and discusses recent results on non-identified charged hadron distributions in Z-tagged jets and charmonium distributions within jets. Future works towards furthering the knowledge of hadronizion are also discussed.

1   Introduction

Three dimensional imaging of the nucleon structure has gained enormous interests over the past decade in anticipation of forthcoming data at the Electron-Ion Collider. As introducing the transverse momentum dependent (TMD) motion of partons inside the nucleon and their interactions with fundamental properties of the nucleon explained one of greatest mysteries, i.e. large transverse single spin asymmetry (TSSA) measured in pp collisions, adding an additional dimension to the description of formation of hadrons is only natural and will enrich the hadronic physics. On the other hand, with the success of Soft-Collinear Effective Theory (SCET) that allows solving multi-scale problems in a re-factorized manner, various observables probing substructure of hadro-produced jets have become calculable retaining predictive power of pQCD. Measurements presented here and other planned measurements involve looking inside jets for this reason.

2   Hadronization of charged particles in jets

In standard collinear perturbative QCD (pQCD) factorization framework, hadronization of charged particles is described by collinear fragmentation functions (FFs), denoted as $D_c^h(z, \mu)$, where z is the momentum fraction of an outgoing parton c carried by a hadron h (see [1] for a review). In this framework, single inclusive hadron production in pp collisions factorizes into hard scattering processes at short distance and dynamics at long distance described by collinear FFs and additionally collinear parton distribution functions (PDFs), denoted by

$f_{a/p}(x,\mu)$, that parameterize the momentum fraction x of an incoming proton carried by a parton a taking a part in hard scattering processes. The FFs and PDFs are non-perturbative (NP) objects that can only be measured in experiments, in contrast to the single inclusive jet production that is perturbatively calculable in pQCD. Jet fragmentation functions (JFFs) are experimental observables that measure the longitudinal momentum fraction of a jet carried by hadrons inside the jets. Within the SCET framework, JFFs are defined in such a way that they can probe collinear FFs. Analogously, TMD JFFs can access standard definition of TMD FFs measured in e+e- and SIDIS [2]. Theoretical development of this kind is crucial in the following sense. TMD factorization breaking is predicted for jet production in p+p [3], in which case non-perturbative TMD distribution and fragmentation functions can no longer be identified as universal objects in different collision processes. As such, interpretation of these measurements has to be carried out with caution and test of factorization breaking and universality of non-perturbative FFs across different collision systems is indispensible. LHCb data provide constraints on these NP physics objects in kinematic regimes complementary to those accessible at central pseudorapidities, i.e. a region of x smaller than 10^{-4} for TMD PDFs. The Drell Yan measurements with the $Z \to \mu\mu$ channel [4–6] at LHCb have been used in recent global analysis [7] of TMDs. Boer-Mulder's TMD PDF that affects the DY angular distributions are also expected to be an independent input to constraining TMD PDFs and test their universality.

The unpolarized TMD FFs depend on z, j_T and theory scale μ, i.e. $\hat{D}_i^h(z, j_T, \mu)$, therefore our main observables naturally involve z and j_T, and additionally r, experimentally defined in Eq. 1 and also found in Ref. [8].

$$z = \frac{p_h \cdot p_{jet}}{p_{jet}^2}, \qquad j_T = \frac{|p_h \times p_{jet}|}{p_{jet}}[GeV/c], \qquad r = \sqrt{d\phi_{h-jet}^2 + d\eta_{h-jet}^2} \qquad (1)$$

where z is the longitudinal momentum fraction of jet momentum carried by a hadron, j_T is the transverse component of hadron momentum with respect to the jet axis and r is the radial distance of a hadron from the jet axis.

The dominant hard partonic process for Z+jet production in the LHCb acceptance is $qg \to Zq$. Z-tagged jets at LHCb are therefore predominantly initiated by quarks and enhance sensitivity to the quark TMD FFs. The asymmetric parton momentum fractions of the incoming protons required for both Z and jet in the forward region and the back-to-back topo further enriches the data with jets initiated by light valence quarks with a large x .

The jet fragmentation function measurements of which results are being presented here are listed in Eq. 2 [9], where the phase space $\mathcal{PS} = d\eta_Z d\eta_{jet} dp_T d^2\mathbf{q_T}$, where $\mathbf{q_T}$ is the vector sum of the transverse momenta of the Z boson and the jet, and $2\mathbf{p_T}$ the difference between the two.

$$F(z_h) = \frac{d\sigma}{d\mathcal{PS}\,dz_h}\Big/\frac{d\sigma}{d\mathcal{PS}}, \qquad \int dz_h\, f(z_h, j_T) = \frac{d\sigma}{d\mathcal{PS}\,dj_T}\Big/\frac{d\sigma}{d\mathcal{PS}}, \qquad (2)$$

$$G(r) = \frac{d\sigma}{d\mathcal{PS}\,dr}\Big/\frac{d\sigma}{d\mathcal{PS}}, \qquad \text{where} \quad h : \pi^{\pm} \text{ or } K^{\pm} \text{ or } p^{\pm}.$$

