

Status and prospects for τ property measurements at Belle II

Michel Hernández Villanueva^{1*}
on behalf of the Belle II collaboration

¹ Deutsches Elektronen-Synchrotron (DESY)
* michel.hernandez.villanueva@desy.de

December 10, 2021

*16th International Workshop on Tau Lepton Physics (TAU2021),
September 27 – October 1, 2021
doi:10.21468/SciPostPhysProc.?*

1 Abstract

2 The Belle II experiment is a major upgrade of the Belle detector, operating at the Su-
3 perKEKB energy-asymmetric e^+e^- collider with a design luminosity of $8 \times 10^{35} \text{ cm}^{-2}\text{s}^{-1}$.
4 To date, the detector has recorded over 200 fb^{-1} of data and aims to collect 50 ab^{-1} , a fac-
5 tor of x50 more than its predecessor. Belle II has a broad τ physics program, from high-
6 precision measurements of SM parameters to searches of new physics via experimental
7 probe of BSM processes. In this work, we review the status of the Belle II experiment,
8 and the prospects for the measurement of the τ lepton mass and lifetime, fundamental
9 inputs in tests of LFU violation.

10

11 Contents

12	1 Introduction	2
13	2 Performance studies with τ leptons	2
14	2.1 Tracking efficiency	2
15	2.2 Particle identification performance	3
16	3 Measurements of τ lepton properties	3
17	3.1 τ lepton mass	3
18	3.2 τ lepton lifetime	5
19	4 Conclusion	6
20	References	7

21
22

1 Introduction

The Belle II experiment is a detector installed on the asymmetric electron-positron collider SuperKEKB, located at the KEK laboratory in Tsukuba, Japan. Working at the center-of-mass energy of 10.58 GeV, SuperKEKB is designed to reach an instantaneous luminosity of $3.1 \times 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$. Belle II represents an upgrade of its predecessor Belle, with major modifications in each of the subsystems aiming to reconstruct collision products in a high-luminosity environment and collect 50 ab^{-1} by the end of the operation. A complete description of the Belle II detector can be found in Ref [1]. In 2021, SuperKEKB has set a new world record reaching a peak luminosity of $L_{\text{peak}} = 3.1 \times 10^{34} \text{ cm}^{-2} \text{ s}^{-1}$. To date, 213 fb^{-1} of collision data have been recorded by Belle II.

At the energy operation of Belle II, the cross-section of the process $e^+e^- \rightarrow \tau^-\tau^-$ (0.92 nb) is similar to the cross-section for the production of B-meson pairs (1.10 nb). The well-defined kinematics in the events with a low-background environment, added to the enormous amount of collisions recorded by Belle II, enables a world-leading τ physics program in the next decade [2], with 45 billions of τ lepton pairs expected in the data set recorded by the end of the operations.

One important step for precision measurements and searches of non-Standard Model physics is the measurement of the τ lepton properties with the highest precision, since the uncertainty in quantities such as the mass or the lifetime has important consequences in searches like lepton flavor universality violation induced by beyond Standard Model interactions.

2 Performance studies with τ leptons

The production of τ pair leptons in Belle II is not only useful for the understanding of the physics involved in the decays but also a powerful tool for performance studies. Unique features of τ lepton allow the study of the efficiency in the detection of tracks, as well as testing the charged-particle identification performances.

2.1 Tracking efficiency

The tracking efficiency is determined with a tag-and-probe method following Ref [3], using $e^-e^+ \rightarrow \tau^-\tau^+$ events where one of the τ leptons decay leptonically into either an electron or a muon, while the other τ lepton decays as $\tau^\pm \rightarrow \pi^\pm \pi^\pm \pi^\pm \nu_\tau + \geq 0 \pi^0$. Such τ decays are labeled as 1-prong and 3-prong respectively, depending on the number of the detectable charged particles in the channel. Three good quality tracks with total charge ± 1 are used to tag events, with the existence of a probe track ensured by charge conservation.

