

# Bethe vectors and recurrence relations for twisted Yangian based models

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## Abstract

We study Olshanski twisted Yangian based models, known as one-dimensional “soliton non-preserving” open spin chains, by means of algebraic Bethe ansatz. The even case, when the bulk symmetry is  $\mathfrak{gl}_{2n}$  and the boundary symmetry is  $\mathfrak{sp}_{2n}$  or  $\mathfrak{so}_{2n}$ , was studied in [GMR19]. In the present work, we focus on the odd case, when the bulk symmetry is  $\mathfrak{gl}_{2n+1}$  and the boundary symmetry is  $\mathfrak{so}_{2n+1}$ . We explicitly construct Bethe vectors and present a more symmetric form of the trace formula. We use the composite model approach and  $Y(\mathfrak{gl}_n)$ -type recurrence relations to obtain recurrence relations for twisted Yangian based Bethe vectors, for both even and odd cases.

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## 47 1 Introduction

48 Twisted Yangian based models, known as one-dimensional “soliton non-preserving” open spin  
 49 chains, were first investigated by means of analytic Bethe ansatz techniques [Doi00, AA<sup>+</sup>05,  
 50 AC<sup>+</sup>06a, AC<sup>+</sup>06b] and more recently in [ADK15]. Such models are known to play a role in  
 51 Yang-Mills theories, where twisted Yangians emerge in the context of integrable boundary  
 52 overlaps [dL<sup>+</sup>19, Gom24] and open fishchains [GJP21].

53 A crucial step in understanding twisted Yangian based models is finding explicit expres-  
 54 sions of Bethe vectors. In the case when the bulk symmetry is  $\mathfrak{gl}_{2n}$  and the boundary symmetry  
 55 is  $\mathfrak{sp}_{2n}$  or  $\mathfrak{so}_{2n}$ , this was achieved in [GMR19] using algebraic Bethe ansatz techniques put for-  
 56 ward in [Rsh85, DVK87]. These techniques apply to the cases, when the  $R$ -matrix intertwining  
 57 monodromy matrices of the model can be written in a six-vertex block-form. The monodromy  
 58 matrix is then also written in a block-form, in terms of matrix operators  $A$ ,  $B$ ,  $C$ , and  $D$ , that are  
 59 matrix analogues of the conventional creation, annihilation and diagonal operators. Exchange  
 60 relations between these matrix operators turn out to be reminiscent of those of the standard  
 61 six-vertex model. Such techniques have been used to study  $\mathfrak{so}_{2n}$ - and  $\mathfrak{sp}_{2n}$ -symmetric spin  
 62 chains in [Rsh91, GP16, GR20a, GR20b, Reg22]. A more general framework of such techniques  
 63 has recently been proposed in [Ger24].

64 In the present paper we extend the results of [GMR19] to the odd case, when the bulk  
 65 symmetry is  $\mathfrak{gl}_{2n+1}$  and the boundary symmetry is  $\mathfrak{so}_{2n+1}$ . This extension is based on a simple  
 66 observation that the generating matrix of the odd twisted Yangian  $Y^+(\mathfrak{gl}_{2n+1})$  can be decom-  
 67 posed into four overlapping  $(n+1) \times (n+1)$ -dimensional matrix operators satisfying the same  
 68 exchange relations as those of  $Y^+(\mathfrak{gl}_{2n+2})$  thus allowing us to employ the same algebraic Bethe  
 69 ansatz approach. However, the overlapping introduces a new challenge since the middle entry  
 70 of the generating matrix is now included in both  $A$  and  $B$  matrix operators leading to an uncer-  
 71 tainty in the  $AB$  exchange relation. This issue is resolved in the technical Lemma 3.8 stating  
 72 action of the middle entry on Bethe vectors. Computing this action requires knowledge of  
 73 recurrence relations for Bethe vectors. We use the composite model techniques together with  
 74 the  $Y(\mathfrak{gl}_n)$ -type recurrence relations found in [HL<sup>+</sup>17b] to obtain the  $Y^\pm(\mathfrak{gl}_{2n})$ - and  $Y^+(\mathfrak{gl}_{2n+1})$ -  
 75 type recurrence relations. The main results of this paper are presented in Theorem 3.9 and  
 76 Propositions 4.4 and 4.6.

77 The first main result, Theorem 3.9, states that Bethe vectors, defined by formula (3.42),  
 78 are eigenvectors of the transfer matrix, defined by formula (3.44), provided Bethe equations  
 79 (3.53) and (3.54) hold. This Theorem is an extension of Theorems 4.3 and 4.4 in [GMR19]  
 80 to the odd case. Commutativity of transfer matrices is shown in Appendix A.2. We also found  
 81 a more symmetric form of the trace formula for Bethe vectors derived in [GMR19]. The new  
 82 formula is presented in Proposition 3.12. Its main ingredient is the so-called “master” creation  
 83 operator, defined by formula (3.55). Low rank examples of the “master” creation operator are  
 84 presented in Example 3.11.

85 The second main result, Propositions 4.4 and 4.6, present recurrence relation for  $Y^\pm(\mathfrak{gl}_{2n})$ -  
86 and  $Y^+(\mathfrak{gl}_{2n+1})$ -based Bethe vectors, respectively. Schematically, they are of the form

$$\begin{aligned} \Psi^{(m_1, \dots, m_n)} &= \sum_{1 \leq i \leq n} s_{i, 2n-i+1} \Psi^{(m_1, \dots, m_{i-1}, m_i-2, \dots, m_{n-1}-2, m_n-1)} \\ &+ \sum_{1 \leq i < j \leq n} (s_{i, 2n-j+1} + s_{j, 2n-i+1}) \Psi^{(m_1, \dots, m_{i-1}, m_i-1, \dots, m_{j-1}-1, m_j-2, \dots, m_{n-1}-2, m_n-1)} \end{aligned} \quad (1.1)$$

87 in the even case and

$$\begin{aligned} \Psi^{(m_1, \dots, m_n)} &= \sum_{1 \leq i \leq n} s_{i, n+1} \Psi^{(m_1, \dots, m_{i-1}, m_i-1, \dots, m_{n-1}-1, m_n-1)} \\ &+ \sum_{1 \leq i < n} (s_{i, n+2} + s_{n, n+i+2}) \Psi^{(m_1, \dots, m_{i-1}, m_i-1, \dots, m_{n-1}-1, m_n-2)} \\ &+ \sum_{1 \leq i \leq n} s_{i, 2n-i+2} \Psi^{(m_1, \dots, m_{i-1}, m_i-2, \dots, m_{n-1}-2, m_n-2)} \\ &+ \sum_{1 \leq i < j < n} (s_{i, 2n-j+2} + s_{j, 2n-i+2}) \Psi^{(m_1, \dots, m_{i-1}, m_i-1, \dots, m_{j-1}-1, m_j-2, \dots, m_{n-1}-2, m_n-2)} \end{aligned} \quad (1.2)$$

88 in the odd case. Here  $m_i$ 's indicate excitation numbers associated with the  $i$ -th simple root of  
89 the boundary symmetry algebra,  $s_{ij}$ 's represent generating series of the twisted Yangian, and  
90 all scalar factors and spectral parameter dependencies are omitted. These relations are com-  
91 patible with the weight grading of twisted Yangian (see Appendix A.1). Repeated application  
92 of relations (1.1) and (1.2) allows us to express Bethe vectors  $\Psi^{(m_1, \dots, m_n)}$  in terms of those with  
93 no level- $n$  excitations, i.e. with  $m_n = 0$ . The latter Bethe vectors obey  $Y(\mathfrak{gl}_n)$ -type recurrence  
94 relations of the form [HL<sup>+</sup>17b]

$$\Psi^{(m_1, \dots, m_{n-1}, 0)} = \sum_{1 \leq i < n} s_{i, n} \Psi^{(m_1, \dots, m_{i-1}, m_i-1, \dots, m_{n-1}-1, 0)} \quad (1.3)$$

95 the explicit form of which is recalled in Appendix A.3. This feature is explained in Remark 3.3.  
96 Recurrence relations (1.1) and (1.2) are rather complex, especially in the odd case. However,  
97 low rank cases, explicitly stated in Examples 4.5 and 4.7, are manageable for practical compu-  
98 tations. Moreover, the known results of  $Y(\mathfrak{gl}_n)$ -based models [HL<sup>+</sup>17a, HL<sup>+</sup>17b, HL<sup>+</sup>18a, HL<sup>+</sup>20]  
99 can be employed after the first step of nesting.

100 The paper is organised as follows. In Section 2 we introduce notation used throughout  
101 the paper and recall the necessary algebraic properties of twisted Yangians. In Section 3 we  
102 present the algebraic Bethe ansatz: Bethe vectors, their eigenvalues and the corresponding  
103 Bethe equations. We consider both even and odd cases simultaneously giving a coherent frame-  
104 work needed for obtaining recurrence relations. In Section 4 we obtain recurrence relations  
105 and present a proof of the technical Lemma 3.8. In Appendix A we recall weight grading of  
106  $Y^\pm(\mathfrak{gl}_N)$ , a recurrence relation for  $Y(\mathfrak{gl}_n)$ -based Bethe vectors, and provide a proof of commu-  
107 tativity of transfer matrices.

## 108 2 Definitions and preliminaries

109 Throughout the manuscript the middle alphabet letters  $i, j, k, \dots$  will be used to denote inte-  
110 ger numbers, letters  $u, v, w, \dots$  will denote either complex numbers or formal parameters, and  
111 letters  $a$  and  $b$  (often decorated with additional indices) will be used to label vector spaces.