Note that the second equation is equivalent to the TMD JFF, denoted by $f(z_h, j_T)$, integrated over z_h. Experimentally, these quantities can be expressed as shown in Eq. 3.

$$F = \frac{1}{N_{Z+jet}}\frac{dN_h}{dz_h}, \qquad \int dz_h f = \frac{1}{N_{Z+jet}}\frac{dN_h}{dj_T}, \qquad G = \frac{1}{N_{Z+jet}}\frac{dN_h}{dr} \qquad (3)$$

2.1 Results

Measured longitudinal momentum distributions z is shown in Fig. 1. In the low z region less than 0.04, a humped-back structure that is attributed to color-coherence effects compatible with parton-hadron duality [10] was is seen. More energy available to produce hadrons in jets from harder scattering explains the increase of the height of a hump with jet p_T. At mid-to-high z, between 0.04 and 0.4, the fragmentation of a jet into hadrons is governed by the scaling behavior. The highest z region is experimentally challenging due to limited statistics and degrading reconstruction efficiency. Theoretically, the presence of large logarithms that have to be re-summed with good precision adds to the challenge. Comparing the collinear JFF results of the LHCb Z+jet to the ATLAS single inclusive jet, an obvious difference seen in the shape of the distributions, an excess at higher z and a depletion at lower z, is as expected based on the gluon dominance in the single inclusive jet production at LHC. A consistency between the LHCb Z+jet and the ATLAS γ+jet results confirms this interpretation.
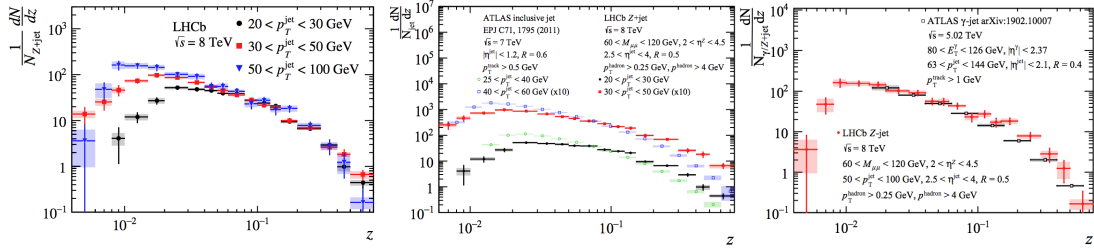
Figure 1: Measured z distributions at $\sqrt{s}$=8 TeV (left) and comparisons with ATLAS inclusive jets (middle) and ATLAS γ-jets (right).

Measured j_T and radial distributions are shown in Fig. 2. The shape of j_T distributions does not show strong dependence on the p_T of the jets, while harder jets tend to produce more particles in the collinear direction of the jets.
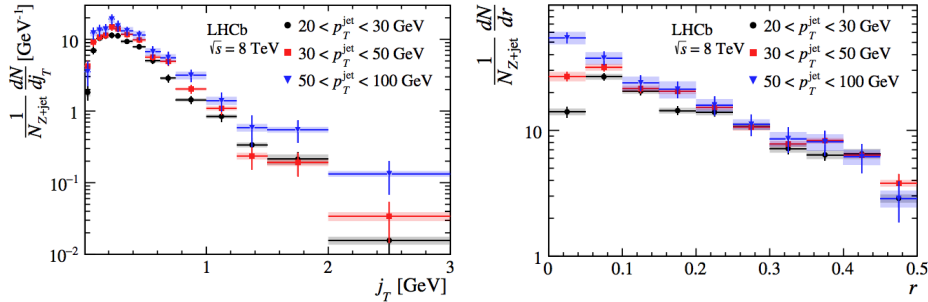
Figure 2: Measured j_T distributions (left) and r distributions (right) at $\sqrt{s}$=8 TeV.

The results are in good agreement with the TMD analysis performed in Ref. [9] in regions where perturbative calculations lack large logarithms that need to be resummed with high precision. This is promising for further studies to be published in the future.

2.2 Future works

Fully differential TMD JFF measurements, i.e. with respect to j_T and z_h are being studied at LHCb along with identified charged hadron distributions. Identified charged hadron distributions provide additional information on the valence versus sea quark's role initiating the parton shower that leads to the creation of jets. The corresponding observable is shown in

Eq. 4. When measured at a different beam energy, the observables can provide insights into the universality of the JFFs.

$$f(z_h, j_T) = \frac{d\sigma}{d\mathcal{PS}dz_h dj_T} \Big/ \frac{d\sigma}{d\mathcal{PS}}, \quad \text{where} \quad h = \text{identified } \pi^\pm, K^\pm, p^\pm. \tag{4}$$

Measurements that are sensitive to electromagnetic charge such as positive-to-negative charge ratio or jet charge [11] can offer a more systematic way of decomposing jets into different parton flavors. Lastly, similar properties in heavy quark initiated jets is also being investigated to compare with light-quark initiated jets.