Precision analyses require to measure the tracking efficiency in data to assign a systematic uncertainty due to the mis-modeling in the Monte Carlo simulation (MC). The data-MC discrepancy is defined as

$$\delta = 1 - \frac{\epsilon_{\text{data}}}{\epsilon_{\text{MC}}}, \quad (1)$$

where ϵ_{data} and ϵ_{MC} are the tracking efficiencies measured in data and simulation. Figure 1 shows the overall discrepancy between data and MC, subdivided by channel and by run period, as well as the combined result, yielding

$$\delta_{\text{overall}}^* = 0.13 \pm 0.16(\text{stat}) \pm 0.89(\text{syst})\%. \quad (2)$$

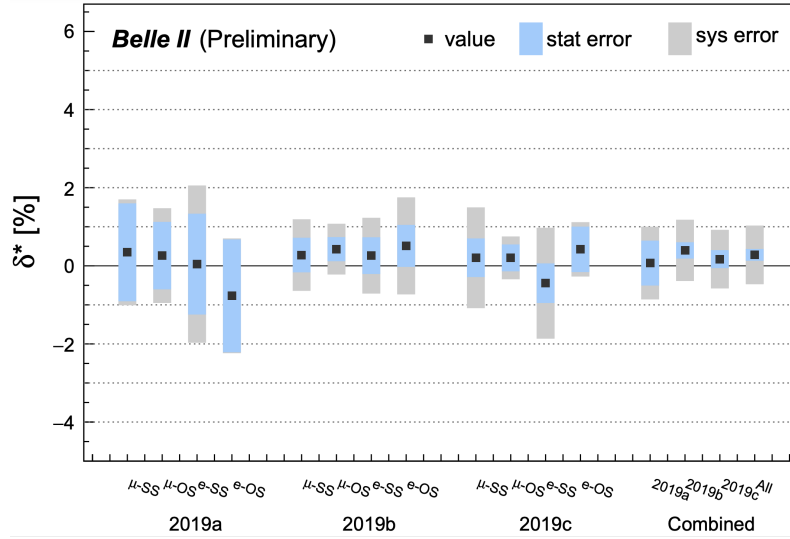


Figure 1: Overall data vs MC discrepancy for tracking efficiency, shown for individual τ pairs with one of them decaying into an electron or a muon, and with the two tagging tracks in the 3-prong side of same-sign (SS) or opposite-sign (OS), as well as for different data taking periods [4]. Statistical and systematic uncertainties are shown in grey and blue, respectively.

60 2.2 Particle identification performance

61 Using a similar approach to that of section 2.1, a tag-and-probe method is used with τ pair
 62 events, selecting pions from one of the τ leptons decaying to a 3-prong channel $\tau^\pm \rightarrow h^\pm h^\pm h^\pm \nu_\tau$.
 63 Tagging two same-charged tracks as pions provide a pure sample of pions in the remaining
 64 opposite-charged track, since the mode with kaons that would pass the tag selection, such as
 65 $\tau^- \rightarrow \pi^- K^+ \pi^- \nu_\tau$, are highly suppressed.

66 The probing track is used to calculate particle identification (ID) efficiency for pions and
 67 $\pi^\pm \rightarrow \ell^\pm$ mis-identification rates, in bins of momentum and polar angle. Figure 2 shows the
 68 example of a bin for electron ID, where the $\pi \rightarrow e$ mis-identification obtained from τ lepton
 69 decays is compared with other methods.

70 3 Measurements of τ lepton properties

71 3.1 τ lepton mass

72 Lepton masses are fundamental parameters of the Standard Model. While the mass of the
 73 electron and the muon are known with very high precision, the short lifetime of the τ lepton
 74 complicates the measurement of its mass. The two methods available for the determination of
 75 m_τ are the scan of the production threshold, used by the BESIII collaboration to provide the
 76 most precise measurement of m_τ [6], and the pseudomass method developed originally by the
 77 ARGUS collaboration [7], with recent updates by Belle and BaBar performing the measurement
 78 of m_τ at the center-of-mass energy of the $\Upsilon(4S)$ resonance [8, 9].