## 112 2.1 Lie algebras

113 Choose  $N \geq 2$ . Let  $\mathfrak{gl}_N$  denote the general linear Lie algebra and let  $e_{ij}$  with  $1 \leq i, j \leq N$  be  
114 the standard basis elements of  $\mathfrak{gl}_N$  satisfying

$$[e_{ij}, e_{kl}] = \delta_{jk}e_{il} - \delta_{il}e_{kj}. \quad (2.1)$$

115 The orthogonal Lie algebra  $\mathfrak{so}_N$  and the symplectic Lie algebra  $\mathfrak{sp}_N$  can be regarded as subal-  
116 gebras of  $\mathfrak{gl}_N$  as follows. For any  $1 \leq i, j \leq N$  set  $\theta_{ij} := \theta_i\theta_j$  with  $\theta_i := 1$  in the orthogonal  
117 case and  $\theta_i := \delta_{i>N/2} - \delta_{i \leq N/2}$  in the symplectic case. Introduce elements  $f_{ij} := e_{ij} - \theta_{ij}e_{\bar{j}\bar{i}}$   
118 with  $\bar{i} := N - i + 1$  and  $\bar{j} := N - j + 1$ . These elements satisfy the relations

$$[f_{ij}, f_{kl}] = \delta_{jk}f_{il} - \delta_{il}f_{kj} + \theta_{ij}(\delta_{\bar{j}\bar{l}}f_{k\bar{i}} - \delta_{\bar{i}\bar{k}}f_{\bar{j}l}), \quad (2.2)$$

$$f_{ij} + \theta_{ij}f_{\bar{j}\bar{i}} = 0, \quad (2.3)$$

119 which in fact are the defining relations of  $\mathfrak{so}_N$  and  $\mathfrak{sp}_N$ . It will be convenient to denote both  
120 algebras by  $\mathfrak{g}_N$ . Write  $N = 2n$  or  $N = 2n + 1$ . In this work we will focus on the following chain  
121 of Lie algebras

$$\mathfrak{gl}_N \supset \mathfrak{g}_N \supset \mathfrak{gl}_n \supset \mathfrak{gl}_{n-1} \supset \cdots \supset \mathfrak{gl}_2,$$

122 where  $\mathfrak{gl}_n, \mathfrak{gl}_{n-1}, \dots, \mathfrak{gl}_2$  are subalgebras of  $\mathfrak{g}_N$  generated by  $f_{ij}$  with  $1 \leq i, j \leq k$  and  
123  $k = n, n-1, \dots, 2$ , respectively.

## 124 2.2 Matrix operators

125 For any  $k \in \mathbb{N}$  let  $E_{ij}^{(k)} \in \text{End}(\mathbb{C}^k)$  with  $1 \leq i, j \leq k$  denote the standard matrix units with  
126 entries in  $\mathbb{C}$  and let  $E_i^{(k)} \in \mathbb{C}^k$  with  $1 \leq i \leq k$  denote the standard basis vectors of  $\mathbb{C}^k$  so that  
127  $E_{ij}^{(k)} E_l^{(k)} = \delta_{jl}E_i^{(k)}$ . We will frequently use the barred index notation

$$E_{i\bar{j}}^{(k)} := E_{k-i+1, k-j+1}^{(k)}, \quad E_{\bar{i}}^{(k)} := E_{k-i+1}^{(k)}. \quad (2.4)$$

128 Introduce matrix operators

$$I^{(k,k)} := \sum_{i,j} E_{ii}^{(k)} \otimes E_{jj}^{(k)}, \quad P^{(k,k)} := \sum_{i,j} E_{ij}^{(k)} \otimes E_{ji}^{(k)}, \quad Q^{(k,k)} := \sum_{i,j} E_{ij}^{(k)} \otimes E_{\bar{i}\bar{j}}^{(k)}, \quad (2.5)$$

129 where the tensor product is defined over  $\mathbb{C}$ . We will always assume that the summation is over  
130 all admissible values, if not stated otherwise. Note that the operator  $Q^{(k,k)}$  is an idempotent  
131 operator,  $(Q^{(k,k)})^2 = kQ^{(k,k)}$ , obtained by partially transforming the permutation operator  
132  $P^{(k,k)}$  with the transposition  $\omega : E_{ij}^{(k)} \mapsto E_{\bar{j}\bar{i}}^{(k)}$ , that is,  $Q^{(k,k)} = (\text{id} \otimes \omega)(P^{(k,k)}) = (\omega \otimes \text{id})(P^{(k,k)})$ .

133 Next, we introduce a matrix-valued rational function

$$R^{(k,k)}(u) := I^{(k,k)} - u^{-1}P^{(k,k)} \quad (2.6)$$

134 called the *Yang's R-matrix*. It is a solution of the quantum Yang-Baxter equation in  $\mathbb{C}^k \otimes \mathbb{C}^k \otimes \mathbb{C}^k$ :

$$R_{12}^{(k,k)}(u-v)R_{13}^{(k,k)}(u-z)R_{23}^{(k,k)}(v-z) = R_{23}^{(k,k)}(v-z)R_{13}^{(k,k)}(u-z)R_{12}^{(k,k)}(u-v). \quad (2.7)$$

135 Here the subscript notation indicates the tensor spaces the matrix operators act on. We will  
136 use such a subscript notation throughout the manuscript. We will also make use the partially  
137  $\omega$ -transposed  $R$ -matrix

$$\widehat{R}^{(k,k)}(u) := (\text{id} \otimes \omega)(R^{(k,k)}(u)) = I^{(k,k)} - u^{-1}Q^{(k,k)} \quad (2.8)$$

138 satisfying a transposed version of (2.7):

$$R_{12}^{(k,k)}(u-v)\widehat{R}_{23}^{(k,k)}(v-z)\widehat{R}_{13}^{(k,k)}(u-z) = \widehat{R}_{13}^{(k,k)}(u-z)\widehat{R}_{23}^{(k,k)}(v-z)R_{12}^{(k,k)}(u-v). \quad (2.9)$$

### 139 2.3 Twisted Yangian $Y^\pm(\mathfrak{gl}_N)$

140 We briefly recall the necessary details of the “ $\rho$ -shifted” twisted Yangian  $Y^\pm(\mathfrak{gl}_N)$  adhering  
 141 closely to [AC<sup>+</sup>06a, GMR19] (see also [Ols92] and Chapters 2 and 4 in [Mol07]); here the  
 142 upper (resp. lower) sign in  $\pm$  corresponds to the orthogonal (resp. symplectic) case. The  
 143 parameter  $\rho \in \mathbb{C}$  is introduced to accommodate applications to Yang-Mills theories and con-  
 144 densed matter systems, where  $\rho$  plays a role of a boundary parameter, and integrable overlaps,  
 145 where  $\rho$  appears as an integer parameter in the nesting procedure.

146 Twisted Yangian  $Y^\pm(\mathfrak{gl}_N)$  is a unital associative  $\mathbb{C}$ -algebra with generators  $s_{ij}[r]$  where  
 147  $1 \leq i, j \leq N$  and  $r \in \mathbb{N}$ . The defining relations, written in terms of the generating series  
 148  $s_{ij}(u) := \delta_{ij} + \sum_{r \geq 1} s_{ij}[r] u^{-r}$ , where  $u$  is a formal variable, are

$$\begin{aligned} [s_{ij}(u), s_{kl}(v)] &= \frac{1}{u-v} (s_{kj}(u) s_{il}(v) - s_{kj}(v) s_{il}(u)) \\ &\quad - \frac{1}{u-\tilde{v}} (\theta_{j\bar{k}} s_{i\bar{k}}(u) s_{\bar{j}l}(v) - \theta_{\bar{i}l} s_{k\bar{i}}(v) s_{\bar{j}l}(u)) \\ &\quad + \frac{1}{(u-v)(u-\tilde{v})} \theta_{i\bar{j}} (s_{k\bar{i}}(u) s_{\bar{j}l}(v) - s_{k\bar{i}}(v) s_{\bar{j}l}(u)) \end{aligned} \quad (2.10)$$

149 and

$$\theta_{ij} s_{\bar{j}\bar{i}}(\tilde{u}) = s_{ij}(u) \pm \frac{s_{ij}(u) - s_{ij}(\tilde{u})}{u - \tilde{u}}. \quad (2.11)$$

150 Here  $\bar{i} = N - i + 1$ ,  $\bar{j} = N - j + 1$ , etc., and  $\tilde{u} := -u - \rho$ ,  $\tilde{v} := -v - \rho$ . These relations can be  
 151 cast in a matrix form as follows. Combine the series  $s_{ij}(u)$  into the generating matrix

$$S^{(N)}(u) := \sum_{i,j} E_{ij}^{(N)} \otimes s_{ij}(u) \quad (2.12)$$

152 The defining relations (2.10) and (2.11) are then equivalent to the twisted reflection equation

$$\begin{aligned} R_{12}^{(N,N)}(u-v) S_1^{(N)}(u) \widehat{R}_{12}^{(N,N)}(\tilde{v}-u) S_2^{(N)}(v) \\ = S_2^{(N)}(v) \widehat{R}_{12}^{(N,N)}(\tilde{v}-u) S_1^{(N)}(u) R_{12}^{(N,N)}(u-v) \end{aligned} \quad (2.13)$$

153 and the symmetry relation

$$\omega(S^{(N)}(\tilde{u})) = S^{(N)}(u) \pm \frac{S^{(N)}(u) - S^{(N)}(\tilde{u})}{u - \tilde{u}}. \quad (2.14)$$

### 154 2.4 Block decomposition

155 Set  $\hat{n} := n$  when  $N = 2n$  and  $\hat{n} := n + 1$  when  $N = 2n + 1$ . Then define  $\hat{n} \times \hat{n}$  dimensional  
 156 matrix operators

$$\begin{aligned} A^{(\hat{n})}(u) &= \sum_{i,j} E_{ij}^{(\hat{n})} \otimes s_{ij}(u), & B^{(\hat{n})}(u) &= \sum_{i,j} E_{ij}^{(\hat{n})} \otimes s_{i,n+j}(u), \\ C^{(\hat{n})}(u) &= \sum_{i,j} E_{ij}^{(\hat{n})} \otimes s_{n+i,j}(u), & D^{(\hat{n})}(u) &= \sum_{i,j} E_{ij}^{(\hat{n})} \otimes s_{n+i,n+j}(u). \end{aligned} \quad (2.15)$$