3 Formation of quarkonia in jets

Production mechanisms of heavy quarkonia is not well understood. LHCb published a slew of results on charmonium/bottomonium cross sections and polarizations [?, 12–15]. Non-relativistic QCD (NRQCD), the most common approach to this problem, is generally successful in describing quarkonuia productions across different experiments. In this framework, the long distance matrix elements (LDMEs) that describe the nonperturbative transition of the $c\bar{c}$ pair in a state of definite color and angular momentum quantum number into a final state containing J/ψ are extracted from data. The challenge is that almost all of the world data on quarkonia polarization contradict the high transverse polarization predicted from NRQCD using extracted LDMEs. An alternative approach to understanding quarkonium production mechanism is to measure quarkonium in jets. The usual observable discussed in the previous section is to measure is the collinear JFF of quarkonia, i.e. longitudinal momentum fraction of the jets carried by quarkonia $z = p_T^{J/\psi}/p_T^{jet}$. For the first time, LHCb measured this for charmonium, prompt as well as from b-hadron decay [16].

3.1 Results

Measured z for prompt J/ψ and from b-hadron decay are shown in Fig. 3. The data are compared to the leading-order (LO) NRQCD-based prediction as implemented in PYTHIA8. This prediction used LDMEs determined empirically both for color-octet and color-singlet intermediate states. While b quark fragmentation into a jet with a b-hadron that subsequently decays into a charmonium state is fairly well described by the LO NRQCD predictions, the prompt J/ψ counterpart in data takes much smaller momentum fraction of jets within which they are created than the PYTHIA8 predictions.
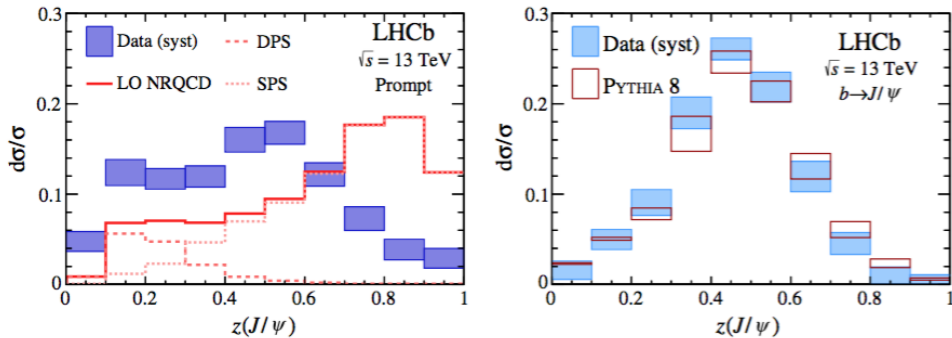
Figure 3: Measured z distributions for prompt J/ψ (left) and from b-decay (right) at $\sqrt{s}$=13 TeV.

To reconcile the discrepancy between the data and the Pythia8 predictions, Ref. [17] performed a careful analysis using NRQCD. It was noted that the parton shower of the heavy quark-antiquark pair produced in a color-octet state was treated as a color singlet particle splitting into quark-antiquark pair 2 $P_{qq}(z)$, of which probability peaks at z=1. At LHC, quarkonia production is predominantly by gluon fragmentation. In [17], two methods were introduced to address this point, namely gluon-fragmentation-improved-PYTHIA and fragmenting jet function (FJF). Given the multi-scale nature of the quarkonium production, the latter approach is an appropriate solution, extended from JFJ defined for hadron JFF measurements. Both methods capture the softer z seen in the LHCb prompt J/ψ data.

3.2   Future works

The z measurements can be repeated for bottomonium. The z-dependent polarization of charmonium as well as bottomonium in jets are being measured at LHCb. More differential measurements of this kind are expected to provide a better handle on understanding and constraining different production mechanisms of quarkonia. Separation between heavy-quark and gluon fragmentation is also an experimentally challenging question of interest.

4   Conclusion

LHCb have successfully performed measurements studying properties/mechanisms of hadronization using jet substructure technique, charmonium for the first time, followed by non-identified hadrons in jets. A number of follow-up measurements are being carried out towards gaining insights into the flavor dependence of the quark that initiates jets that produce charged hadrons and the full TMD picture of the hadronization process. Quarkonium production mechanisms are still not well understood. Nonetheless, a new jet substructure type of measurements opened up ways to test the current state-of-the-art theory and raised more questions to be answered. Other measurements being explored include transverse polarization of Λ baryons towards accessing TMD FF and Z-hadron correlations for testing TMD factorization breaking.

Acknowledgements

Funding information   The author is supported by the U.S. National Science Foundation.

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