79 Belle II has performed a measurement of the τ lepton mass with the data recorded during
 80 2019 - 2020, corresponding to an integrated luminosity of 8.8 fb^{-1} . Using $e^+e^- \rightarrow \tau^+\tau^-$ events
 81 at the center-of-mass energy of 10.58 GeV, with one of the τ leptons decaying in a 3-prong
 82 mode $\tau^\pm \rightarrow \pi^\pm \pi^\pm \pi^\pm \nu_\tau$, and the other in a 1-prong channel $\tau^\pm \rightarrow \ell^\pm \bar{\nu}_\ell \nu_\tau$ or $\tau^\pm \rightarrow \pi^\pm \nu_\tau$.
 83 Since in Belle II the collisions happen far from the production threshold, the pseudomass

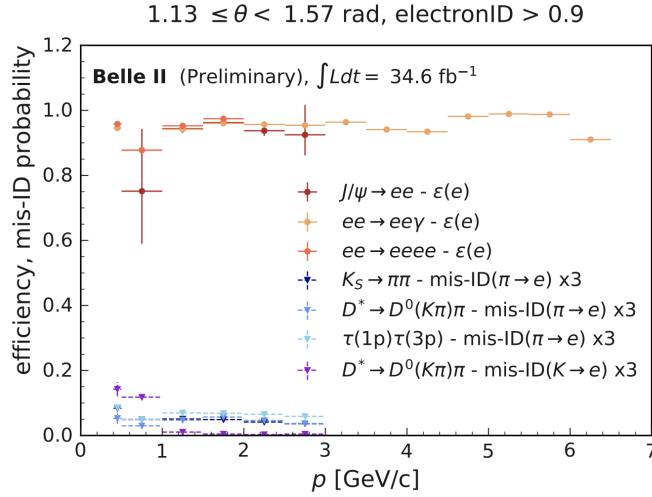


Figure 2: Example of a bin of the polar angle θ for electron ID with efficiencies and hadron-lepton mis-identification rates from all measurements [5]. The mis-identification rates have been inflated by a factor of 3 for illustration. $\pi \rightarrow e$ mis-identification obtained from τ lepton decays is shown in cyan.

84 technique is used to measure the mass of the τ lepton. To construct the pseudomass, the
85 kinematic properties of a semileptonic decay $\tau^\pm \rightarrow h^\pm + \nu_\tau$ are used,

$$m_\tau^2 = (p_h + p_\nu)^2 = 2E_h(E_\tau - E_h) + m_h^2 - 2|\vec{p}_h|(E_\tau - E_h) \cos(\vec{p}_h, \vec{p}_\nu), \quad (3)$$

86 where E_i , p_i and m_i are the energy, 4-momentum and invariant mass for the i -particle, h is the
87 hadronic system coming from the τ lepton, and $m_\nu = 0$ is assumed. As the neutrinos in the
88 event escape from detection and \vec{p}_ν is not known, the approximation $\cos(\vec{p}_\nu, \vec{p}_h) = 1$ is taken,
89 resulting in

$$M_{\min}^2 = 2E_h(E_\tau - E_h) + m_h^2 - 2|\vec{p}_h|(E_\tau - E_h) < m_\tau^2, \quad (4)$$

90 with M_{\min} being called the pseudomass. The position of the cutoff in M_{\min} estimates the value
91 of the mass, while the distribution is smeared by the resolution in the detectors and initial-state
92 radiation (ISR) / final-state radiation (FSR) effects.