157 These operators are matrix analogues of the conventional  $a$ ,  $b$ ,  $c$  and  $d$  operators of the six-  
 158 vertex type algebraic Bethe ansatz. The exchange relations that we will need are [GMR19]:

$$\begin{aligned} A_b^{(\hat{n})}(v) B_a^{(\hat{n})}(u) &= R_{ab}^{(\hat{n},\hat{n})}(u-v) B_a^{(\hat{n})}(u) \widehat{R}_{ab}^{(\hat{n},\hat{n})}(\tilde{v}-u) A_b^{(\hat{n})}(v) \\ &\quad + \frac{P_{ab}^{(\hat{n},\hat{n})} B_a^{(\hat{n})}(v) \widehat{R}_{ab}^{(\hat{n},\hat{n})}(\tilde{v}-u) A_b^{(\hat{n})}(u)}{u-v} \mp \frac{B_b^{(\hat{n})}(v) Q_{ab}^{(\hat{n},\hat{n})} D_a^{(\hat{n})}(u)}{u-\tilde{v}}, \end{aligned} \quad (2.16)$$

$$\begin{aligned} R_{ab}^{(\hat{n},\hat{n})}(u-v) B_a^{(\hat{n})}(u) \widehat{R}_{ab}^{(\hat{n},\hat{n})}(\tilde{v}-u) B_b^{(\hat{n})}(v) \\ = B_b^{(\hat{n})}(v) \widehat{R}_{ab}^{(\hat{n},\hat{n})}(\tilde{v}-u) B_a^{(\hat{n})}(u) R_{ab}^{(\hat{n},\hat{n})}(u-v), \end{aligned} \quad (2.17)$$

$$\begin{aligned}
& R_{ab}^{(\hat{n},\hat{n})}(u-v)A_a^{(\hat{n})}(u)A_b^{(\hat{n})}(v) - A_b^{(\hat{n})}(v)A_a^{(\hat{n})}(u)R_{ab}^{(\hat{n},\hat{n})}(u-v) \\
&= \mp \frac{R_{ab}^{(\hat{n},\hat{n})}(u-v)B_a^{(\hat{n})}(u)Q_{ab}^{(\hat{n},\hat{n})}C_b^{(\hat{n})}(v) - B_b^{(\hat{n})}(v)Q_{ab}^{(\hat{n},\hat{n})}C_a^{(\hat{n})}(u)R_{ab}^{(\hat{n},\hat{n})}(u-v)}{u-\tilde{v}}, \quad (2.18)
\end{aligned}$$

$$\begin{aligned}
C_a^{(\hat{n})}(u)A_b^{(\hat{n})}(v) &= A_b^{(\hat{n})}(v)\widehat{R}_{ab}^{(\hat{n},\hat{n})}(\tilde{v}-u)C_a^{(\hat{n})}(u)R_{ab}^{(\hat{n},\hat{n})}(u-v) \\
&+ \frac{P_{ab}^{(\hat{n},\hat{n})}A_a^{(\hat{n})}(u)\widehat{R}_{ab}^{(\hat{n},\hat{n})}(\tilde{v}-u)C_b^{(\hat{n})}(v)}{u-v} \mp \frac{D_a^{(\hat{n})}(u)Q_{ab}^{(\hat{n},\hat{n})}C_b^{(\hat{n})}(v)}{u-\tilde{v}} \quad (2.19)
\end{aligned}$$

159 and

$$\widehat{D}^{(\hat{n})}(\tilde{u}) = A^{(\hat{n})}(u) \pm \frac{A^{(\hat{n})}(u) - A^{(\hat{n})}(\tilde{u})}{u-\tilde{u}}, \quad \pm \widehat{B}^{(\hat{n})}(\tilde{u}) = B^{(\hat{n})}(u) \pm \frac{B^{(\hat{n})}(u) - B^{(\hat{n})}(\tilde{u})}{u-\tilde{u}}. \quad (2.20)$$

160 Here indices  $a$  and  $b$  label two distinct copies of  $\text{End}(\mathbb{C}^{\hat{n}})$ , and  $\widehat{D}^{(\hat{n})}(\tilde{u})$ ,  $\widehat{B}^{(\hat{n})}(\tilde{u})$  are  $\omega$ -transposed  
161 matrices. Taking matrix coefficients of (2.16)–(2.20) one obtains relations among generating  
162 series that coincide with those given by the defining relations (2.10) and (2.11).

163 *Remark 2.1.* In the  $\hat{n} = n + 1$  case all four operators in (2.15) are “overlapping”. For example,  
164 when  $N = 3$ , we have  $\hat{n} = n + 1 = 2$  giving

$$\begin{aligned}
A^{(\hat{n})}(u) &= \begin{pmatrix} s_{11}(u) & s_{12}(u) \\ s_{21}(u) & s_{22}(u) \end{pmatrix}, & B^{(\hat{n})}(u) &= \begin{pmatrix} s_{12}(u) & s_{13}(u) \\ s_{22}(u) & s_{23}(u) \end{pmatrix}, \\
C^{(\hat{n})}(u) &= \begin{pmatrix} s_{21}(u) & s_{22}(u) \\ s_{31}(u) & s_{32}(u) \end{pmatrix}, & D^{(\hat{n})}(u) &= \begin{pmatrix} s_{22}(u) & s_{23}(u) \\ s_{32}(u) & s_{33}(u) \end{pmatrix}.
\end{aligned}$$

165 We will mostly be interested in the  $A$  and  $B$  operators. The  $A$  operator will be used to construct a  
166 transfer matrix of the spin chain and the  $B$  operator will be used to construct creation operators.  
167 Both  $A$  and  $B$  operators include generating series  $s_{i\hat{n}}(u)$  with  $1 \leq i \leq n$  associated with the short  
168 root of  $\mathfrak{so}_{2n+1}$ . These series will be used to construct level- $n$  creation operator and should only  
169 be considered as elements of the  $B$  operator. Likewise, the “middle” generating series  $s_{\hat{n}\hat{n}}(u)$   
170 is also included in both  $A$  and  $B$  operators (and  $C$  and  $D$ ), but should only be considered as  
171 an element of the  $A$  operator. These issues will be resolved by restricting to the upper-left  
172  $(n-1) \times (n-1)$ -dimensional submatrix of the  $A$  operator (such a restriction is compatible with  
173 the  $AB$  exchange relation, see Lemma 3.5) and by explicitly computing the action of  $s_{\hat{n}\hat{n}}(u)$  on  
174 level- $n$  Bethe vectors (see Lemma 3.8).

## 175 3 Bethe ansatz

### 176 3.1 Quantum space

177 We study spin chains with the full quantum space given by

$$L^{(n)} := L(\boldsymbol{\lambda}^{(1)}) \otimes \cdots \otimes L(\boldsymbol{\lambda}^{(\ell)}) \otimes M(\boldsymbol{\mu}) \quad (3.1)$$

178 where  $\ell \in \mathbb{N}$  is the length of the chain, each  $L(\boldsymbol{\lambda}^{(i)})$  and  $M(\boldsymbol{\mu})$  are finite-dimensional irreducible  
179 highest-weight representations of  $\mathfrak{gl}_N$  and  $\mathfrak{g}_N$ , respectively, and the  $N$ -tuples  $\boldsymbol{\lambda}^{(1)}$  and  $\boldsymbol{\mu}$  are  
180 their highest weights. We will say that  $L^{(n)}$  is a *level- $n$  quantum space*.

181 The space  $L^{(n)}$  can be equipped with a structure of a left  $Y^\pm(\mathfrak{g}_N)$ -module as follows. Intro-  
182 duce Lax operators

$$\mathcal{L}^{(N)}(u) := \sum_{i,j} E_{ij}^{(N)} \otimes (\delta_{ij} - u^{-1}e_{ji}), \quad (3.2)$$

$$\mathcal{M}^{(N)}(u) := \sum_{i,j} E_{ij}^{(N)} \otimes (\delta_{ij} - u^{-1}f_{ji}). \quad (3.3)$$

183 Choose an  $\ell$ -tuple  $\mathbf{c} = (c_1, \dots, c_\ell)$  of distinct complex parameters. Then for any  $\xi \in L^{(n)}$  the  
 184 action of  $Y^\pm(\mathfrak{gl}_N)$  is given by

$$S_a^{(N)}(u) \cdot \xi = \prod_i^{\rightarrow} \mathcal{L}_{ai}^{(N)}(u - c_i) \mathcal{M}_{a, \ell+1}^{(N)}(u + (\rho \pm 1)/2) \prod_i^{\leftarrow} \widehat{\mathcal{L}}_{ai}^{(N)}(\tilde{u} - c_i) \cdot \xi \quad (3.4)$$

185 where the subscript  $a$  labels the matrix space of  $S^{(N)}$  and the subscripts  $i = 1, \dots, \ell$  and  $\ell + 1$   
 186 label the individual tensorands of the space  $L^{(n)}$ , which we call *bulk* and *boundary* quantum  
 187 spaces. The bulk spaces are *evaluation representations* of  $Y(\mathfrak{gl}_N)$  and the boundary space is  
 188 an *evaluation representation* of  $Y^\pm(\mathfrak{gl}_N)$ . Moreover, since  $L^{(n)}$  is finite-dimensional, the formal  
 189 variable  $u$  can be evaluated to any complex number, not equal to any  $c_i$ ,  $\tilde{c}_i$ , and  $-(\rho \pm 1)/2$ .

190 Let  $1_{\lambda^{(i)}}$  and  $1_\mu$  denote highest-weight vectors of  $L(\lambda^{(i)})$  and  $M(\mu)$ , respectively. Set

$$\eta := 1_{\lambda^{(1)}} \otimes \dots \otimes 1_{\lambda^{(\ell)}} \otimes 1_\mu. \quad (3.5)$$

191 Then  $s_{ij}(u) \cdot \eta = 0$  if  $i > j$  and  $s_{ii}(u) \cdot \eta = \mu_i(u) \eta$  where

$$\mu_i(u) := \frac{u + (\rho \pm 1)/2 - \mu_i}{u + (\rho \pm 1)/2} \prod_{j \leq \ell} \frac{u - c_j - \lambda_i^{(j)}}{u - c_i} \cdot \frac{\tilde{u} - c_j - \lambda_i^{(j)}}{\tilde{u} - c_i}. \quad (3.6)$$

192 Note that  $\mu_{N-i+1} = -\mu_i$  and  $\mu_{\hat{n}} = 0$  when  $\hat{n} = n + 1$ .

193 An important property of  $L^{(n)}$  is that the subspace  $(L^{(n)})^0 \subset L^{(n)}$ , annihilated by  $s_{ij}(u)$   
 194 with  $i > n$ ,  $j \leq \hat{n}$  and  $i > j$ , is isomorphic to an  $(\ell + 1)$ -fold tensor product of irreducible  $\mathfrak{gl}_n$   
 195 representations. Its subspace  $(L^{(n)})^1 \subset (L^{(n)})^0$ , annihilated by  $s_{ni}(u)$  with  $i < n$ , is isomorphic  
 196 to an  $(\ell + 1)$ -fold tensor product of irreducible  $\mathfrak{gl}_{n-1}$  representations. This can be continued  
 197 to give the following chain of (sub)spaces

$$L^{(n)} \supset (L^{(n)})^0 \supset (L^{(n)})^1 \supset \dots \supset (L^{(n)})^{n-1} \quad (3.7)$$

198 where  $(L^{(n)})^0, (L^{(n)})^1, \dots, (L^{(n)})^{n-1}$  are isomorphic to  $(\ell + 1)$ -fold tensor products of irreducible  
 199 finite-dimensional  $\mathfrak{gl}_n, \mathfrak{gl}_{n-1}, \dots, \mathfrak{gl}_2$  representations, respectively. This property ensures that  
 200 nested algebraic Bethe ansatz techniques can be applied.

### 201 3.2 Nested quantum spaces

202 Choose an  $n$ -tuple  $\mathbf{m} := (m_1, \dots, m_n)$  of non-negative integers, the excitation (magnon) num-  
 203 bers. For each  $m_k$  assign an  $m_k$ -tuple  $\mathbf{u}^{(k)} := (u_1^{(k)}, \dots, u_{m_k}^{(k)})$  of complex parameters (off-shell  
 204 Bethe roots) and an  $m_k$ -tuple  $\mathbf{a}^k := (a_1^k, \dots, a_{m_k}^k)$  of labels, except that for  $m_n$  we assign two  
 205  $m_n$ -tuples of labels,  $\hat{\mathbf{a}} := (\hat{a}_1, \dots, \hat{a}_{m_n})$  and  $\check{\mathbf{a}} := (\check{a}_1, \dots, \check{a}_{m_n})$ . We will often use the following  
 206 shorthand notation:

$$\mathbf{u}^{(k\dots l)} := (\mathbf{u}^{(k)}, \mathbf{u}^{(k+1)}, \dots, \mathbf{u}^{(l)}). \quad (3.8)$$

207 We will assume that  $\mathbf{u}^{(k\dots k)} = \mathbf{u}^k$  and that  $\mathbf{u}^{(k\dots l)}$  is an empty tuple if  $k > l$  so that, for instance,

$$f(\mathbf{u}^{(1\dots k)}, \mathbf{u}^{(k\dots l)}) = f(\mathbf{u}^{(1\dots k)})$$

208 for any function or operator  $f$  when  $k \geq l$ . For any tuples  $\mathbf{u}$  and  $\mathbf{v}$  of complex parameters we  
 209 set

$$f^\pm(u_i, v_j) := \frac{u_i - v_j \pm 1}{u_i - v_j}, \quad f^\pm(\mathbf{u}, \mathbf{v}) := \prod_{i,j} f^\pm(u_i, v_j), \quad \frac{1}{\mathbf{u} - \mathbf{v}} := \prod_{i,j} \frac{1}{u_i - v_j} \quad (3.9)$$

210 where the products are over all admissible indices  $i$  and  $j$ .

211 Let  $V_{a_i^k}^{(k)}$  denote a copy of  $\mathbb{C}^k$  labelled by “ $a_i^k$ ” and let  $W_{a^k}^{(k)}$  be defined by

$$W_{a^k}^{(k)} := V_{a_1^k}^{(k)} \otimes \cdots \otimes V_{a_{m_k}^k}^{(k)} \cong (\mathbb{C}^k)^{\otimes m_k}. \quad (3.10)$$

212 Labels  $a_i^k$  will be used to trace the action of matrix operators. We illustrate this property with  
 213 an example. Let  $\xi = \xi_{a_1^k} \otimes \cdots \otimes \xi_{a_{m_k}^k} \in W_{a^k}^{(k)}$  and let  $M_{a_j^k}^{(k)} \in \text{End}(V_{a_j^k}^{(k)})$  be a matrix operator  
 214 acting in the space labelled  $a_j^k$ . Then

$$M_{a_j^k}^{(k)} \xi = \xi_{a_1^k} \otimes \cdots \otimes \xi_{a_{j-1}^k} \otimes \left( M_{a_j^k}^{(k)} \xi_{a_j^k} \right) \otimes \xi_{a_{j+1}^k} \otimes \cdots \otimes \xi_{a_{m_k}^k}.$$

215 Let  $V_{\hat{a}_i}^{(\hat{n})}, V_{\hat{a}_i}^{(\hat{n})} \cong \mathbb{C}^{\hat{n}}$  and  $W_{\hat{a}}^{(\hat{n})}, W_{\hat{a}}^{(\hat{n})} \cong (\mathbb{C}^{\hat{n}})^{\otimes m_n}$  be defined analogously to (3.10). We define  
 216 a *level-( $n-1$ ) quantum space* by

$$L^{(n-1)} := W_{\hat{a}}^{(\hat{n})} \otimes W_{\hat{a}}^{(\hat{n})} \otimes (L^{(n)})^0. \quad (3.11)$$

217 When  $\hat{n} = n + 1$ , we additionally introduce “reduced” vector spaces

$$\overline{W}_{\hat{a}}^{(\hat{n})} := \overline{V}_{\hat{a}_1}^{(\hat{n})} \otimes \cdots \otimes \overline{V}_{\hat{a}_{m_n}}^{(\hat{n})}, \quad \overline{W}_{\hat{a}}^{(\hat{n})} := \overline{V}_{\hat{a}_1}^{(\hat{n})} \otimes \cdots \otimes \overline{V}_{\hat{a}_{m_n}}^{(\hat{n})} \quad (3.12)$$

218 where

$$\overline{V}_{\hat{a}_i}^{(\hat{n})} := \text{span}_{\mathbb{C}}\{E_j^{(\hat{n})} : 2 \leq j \leq \hat{n}\} \subset V_{\hat{a}_i}^{(\hat{n})}, \quad \overline{V}_{\hat{a}_i}^{(\hat{n})} := \text{span}_{\mathbb{C}}\{E_1^{(\hat{n})}\} \subset V_{\hat{a}_i}^{(\hat{n})}. \quad (3.13)$$

219 Specifically,  $\overline{W}_{\hat{a}}^{(\hat{n})}$  is isomorphic to  $(\mathbb{C}^n)^{\otimes m_n}$  and  $\overline{W}_{\hat{a}}^{(\hat{n})}$  a 1-dimensional vector space. We then  
 220 define a *reduced level-( $n-1$ ) quantum space* by

$$\overline{L}^{(n-1)} := \overline{W}_{\hat{a}}^{(\hat{n})} \otimes \overline{W}_{\hat{a}}^{(\hat{n})} \otimes (L^{(n)})^0 \subset L^{(n-1)}. \quad (3.14)$$

221 The spaces  $L^{(n-1)}$  and  $\overline{L}^{(n-1)}$  will serve as the full (nested) quantum spaces of the  $Y(\mathfrak{gl}_n)$ -  
 222 based models obtained after the first step of nesting in the even and odd cases, respectively;  
 223 see Remark 3.3.

224 Then, for each  $k = n - 2, n - 3, \dots, 1$  we define a *level- $k$  quantum space* by

$$L^{(k)} := W_{a^{k+1}}^{(k+1)} \otimes (L^{(k+1)})^0 \quad (3.15)$$

225 where  $(L^{(k+1)})^0$  is a *level-( $k+1$ ) vacuum subspace* given by

$$(L^{(k+1)})^0 := (W_{a^{k+2}}^{(k+2)})^0 \otimes \cdots \otimes (W_{a^{n-1}}^{(n-1)})^0 \otimes (W_{\hat{a}}^{(\hat{n})})^0 \otimes (W_{\hat{a}}^{(\hat{n})})^0 \otimes (L^{(n)})^{n-k-1} \subset L^{(k+1)} \quad (3.16)$$

226 where

$$(W_{a^{k+2}}^{(k+2)})^0 \subset W_{a^{k+2}}^{(k+2)}, \quad \dots, \quad (W_{a^{n-1}}^{(n-1)})^0 \subset W_{a^{n-1}}^{(n-1)}, \quad (W_{\hat{a}}^{(\hat{n})})^0 \subset W_{\hat{a}}^{(\hat{n})}, \quad (W_{\hat{a}}^{(\hat{n})})^0 \subset W_{\hat{a}}^{(\hat{n})}$$

227 are 1-dimensional subspaces spanned by vectors

$$E_1^{(k+2)} \otimes \cdots \otimes E_1^{(k+2)}, \quad \dots, \quad E_1^{(n-1)} \otimes \cdots \otimes E_1^{(n-1)}, \quad E_{\hat{1}}^{(\hat{n})} \otimes \cdots \otimes E_{\hat{1}}^{(\hat{n})}, \quad E_{\hat{1}}^{(\hat{n})} \otimes \cdots \otimes E_{\hat{1}}^{(\hat{n})}$$

228 respectively. When  $\hat{n} = n + 1$ , note that  $(L^{(n-1)})^0 \subset \overline{L}^{(n-1)}$ . Moreover,  $(L^{(k+1)})^0 \cong (L^{(n)})^{n-k-1}$   
 229 for  $1 \leq k \leq n - 2$ . The spaces  $L^{(k)}$  will serve as the full (nested) quantum spaces of the  
 230  $Y(\mathfrak{gl}_{k+1})$ -based models obtained after  $n - k$  steps of nesting.