93 Figure 3 shows the distribution of M_{\min} for the τ decaying to the 3-prong mode in both
94 data and MC, and a magnification into the region of interest $1.7 < M_{\min} < 1.85$ GeV/ c^2 where
95 the cutoff is observed. An empirical probability density function (p.d.f.) is built to determine
96 the mass of the τ lepton,

$$F(M_{\min}, \vec{P}) = (P_3 + P_4 \cdot M_{\min}) \cdot \tan^{-1}[(M_{\min} - P_1)/P_2] + P_5 M_{\min} + 1, \quad (5)$$

97 where P_1 is an estimator of m_τ . Figure 4 shows the result of a maximum likelihood fit per-
98 formed with (5), yielding after correction factors $m_\tau = 1777.28 \pm 0.75(\text{stat})$ MeV/ c^2 .

99 Main sources of systematic uncertainties are the momentum shift due to imperfections
100 on the magnetic field map (0.29 MeV/ c^2), the determination of the bias in the estimator P_1
101 (0.12 MeV/ c^2), the choice of the p.d.f. versus alternative functions (0.08 MeV/ c^2), the choice
102 of the fit window (0.04 MeV/ c^2) and the energy beam determination from the beam-energy-
103 constrained mass of fully reconstructed neutral and charged B decays (0.03 MeV/ c^2). Details
104 about the estimation for each source can be found in [10]. Therefore, the mass of the τ lepton
105 measured is

$$m_\tau = 1777.28 \pm 0.75(\text{stat.}) \pm 0.33(\text{syst.}) \text{ MeV}/c^2. \quad (6)$$

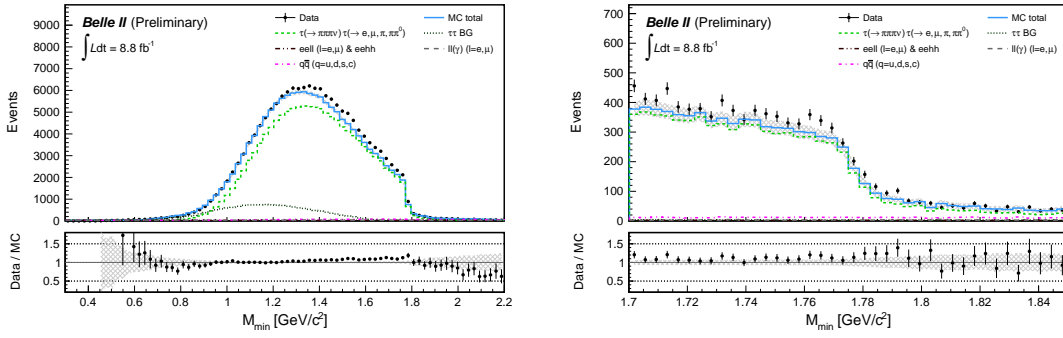


Figure 3: Distribution of the pseudomass for the $\tau^\pm \rightarrow \pi^\pm \pi^\pm \pi^\pm \nu_\tau$ mode in both data and MC, and the magnification into the region of interest 1.7 to 1.85 GeV/c^2 where the cutoff is observed [10].

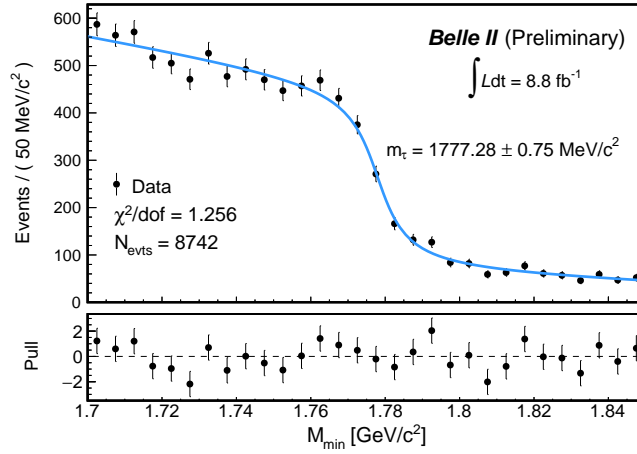


Figure 4: Distribution of the pseudomass for the $\tau^\pm \rightarrow \pi^\pm \pi^\pm \pi^\pm \nu_\tau$ mode in the region of interest 1.7 to 1.85 GeV/c^2 [10]. The blue solid line represents the p.d.f. described in the text, with parameters determined by a maximum likelihood fit.