### 231 3.3 Monodromy matrices

232 We will say that the matrix  $S^{(N)}(u)$ , acting in the space  $L^{(n)}$  via (3.4), is a *level- $n$  monodromy*  
 233 *matrix*. In this setting, we will treat  $u$  as a non-zero complex number not equal to any  $c_i$ ,  $\tilde{c}_i$  and  
 234  $-(\rho \pm 1)/2$ . We define a *level- $(n-1)$  nested monodromy matrix*, acting in the space  $L^{(n-1)}$ , by

$$T_a^{(\hat{n})}(v; \mathbf{u}^{(n)}) := \prod_{i \leq m_n}^{\leftarrow} \widehat{R}_{\hat{a}_i a}^{(\hat{n}, \hat{n})}(u_i^{(n)} - v) \prod_{i \leq m_n}^{\leftarrow} \widehat{R}_{\hat{a}_i a}^{(\hat{n}, \hat{n})}(\tilde{u}_i^{(n)} - v) A_a^{(\hat{n})}(v). \quad (3.17)$$

235 When  $\hat{n} = n + 1$ , we introduce a *reduced level- $(n-1)$  nested monodromy matrix*, acting in the  
 236 space  $\overline{L}^{(n-1)}$ , by

$$\overline{T}_a^{(n)}(v; \mathbf{u}^{(n)}) := \prod_{i \leq m_n}^{\leftarrow} \overline{\widehat{R}}_{\hat{a}_i a}^{(n, n)}(u_i^{(n)} - v) [A_a^{(\hat{n})}(v)]^{(n)} \quad (3.18)$$

237 where  $\overline{\widehat{R}}_{\hat{a}_i a}^{(n, n)}$  is the restriction of  $\widehat{R}_{\hat{a}_i a}^{(n, n)}$  to  $\overline{V}_{\hat{a}_i}^{(\hat{n})} \otimes V_a^{(n)} \subset V_{\hat{a}_i}^{(\hat{n})} \otimes V_a^{(\hat{n})}$  (recall (2.8) and (3.13)),  
 238 and the notation  $[ \ ]^{(n)}$  means the restriction to the upper-left  $(n \times n)$ -dimensional submatrix;  
 239 this notation will be used throughout the manuscript.

240 **Lemma 3.1.** *When  $\hat{n} = n + 1$ , in the space  $\overline{L}^{(n-1)}$  we have the equality of operators*

$$[\overline{T}_a^{(n)}(v; \mathbf{u}^{(n)})]^{(n)} = \overline{T}_a^{(n)}(v; \mathbf{u}^{(n)}). \quad (3.19)$$

241 Moreover, the space  $\overline{L}^{(n-1)}$  is stable under the action of  $\overline{T}_a^{(n)}(v; \mathbf{u}^{(n)})$ .

242 *Proof.* From (2.8) observe that

$$[\widehat{R}_{ba}^{(\hat{n}, \hat{n})}(v)]_{kl} E_j^{(\hat{n})} = \delta_{kl} E_j^{(\hat{n})} - v^{-1} \delta_{\hat{n}-l+1, j} E_{\hat{n}-k+1}^{(\hat{n})} \quad (3.20)$$

243 where  $[ \ ]_{kl}$  selects the  $(k, l)$ -th matrix element of  $\widehat{R}_{ba}^{(\hat{n}, \hat{n})}$  in the  $a$ -space; this notation will be  
 244 used throughout the manuscript. Therefore, for any  $1 \leq k, l \leq n$  and any  $\eta \in \overline{W}_a^{(\hat{n})}$ ,  $\zeta \in \overline{W}_a^{(\hat{n})}$ ,  
 245  $\xi \in (L^{(n)})^0$ , viz. (3.14), we have

$$\begin{aligned} & [\overline{T}_a^{(\hat{n})}(v; \mathbf{u}^{(n)})]_{kl} \cdot \eta \otimes \zeta \otimes \xi \\ &= \sum_{p, r} \left[ \prod_{i \leq m_n}^{\leftarrow} \widehat{R}_{\hat{a}_i a}^{(\hat{n}, \hat{n})}(u_i^{(n)} - v) \right]_{kp} \cdot \eta \otimes \left[ \prod_{i \leq m_n}^{\leftarrow} \widehat{R}_{\hat{a}_i a}^{(\hat{n}, \hat{n})}(\tilde{u}_i^{(n)} - v) \right]_{pr} \cdot \zeta \otimes s_{rl}(v) \cdot \xi \\ &= \sum_{p \leq n} \left[ \prod_{i \leq m_n}^{\leftarrow} \widehat{R}_{\hat{a}_i a}^{(\hat{n}, \hat{n})}(u_i^{(n)} - v) \right]_{kp} \cdot \eta \otimes \zeta \otimes s_{pl}(v) \cdot \xi \end{aligned} \quad (3.21)$$

246 since  $s_{\hat{n}l}(v) \cdot \xi = 0$  by definition of  $(L^{(n)})^0$ , and, by (3.20),

$$\left[ \prod_{i \leq m_n}^{\leftarrow} \widehat{R}_{\hat{a}_i a}^{(\hat{n}, \hat{n})}(\tilde{u}_i^{(n)} - v) \right]_{pr} \cdot \zeta = \delta_{pr} \zeta$$

247 when  $r < \hat{n}$  because  $\zeta$  is a scalar multiple of  $E_1^{(\hat{n})} \otimes \cdots \otimes E_1^{(\hat{n})}$ . But

$$\left[ \prod_{i \leq m_n}^{\leftarrow} \widehat{R}_{\hat{a}_i a}^{(\hat{n}, \hat{n})}(u_i^{(n)} - v) \right]_{kp} \cdot \eta \notin \overline{W}_a^{(\hat{n})}$$

248 when  $k, p \leq n$  only if the product includes  $[\widehat{R}_{\hat{a}_i a}^{(\hat{n}, \hat{n})}(u_i^{(n)} - v)]_{\hat{n}r}$  with  $r \leq n$ , but then it must also  
 249 include  $[\widehat{R}_{\hat{a}_i a}^{(\hat{n}, \hat{n})}(u_i^{(n)} - v)]_{r\hat{n}}$  which acts by zero on  $\eta$  since the spaces  $\overline{V}_{\hat{a}_i}^{(\hat{n})}$  have no  $E_1^{(\hat{n})}$ 's. Thus

$$\left[ \prod_{i \leq m_n}^{\leftarrow} \widehat{R}_{\hat{a}_i a}^{(\hat{n}, \hat{n})}(u_i^{(n)} - v) \right]_{kp} \cdot \eta = \left[ \prod_{i \leq m_n}^{\leftarrow} \overline{\widehat{R}}_{\hat{a}_i a}^{(\hat{n}, \hat{n})}(u_i^{(n)} - v) \right]_{kp} \cdot \eta \in \overline{W}_{\hat{a}}^{(\hat{n})} \quad (3.22)$$

250 implying (3.19). To prove the second part of the claim, notice that  $(L^{(n)})^0$  is stable under the  
 251 action of  $s_{pl}(u)$  with  $1 \leq p, l \leq n$ . Indeed, by definition, it is the subspace of  $L^{(n)}$  annihilated by  
 252  $s_{\bar{i}j}(u)$  with  $\bar{i} > n, j \leq \hat{n}$  and  $\bar{i} > j$ . Assuming  $1 \leq i, j, k, l \leq n$ , (2.10) gives  $s_{\bar{i}j}(u)s_{kl}(v) = 0$  in the  
 253 space  $(L^{(n)})^0$  thus proving its stability. The stability of  $\overline{L}^{(n-1)}$  under the action of  $T_a^{(n)}(v; \mathbf{u}^{(n)})$   
 254 then follows immediately from (3.21) and (3.22).  $\square$

255 Next, for each  $k = n-1, n-2, \dots, 2$ , we define a *level-(k-1) nested monodromy matrix*,  
 256 acting in the space  $L^{(k-1)}$ , by

$$T_a^{(k)}(v; \mathbf{u}^{(k \dots n)}) := \prod_{i \leq m_k}^{\leftarrow} \widehat{R}_{\hat{a}_i a}^{(k, k)}(u_i^{(k)} - v) [T_a^{(k+1)}(v; \mathbf{u}^{(k+1 \dots n)})]^{(k)} \quad (3.23)$$

257 where  $T_a^{(k+1)}$  should be  $\overline{T_a^{(k+1)}}$  when  $\hat{n} = n+1$  and  $k = n$ .

258 **Lemma 3.2.** For each  $2 \leq k \leq n$ , the space  $L^{(k-1)}$  is stable under the action of  $T_a^{(k)}(v; \mathbf{u}^{(k \dots n)})$  and

$$\begin{aligned} R_{ab}^{(k, k)}(v-w) T_a^{(k)}(v; \mathbf{u}^{(k \dots n)}) T_b^{(k)}(w; \mathbf{u}^{(k \dots n)}) \\ = T_b^{(k)}(w; \mathbf{u}^{(k \dots n)}) T_a^{(k)}(v; \mathbf{u}^{(k \dots n)}) R_{ab}^{(k, k)}(v-w) \end{aligned} \quad (3.24)$$

259 in this space, except, when  $\hat{n} = n+1$  and  $k = n$ ,  $L^{(k-1)}$  should be  $\overline{L}^{(k-1)}$  and  $T^{(k)}$  should be  $\overline{T^{(k)}}$ .

260 *Proof.* When  $k = n$  and  $\hat{n} = n$ , this was shown in Proposition 3.13 in [GMR19]. When  $k = n$   
 261 and  $\hat{n} = n+1$ , the first part of the claim follows from Lemma 3.1; the second part follows from  
 262 the observation that

$$R_{ab}^{(n, n)}(u-v) [A_a^{(\hat{n})}(u)]^{(n)} [A_b^{(\hat{n})}(v)]^{(n)} = [A_b^{(\hat{n})}(v)]^{(n)} [A_a^{(\hat{n})}(u)]^{(n)} R_{ab}^{(n, n)}(u-v) \quad (3.25)$$

263 in the space  $\overline{L}^{(n-1)}$  and application of the transposed quantum Yang-Baxter equation (2.9).  
 264 The (3.25) follows from (2.18) or directly from (2.10) upon restricting to  $1 \leq i, j, k, l \leq n$ .  
 265 The  $k < n$  cases then follow by the standard arguments.  $\square$

266 *Remark 3.3.* Lemma 3.2 together with (3.17), (3.18) say that  $Y^\pm(\mathfrak{gl}_{2n})$ - and  $Y^+(\mathfrak{gl}_{2n+1})$ -based  
 267 models, after the first step of nesting, are equivalent to  $Y(\mathfrak{gl}_n)$ -based models with off-shell  
 268 Bethe roots given by  $\mathbf{v}^{(1 \dots n-2)} := \mathbf{u}^{(1 \dots n-2)}$  and  $\mathbf{v}^{(n)} := (\mathbf{u}^{(n)}, \tilde{\mathbf{u}}^{(n)})$  in the even case, and  
 269  $\mathbf{v}^{(n)} := \mathbf{u}^{(n)}$  in the odd case. This property will be explored in Section 4.

### 270 3.4 Creation operators

271 We define a *level-n creation operator* by

$$\mathfrak{B}^{(n)}(\mathbf{u}^{(n)}) := \prod_{1 \leq i \leq m_n}^{\leftarrow} \left( \mathfrak{t}_{\hat{a}_i \hat{a}_i}^{(n)}(u_i^{(n)}) \prod_{i < j \leq m_n}^{\rightarrow} \frac{R_{\hat{a}_i \hat{a}_j}^{(\hat{n}, \hat{n})}(\tilde{u}_i^{(n)} - u_j^{(n)})}{(f^-(\tilde{u}_i^{(n)}, u_j^{(n)}))^{\delta_{\hat{n}n}}} \right) \quad (3.26)$$

272 where

$$\mathfrak{t}_{\hat{a}_i \hat{a}_i}^{(n)}(u_i^{(n)}) := \sum_{k, l \leq \hat{n}} (E_k^{(\hat{n})})^* \otimes (E_l^{(\hat{n})})^* \otimes [B_a^{(\hat{n})}(u_i^{(n)})]_{\bar{k}, l} \in (V_{\hat{a}_i}^{(\hat{n})})^* \otimes (V_{\hat{a}_i}^{(\hat{n})})^* \otimes \text{End}(L^{(n)}) \quad (3.27)$$

273 and  $B_a^{(\hat{n})}(u_i^{(n)})$  is the  $B$ -block of the operator in the right hand side of (3.4). The  $R$ -matrices  
 274 in (3.26) are necessary for the wanted order of the  $\check{R}$ -matrices in (3.17), which in turn is  
 275 necessary for Lemma 3.2 to hold. The denominator is an overall normalisation factor.

276 From (3.26) it is clear that  $\mathcal{B}^{(n)}(\mathbf{u}^{(n)})$  satisfies the recurrence relation

$$\mathcal{B}^{(n)}(\mathbf{u}^{(n)}) = \mathfrak{t}_{\check{a}_{m_n} \check{a}_{m_n}}^{(n)}(u_{m_n}^{(n)}) \mathcal{B}^{(n)}(\mathbf{u}^{(n)} \setminus u_{m_n}^{(n)}) \mathcal{R}^{(\hat{n})}(u_{m_n}^{(n)}; \mathbf{u}^{(n)} \setminus u_{m_n}^{(n)}) \quad (3.28)$$

277 where  $\mathcal{B}^{(n)}(\mathbf{u}^{(n)} \setminus u_{m_n}^{(n)})$  is defined via (3.26) except the ranges of products are  $1 \leq i < m_n$  and  
 278  $i < j < m_n$ , and

$$\mathcal{R}^{(\hat{n})}(u_{m_n}^{(n)}; \mathbf{u}^{(n)} \setminus u_{m_n}^{(n)}) := \prod_{1 \leq i < m_n}^{\leftarrow} \frac{R_{\check{a}_i \check{a}_{m_n}}^{(\hat{n}, \hat{n})}(\tilde{u}_i^{(n)} - u_{m_n}^{(n)})}{(f^-(\tilde{u}_i^{(n)}, u_{m_n}^{(n)}))^{\delta_{\hat{n}n}}}. \quad (3.29)$$

279 We will later meet operators  $\mathcal{B}^{(n)}(\mathbf{u}^{(n)} \setminus u_l^{(n)})$  and  $\mathcal{R}^{(\hat{n})}(u_l^{(n)}; \mathbf{u}^{(n)} \setminus u_l^{(n)})$  for any  $l$  that are defined  
 280 analogously except  $u_i^{(n)}$  (resp.  $\tilde{u}_i^{(n)}$ ) should be replaced with  $u_{i+1}^{(n)}$  (resp.  $\tilde{u}_{i+1}^{(n)}$ ) for all  $l \leq i < m_n$ .

281 Next, for each  $k = n-1, n-2, \dots, 1$  we define a *level- $k$  creation operator* by

$$\mathcal{B}^{(k)}(\mathbf{u}^{(k)}; \mathbf{u}^{(k+1\dots n)}) := \prod_{1 \leq i \leq m_k}^{\leftarrow} \mathfrak{t}_{a_i^k}^{(k)}(u_i^{(k)}; \mathbf{u}^{(k+1\dots n)}) \quad (3.30)$$

282 where

$$\mathfrak{t}_{a_i^k}^{(k)}(u_i^{(k)}; \mathbf{u}^{(k+1\dots n)}) := \sum_{1 \leq j \leq k} (E_j^{(k)})_{a_i^k}^* \otimes [T_a^{(k+1)}(u_i^{(k)}; \mathbf{u}^{(k+1\dots n)})]_{j, k+1} \in (V_{a_i^k}^{(k)})^* \otimes \text{End}(L^{(k)}). \quad (3.31)$$

283 Note that  $T_a^{(n)}(u_i^{(n-1)}; \mathbf{u}^{(n)})$  should be replaced with  $\overline{T_a^{(n)}}(u_i^{(n-1)}; \mathbf{u}^{(n)})$  when  $\hat{n} = n+1$ .

284 Parameters of creation operators may be permuted using the following standard result,  
 285 which follows from (2.17); see Lemma 3.6 in [GMR19].

286 **Lemma 3.4.** *The level- $n$  creation operator satisfies*

$$\mathcal{B}^{(n)}(\mathbf{u}^{(n)}) = \mathcal{B}^{(n)}(\mathbf{u}_{i \leftrightarrow i+1}^{(n)}) \check{R}_{\check{a}_{i+1} \check{a}_i}^{(\hat{n}, \hat{n})}(u_i^{(n)} - u_{i+1}^{(n)}) \check{R}_{\check{a}_{i+1} \check{a}_i}^{(\hat{n}, \hat{n})}(u_{i+1}^{(n)} - u_i^{(n)}). \quad (3.32)$$

287 For each  $1 \leq k \leq n-1$  the level- $k$  creation operator satisfies

$$\mathcal{B}^{(k)}(\mathbf{u}^{(k)}; \mathbf{u}^{(k+1\dots n)}) = \mathcal{B}^{(k)}(\mathbf{u}_{i \leftrightarrow i+1}^{(k)}; \mathbf{u}^{(k+1\dots n)}) \check{R}_{a_{i+1}^k a_i^k}^{(k, k)}(u_i^{(k)} - u_{i+1}^{(k)}). \quad (3.33)$$

288 Here the “check”  $\check{R}$ -matrices are defined by

$$\check{R}_{ab}^{(k, k)}(u) := \frac{u}{u-1} P_{ab}^{(k, k)} R_{ab}^{(k, k)}(u) \quad (3.34)$$

289 and  $\mathbf{u}_{i \leftrightarrow i+1}^{(k)}$  denotes the tuple  $\mathbf{u}^{(k)}$  with parameters  $u_i^{(k)}$  and  $u_{i+1}^{(k)}$  interchanged.

290 Recall the notation  $\tilde{v} = -v - \rho$  and introduce the following notation for a symmetrised  
 291 combination of functions or operators

$$\{f(v)\}^v := f(v) + f(\tilde{v}) \quad (3.35)$$

292 and a rational function

$$p(v) := 1 \pm \frac{1}{v - \tilde{v}} \quad (3.36)$$

293 representing the right hand side of the symmetry relation (2.14). The Lemma below rephrases  
 294 the results obtained in [GMR19] in a compact form.

295 **Lemma 3.5.** *The AB exchange relation for the level- $n$  creation operator (3.26) is*

$$\begin{aligned}
 & \{p(v)A_a^{(\hat{n})}(v)\}^v \mathfrak{B}^{(n)}(\mathbf{u}^{(n)}) \\
 &= \mathfrak{B}^{(n)}(\mathbf{u}^{(n)}) \{p(v)T_a^{(\hat{n})}(v; \mathbf{u}^{(n)})\}^v \\
 &+ \sum_i \frac{1}{p(u_i^{(n)})} \left\{ \frac{p(v)}{u_i^{(n)} - v} \mathfrak{t}_{\hat{a}_{m_n} \hat{a}_{m_n}}^{(n)}(v) \right\}^v \mathfrak{B}^{(n)}(\mathbf{u}^{(n)} \setminus u_i^{(n)}) \mathfrak{R}^{(\hat{n})}(u_i^{(n)}; \mathbf{u}^{(n)} \setminus u_i^{(n)}) \\
 &\quad \times \operatorname{Res}_{w \rightarrow u_i^{(n)}} \{p(w)T_a^{(\hat{n})}(w; \mathbf{u}_{\sigma_i}^{(n)})\}^w \prod_{j>i} \check{R}_{\hat{a}_j \hat{a}_{j-1}}^{(\hat{n}, \hat{n})}(u_i^{(n)} - u_j^{(n)}) \check{R}_{\hat{a}_j \hat{a}_{j-1}}^{(\hat{n}, \hat{n})}(u_j^{(n)} - u_i^{(n)}) \quad (3.37)
 \end{aligned}$$

296 where  $\mathbf{u}^{(n)} \setminus u_i^{(n)} := (u_1, \dots, u_{i-1}, u_{i+1}, \dots, u_{m_n})$  and  $\mathbf{u}_{\sigma_i}^{(n)} := (u_1, \dots, u_{i-1}, u_{i+1}, \dots, u_{m_n}, u_i)$ .