106 The systematic uncertainty of m_τ in the Belle II measurement is compatible with the system-
 107 atic uncertainty reported by Belle [8]. The analysis is being updated to use the most recent
 108 data collected, where an improved understanding of the detector could reduce the current
 109 systematic uncertainty.

110 3.2 τ lepton lifetime

111 Since the τ lepton discovery, a large variety of methods have been applied to determine its
 112 lifetime. The most precise measurement has been provided by the Belle collaboration [11]
 113 with $\tau^+\tau^-$ events where both τ -leptons decay to three charged pions and a neutrino. The
 114 technique proposed at Belle II exploits the small beam spot size at the interaction region,
 115 obtained through the nano-beam scheme of SuperKEKB, where the vertical size of the beam
 116 is $\sigma_y < 1\mu\text{m}$ and the primary vertex position \vec{IP} on the transverse plane is treated as known.
 117 The advantage with respect to the method used by Belle is that only one of the leptons in the
 118 τ pair produced from the collision requires to decay to three charged particles, while the other
 119 can be reconstructed via 1-prong channels with highest branching ratios.

120 In the center-of-mass frame, the momentum \vec{p}_τ^{CM} of a τ lepton produced from a e^+e^- colli-

121 sion can be determined by the intersection of two cones when both leptons decay semilepton-
 122 ically, up to a two-fold ambiguity. Using the average of the two possible solutions for \vec{p}_τ^{CM} , the
 123 lifetime of the τ lepton is determined from the distribution of the proper time t , obtained from
 124 the reconstructed decay length ℓ_τ and the estimated \vec{p}_τ^{CM} boosted to the laboratory frame,

$$t = \ell_\tau \frac{m_\tau}{|\vec{p}_\tau|c}. \quad (7)$$

125 The decay length ℓ_τ is obtained from the relation

$$\vec{IP} + \ell_\tau \hat{n}_\tau - \vec{v} = 0, \quad (8)$$

126 where IP_y is fixed, and $\hat{n}_\tau = \vec{p}_\tau/|\vec{p}_\tau|$ is the direction of the τ lepton in the laboratory frame.
 127 The decay vertex position \vec{v} of a 3-prong τ decay is determined by the intersection of the tracks
 128 coming from the τ .

129 A sensitivity study for the determination of the τ lifetime has been performed in Belle II,
 130 with a simulated data set corresponding to 200 fb^{-1} at the $\Upsilon(4S)$ collision energy. τ pair
 131 events are selected with one decaying in a 3-prong mode $\tau^\pm \rightarrow \pi^\pm \pi^\pm \pi^\pm \nu_\tau$, and the other to
 132 a 1-prong mode $\tau^\pm \rightarrow (\rho^\pm \rightarrow \pi^\pm \pi^0) \nu_\tau$. Figure 5 shows the distribution of the reconstructed
 133 proper time in MC events. The extracted lifetime from the fit is $\tau_\tau = 287.2 \pm 0.5(\text{stat.}) \text{ fs}$.
 134 Using an integrated luminosity 3.5 lower, the statistical uncertainty is compatible with the
 135 most precise measurement performed by Belle.