297 *Proof.* From [GMR19], relations (2.16) and (2.20) and properties of the  $Q^{(\hat{n}, \hat{n})}$  matrix operator  
 298 (viz. (2.5)) lead to the following exchange relation with a single creation operator

$$\begin{aligned}
 \{p(v)A_a^{(\hat{n})}(v)\}^v \mathfrak{t}_{\hat{a}_i \hat{a}_i}^{(n)}(u_i^{(n)}) &= \mathfrak{t}_{\hat{a}_i \hat{a}_i}^{(n)}(u_i^{(n)}) \{p(v)T_a^{(\hat{n})}(v; u_i^{(n)})\}^v \\
 &+ \frac{1}{p(u_i^{(n)})} \left\{ \frac{p(v)}{u_i^{(n)} - v} \mathfrak{t}_{\hat{a}_i \hat{a}_i}^{(n)}(v) \right\}^v \operatorname{Res}_{w \rightarrow u_i^{(n)}} \{p(w)T_a^{(\hat{n})}(w; u_i^{(n)})\}^w \quad (3.38)
 \end{aligned}$$

299 where  $T_a^{(\hat{n})}(v; u_i^{(n)}) = \widehat{R}_{\hat{a}_i a}(u_i^{(n)} - v) \widehat{R}_{\hat{a}_i a}(\tilde{u}_i^{(n)} - v) A_a^{(\hat{n})}(v)$ . We extend this to the creation operator  
 300 for  $m_n$  excitations by the standard argument. Indeed, the right hand side of the equation  
 301 consists of terms with  $A_a^{(\hat{n})}(u)$  as the rightmost operator, for  $u$  equal to each of  $v, u_1^{(n)}, \dots, u_{m_n}^{(n)}$   
 302 and the corresponding tilded elements. Due to the  $w \mapsto \tilde{w}$  symmetry of  $\{p(w)A_a^{(\hat{n})}(w)\}^w$  in  
 303 (3.38), it is sufficient to find those terms corresponding to  $v, u_1^{(n)}, \dots, u_{m_n}^{(n)}$ .

304 First, we find the term corresponding to  $v$  to be  $\mathfrak{B}^{(n)}(\mathbf{u}^{(n)}) \{p(v)T_a^{(\hat{n})}(v; \mathbf{u}^{(n)})\}^v$ . The re-  
 305 quired order of  $\widehat{R}$ -matrices inside  $T_a^{(\hat{n})}(v; \mathbf{u}^{(n)})$  is a result of Yang-Baxter moves through the  
 306  $R$ -matrices inside  $\mathfrak{B}^{(n)}(\mathbf{u}^{(n)})$ . Using factorisation (3.28) we find the term corresponding to  
 307  $u_{m_n}^{(n)}$  to be

$$\begin{aligned}
 & \frac{1}{p(u_{m_n}^{(n)})} \left\{ \frac{p(v)}{u_{m_n}^{(n)} - v} \mathfrak{t}_{\hat{a}_{m_n} \hat{a}_{m_n}}^{(n)}(v) \right\}^v \mathfrak{B}^{(n)}(\mathbf{u}^{(n)} \setminus u_{m_n}^{(n)}) \\
 & \quad \times \mathfrak{R}^{(\hat{n})}(u_{m_n}^{(n)}; \mathbf{u}^{(n)} \setminus u_{m_n}^{(n)}) \operatorname{Res}_{w \rightarrow u_{m_n}^{(n)}} \{p(w)T_a^{(\hat{n})}(w; \mathbf{u}^{(n)})\}^w.
 \end{aligned}$$

308 This is because, after applying (3.38) to  $\mathfrak{t}_{\hat{a}_{m_n} \hat{a}_{m_n}}^{(n)}(u_{m_n}^{(n)})$ , there can be no further contributions  
 309 from the parameter-swapped term in the subsequent applications of (3.38).

310 To find the remaining terms, we note that Lemma 3.4 allows us to apply any permutation to  
 311 the spectral parameters of the level- $n$  creation operator before applying the above argument.  
 312 By applying the permutation  $\sigma_i : (1, \dots, i-1, i, i+1, \dots, m_n) \mapsto (1, \dots, i-1, i+1, \dots, m_n, i)$ ,  
 313 we obtain the term corresponding to  $u_i^{(n)}$ .  $\square$

314 The Lemma below states  $Y(\mathfrak{gl}_{k+1})$ -based column-nested AB and DB exchange relations.  
 315 They follow from Lemma 3.2 using standard arguments, see e.g. [BR08].

316 **Lemma 3.6.** *The exchange relation for the level- $k$  creation operator (3.30) is*

$$\begin{aligned}
& [T_a^{(k+1)}(v; \mathbf{u}^{(k+1\dots n)})]^{(k)} \mathcal{B}^{(k)}(\mathbf{u}^{(k)}; \mathbf{u}^{(k+1\dots n)}) \\
&= \mathcal{B}^{(k)}(\mathbf{u}^{(k)}; \mathbf{u}^{(k+1\dots n)}) T_a^{(k)}(v; \mathbf{u}^{(k\dots n)}) \\
&+ \sum_i \frac{1}{u_i^{(k)} - v} \theta_{a_{m_k}^{(k)}}^{(k)}(v; \mathbf{u}^{(k+1\dots n)}) \mathcal{B}^{(k)}(\mathbf{u}^{(k)} \setminus u_i^{(k)}; \mathbf{u}^{(k+1\dots n)}) \\
&\quad \times \operatorname{Res}_{w \rightarrow u_i^{(k)}} T_a^{(k)}(w; (\mathbf{u}_{\sigma_i}^{(k)}, \mathbf{u}^{(k+1\dots n)})) \prod_{j>i}^{\rightarrow} \check{R}_{a_j^k a_{j-1}^k}^{(k,k)}(u_i^{(k)} - u_j^{(k)}). \tag{3.39}
\end{aligned}$$

317 *Moreover,*

$$\begin{aligned}
& [T_a^{(k+1)}(v; \mathbf{u}^{(k+1\dots n)})]_{k+1, k+1} \mathcal{B}^{(k)}(\mathbf{u}^{(k)}; \mathbf{u}^{(k+1\dots n)}) \\
&= \mathcal{B}^{(k)}(\mathbf{u}^{(k)}; \mathbf{u}^{(k+1\dots n)}) f^-(v; \mathbf{u}^{(k)}) [T_a^{(k+1)}(v; \mathbf{u}^{(k+1\dots n)})]_{k+1, k+1} \\
&+ \sum_i \frac{1}{u_i^{(k)} - v} \theta_{a_{m_k}^{(k)}}^{(k)}(v; \mathbf{u}^{(k+1\dots n)}) \mathcal{B}^{(k)}(\mathbf{u}^{(k)} \setminus u_i^{(k)}; \mathbf{u}^{(k+1\dots n)}) \\
&\quad \times \operatorname{Res}_{w \rightarrow u_i^{(k)}} f^-(w; \mathbf{u}^{(k)}) [T_a^{(k+1)}(w; \mathbf{u}^{(k+1\dots n)})]_{k+1, k+1} \prod_{j>i}^{\rightarrow} \check{R}_{a_j^k a_{j-1}^k}^{(k,k)}(u_i^{(k)} - u_j^{(k)}). \tag{3.40}
\end{aligned}$$

318 *Here we used the notation*

$$\mathcal{B}^{(k)}(\mathbf{u}^{(k)} \setminus u_i^{(k)}; \mathbf{u}^{(k+1\dots n)}) = \prod_{1 \leq j < i}^{\leftarrow} \theta_{a_j^k}^{(k)}(u_j^{(k)}; \mathbf{u}^{(k+1\dots n)}) \prod_{i \leq j < m_k}^{\leftarrow} \theta_{a_j^k}^{(k)}(u_{j+1}^{(k)}; \mathbf{u}^{(k+1\dots n)}).$$

### 319 3.5 Bethe vectors

320 Recall (3.5) and define a *nested vacuum vector* by

$$\eta^m := (E_1^{(1)})^{\otimes m_1} \otimes \dots \otimes (E_1^{(n-1)})^{\otimes m_{n-1}} \otimes (E_1^{(\hat{n})})^{\otimes m_n} \otimes (E_1^{(\hat{n})})^{\otimes m_n} \otimes \eta. \tag{3.41}$$

321 Note that  $E_1^{(\hat{n})} = E_2^{(n+1)}$  when  $\hat{n} = n + 1$ . For each  $1 \leq k \leq n$  we define a *level- $k$*  (off-shell)

322 Bethe vector with (off-shell) Bethe roots  $\mathbf{u}^{(1\dots k)}$  and free parameters  $\mathbf{u}^{(k+1\dots n)}$  by

$$\Psi(\mathbf{u}^{(1\dots k)} | \mathbf{u}^{(k+1\dots n)}) := \prod_{i \leq k}^{\leftarrow} \mathcal{B}^{(i)}(\mathbf{u}^{(i)}; \mathbf{u}^{(i+1\dots n)}) \cdot \eta^m. \tag{3.42}$$

323 We will say that vector  $\eta^m$  is the *reference vector* of this Bethe vector. Note that, by construction,  
324  $\Psi(\mathbf{u}^{(1\dots k)} | \mathbf{u}^{(k+1\dots n)}) \in L^{(k)}$  except when  $\hat{n} = n + 1$  and  $k = n - 1$ ,  $\Psi(\mathbf{u}^{(1\dots n-1)} | \mathbf{u}^{(n)}) \in \bar{L}^{(n-1)}$ .