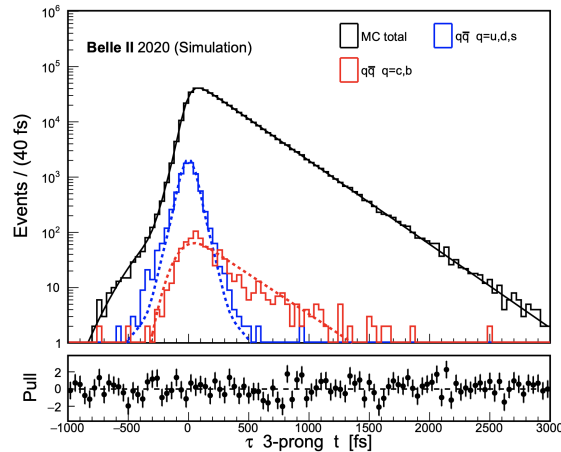


Figure 5: Fit on the distribution of the reconstructed proper time for events with a 3-prong τ lepton selection. Background contribution from $q\bar{q}$, $q = u, d, s$ events is subtracted with a fit to a linear combination of three Gaussians, while $q\bar{q}$, $q = c, b$ is fitted with a convolution of a Gaussian and an exponential function.

136 4 Conclusion

137 Precision measurements of fundamental τ lepton parameters are an important part of the
 138 world-leading Belle II physics program towards the searches of non-standard-model physics.
 139 τ lepton mass studies performed with the early data show a good performance of the detector,
 140 with perspectives to reduced systematic uncertainties in the update for the measurement of m_τ
 141 using the pseudomass technique. The lifetime measurements at Belle II show the potential of
 142 the nano-beam scheme with an upgraded vertex detection system. First studies are promising,
 143 with an update in the measurement feasible in the coming months.

144 **References**

- 145 [1] T. Abe *et al.*, *Belle II Technical Design Report* (2010), [1011.0352](#).
- 146 [2] W. Altmannshofer *et al.*, *The Belle II Physics Book*, PTEP **2019**(12), 123C01 (2019),
147 doi:[10.1093/ptep/ptz106](#), [Erratum: PTEP 2020, 029201 (2020)], [1808.10567](#).
- 148 [3] T. Allmendinger *et al.*, *Track Finding Efficiency in BaBar*, Nucl. Instrum. Meth. A **704**, 44
149 (2013), doi:[10.1016/j.nima.2012.11.184](#), [1207.2849](#).
- 150 [4] The Belle II Collaboration, *Measurement of the tracking efficiency and fake rate with*
151 *$e^+e^- \rightarrow \tau^+\tau^-$ events* (2020), BELLE2-NOTE-PL-2020-014.
- 152 [5] The Belle II collaboration, *Muon and electron identification efficiencies and hadron-lepton*
153 *mis-identification probabilities* (2020), BELLE2-NOTE-PL-2020-027.
- 154 [6] M. Ablikim *et al.*, *Precision measurement of the mass of the τ lepton*, Phys. Rev. D **90**(1),
155 012001 (2014), doi:[10.1103/PhysRevD.90.012001](#), [1405.1076](#).
- 156 [7] H. Albrecht *et al.*, *A Measurement of the tau mass*, Phys. Lett. B **292**, 221 (1992),
157 doi:[10.1016/0370-2693\(92\)90634-G](#).
- 158 [8] K. Abe *et al.*, *Measurement of the mass of the tau-lepton and an upper limit on*
159 *the mass difference between tau+ and tau-*, Phys. Rev. Lett. **99**, 011801 (2007),
160 doi:[10.1103/PhysRevLett.99.011801](#), [hep-ex/0608046](#).
- 161 [9] B. Aubert *et al.*, *Measurements of the tau mass and the mass difference of the tau+ and*
162 *tau- at BABAR*, Phys. Rev. D **80**, 092005 (2009), doi:[10.1103/PhysRevD.80.092005](#),
163 [0909.3562](#).
- 164 [10] F. Abudinén *et al.*, *τ lepton mass measurement at Belle II* (2020), [2008.04665](#).
- 165 [11] K. Belous *et al.*, *Measurement of the τ -lepton lifetime at Belle*, Phys. Rev. Lett. **112**(3),
166 031801 (2014), doi:[10.1103/PhysRevLett.112.031801](#), [1310.8503](#).