325 The Lemma below follows by a repeated application of Lemma 3.4.

326 **Lemma 3.7.** *Bethe vector  $\Psi(\mathbf{u}^{(1\dots k)} | \mathbf{u}^{(k+1\dots n)})$  is invariant under interchange of any two of its*  
327 *Bethe roots,  $u_i^{(l)}$  and  $u_j^{(l)}$ , for all admissible  $i, j$ , and  $l$ .*

328 The last technical result that we will need is the action of  $s_{\hat{n}\hat{n}}(v) = [S_a^{(N)}(v)]_{\hat{n}\hat{n}}$ , viz. (3.4),  
329 on a level- $n$  Bethe vector, when  $\hat{n} = n + 1$ . It is motivated by the following relation in  
330  $Y^+(\mathfrak{gl}_{2n+1})((u^{-1}, v^{-1}))$  for  $1 \leq k \leq n$ :

$$s_{\hat{n}\hat{n}}(v) s_{k\hat{n}}(u) = f^-(v, u) f^+(v, \tilde{u}) s_{k\hat{n}}(u) s_{\hat{n}\hat{n}}(v) - \left\{ \frac{p(v)}{u-v} s_{k\hat{n}}(v) \right\}^v s_{\hat{n}\hat{n}}(u).$$

331 We postpone the proof of the Lemma below to Section 4.3.

332 **Lemma 3.8.** *When  $\hat{n} = n + 1$ ,*

$$\begin{aligned}
s_{\hat{n}\hat{n}}(\nu)\Psi(\mathbf{u}^{(1\dots n)}) &= f^-(\nu, \mathbf{u}^{(n)})f^+(\nu, \tilde{\mathbf{u}}^{(n)})\mu_{\hat{n}}(\nu)\Psi(\mathbf{u}^{(1\dots n)}) \\
&+ \sum_i \frac{1}{p(u_i^{(n)})} \left\{ \frac{p(\nu)}{u_i^{(n)} - \nu} \varrho_{\check{a}_{m_n}\check{a}_{m_n}}^{(n)}(\nu) \right\}^{\nu} \mathcal{B}^{(n)}(\mathbf{u}^{(n)} \setminus u_i^{(n)}) \mathcal{R}^{(\hat{n})}(u_i^{(n)}, \mathbf{u}^{(n)} \setminus u_i^{(n)}) \\
&\times \operatorname{Res}_{w \rightarrow u_i^{(n)}} f^-(w, \mathbf{u}^{(n)})f^+(w, \tilde{\mathbf{u}}^{(n)})\mu_{\hat{n}}(w)\Psi(\mathbf{u}^{(1\dots n-1)} | \mathbf{u}_{\sigma_i}^{(n)}). \tag{3.43}
\end{aligned}$$

### 333 3.6 Transfer matrix and Bethe equations

334 We define the *transfer matrix* by

$$\tau(\nu) := \operatorname{tr}_a (M_a^{(N)} S_a^{(N)}(\nu)) = \operatorname{tr}_a (\alpha_a^{(\hat{n})} [M_a^{(N)}]^{(\hat{n})} \{p(\nu) A_a^{(\hat{n})}(\nu)\}^{\nu}) \tag{3.44}$$

335 where  $M^{(N)} = \sum_i \varepsilon_i E_{ii}^{(N)}$  with  $\varepsilon_i \in \mathbb{C}^\times$  satisfying  $\varepsilon_{N-i+1} = \varepsilon_i$  is a twist matrix, a solution to the  
336 dual twisted reflection equation

$$\begin{aligned}
(M_b^{(N)}(\nu))^{t_b} \widehat{R}_{ab}^{(N,N)}(u - \tilde{\nu}) (M_a^{(N)}(u))^{t_a} R_{ab}^{(N,N)}(\nu - u) \\
= R_{ab}^{(N,N)}(\nu - u) (M_a^{(N)}(u))^{t_a} \widehat{R}_{ab}^{(N,N)}(u - \tilde{\nu}) (M_b^{(N)}(\nu))^{t_b} \tag{3.45}
\end{aligned}$$

337 ensuring commutativity of transfer matrices, see Appendix A.2. Here  $t$  denotes the usual  
338 matrix transposition. The right hand side of (3.44) follows from the symmetry relation (2.20);  
339 the  $\alpha^{(\hat{n})}$  is a diagonal matrix with entries  $\alpha_k = 1$  for all  $k$  except  $\alpha_{\hat{n}} = 1/2$  when  $\hat{n} = n + 1$ ,  
340 which resolves the double-counting of  $s_{\hat{n}\hat{n}}(\nu)$ .

341 **Theorem 3.9.** *The Bethe vector  $\Psi(\mathbf{u}^{(1\dots n)})$  is an eigenvector of  $\tau(\nu)$  with the eigenvalue*

$$\Lambda(\nu; \mathbf{u}^{(1\dots n)}) := \sum_{k \leq \hat{n}} \alpha_k \varepsilon_k \{p(\nu) \Gamma_k(\nu; \mathbf{u}^{(1\dots n)})\}^{\nu} \tag{3.46}$$

342 where  $p(\nu)$  is given by (3.36) and

$$\Gamma_k(\nu; \mathbf{u}^{(1\dots n)}) := f^-(\nu, \mathbf{u}^{(k-1)})f^+(\nu, \mathbf{u}^{(k)})\mu_k(\nu) \quad \text{for } k < \hat{n} \tag{3.47}$$

343 and

$$\Gamma_{\hat{n}}(\nu; \mathbf{u}^{(1\dots n)}) := \begin{cases} f^-(\nu, \mathbf{u}^{(n-1)})f^+(\nu, \mathbf{u}^{(n)})f^+(\nu, \tilde{\mathbf{u}}^{(n)})\mu_n(\nu) & \text{when } \hat{n} = n, \\ f^-(\nu, \mathbf{u}^{(n)})f^+(\nu, \tilde{\mathbf{u}}^{(n)})\mu_{n+1}(\nu) & \text{when } \hat{n} = n + 1 \end{cases} \tag{3.48}$$

344 provided  $\operatorname{Res}_{\nu \rightarrow u_j^{(k)}} \Lambda(\nu; \mathbf{u}^{(1\dots n)}) = 0$  for all admissible  $k$  and  $j$ ; these equations are called Bethe  
345 equations.

346 *Proof.* When  $\hat{n} = n$ , this is a restatement of Theorems 4.3 and 4.4 in [GMR19]. We will briefly  
347 recall the main steps of the proofs therein. They will provide a backbone of the proof of the  
348 more complex  $\hat{n} = n + 1$  case.

349 *The  $\hat{n} = n$  case.* We start by noticing that

$$\prod_{i < j \leq m_n}^{\rightarrow} \check{R}_{\check{a}_j \check{a}_{j-1}}^{(\hat{n}, \hat{n})}(u_i^{(n)} - u_j^{(n)}) \check{R}_{\check{a}_j \check{a}_{j-1}}^{(\hat{n}, \hat{n})}(u_j^{(n)} - u_i^{(n)}) \Psi(\mathbf{u}^{(1\dots n-1)} | \mathbf{u}^{(n)}) = \Psi(\mathbf{u}^{(1\dots n-1)} | \mathbf{u}_{\sigma_i}^{(n)}) \tag{3.49}$$

350 where  $\mathbf{u}_{\sigma_i}^{(n)} = (u_1, \dots, u_{i-1}, u_{i+1}, \dots, u_{m_n}, u_i)$ . This identity is a consequence of Yang-Baxter  
351 moves and the identities

$$\check{R}_{\check{a}_j \check{a}_{j-1}}^{(\hat{n}, \hat{n})}(u_i^{(n)} - u_j^{(n)}) \cdot \eta^m = \eta^m, \quad \check{R}_{\check{a}_j \check{a}_{j-1}}^{(\hat{n}, \hat{n})}(u_j^{(n)} - u_i^{(n)}) \cdot \eta^m = \eta^m \quad (3.50)$$

352 which are computed using (3.20) and (3.41).

353 Next, using (3.42) and (3.44), we write

$$\tau(v) \Psi(\mathbf{u}^{(1\dots n)}) = \text{tr}_a \left( [M_a^{(N)}]^{(n)} \{p(v) A_a^{(n)}(v)\}^v \mathcal{B}^{(n)}(\mathbf{u}^{(n)}) \right) \Psi(\mathbf{u}^{(1\dots n-1)} | \mathbf{u}^{(n)}).$$

354 Lemma 3.5 allows us to exchange  $\{p(v) A_a^{(n)}(v)\}^v$  and  $\mathcal{B}^{(n)}(\mathbf{u}^{(n)})$ . Applying (3.49) to the result  
355 gives

$$\begin{aligned} \tau(v) \Psi(\mathbf{u}^{(1\dots n)}) &= \mathcal{B}^{(n)}(\mathbf{u}^{(n)}) \tau(v; \mathbf{u}^{(n)}) \Psi(\mathbf{u}^{(1\dots n-1)} | \mathbf{u}^{(n)}) \\ &+ \sum_i \frac{1}{p(u_i^{(n)})} \left\{ \frac{p(v)}{u_i^{(n)} - v} \check{t}_{\check{a}_{m_n} \check{a}_{m_n}}^{(n)}(v) \right\}^v \mathcal{B}^{(n)}(\mathbf{u}^{(n)} \setminus u_i^{(n)}) \\ &\times \mathcal{R}^{(\hat{n})}(u_i^{(n)}; \mathbf{u}^{(n)} \setminus u_i^{(n)}) \text{Res}_{w \rightarrow u_i^{(n)}} \tau(w; \mathbf{u}_{\sigma_i}^{(n)}) \Psi(\mathbf{u}^{(1\dots n-1)} | \mathbf{u}_{\sigma_i}^{(n)}) \end{aligned} \quad (3.51)$$

356 where

$$\tau(v; \mathbf{u}^{(n)}) := \text{tr}_a \left( [M_a^{(N)}]^{(n)} \{p(v) T_a^{(n)}(v; \mathbf{u}^{(n)})\}^v \right)$$

357 is a nested transfer matrix. It remains to compute the action of  $\tau(v; \mathbf{u}^{(n)})$  on the nested Bethe  
358 vector  $\Psi(\mathbf{u}^{(1\dots n-1)} | \mathbf{u}^{(n)}) \in L^{(n-1)}$ . By Lemma 3.19, this can be achieved using  $Y(\mathfrak{gl}_n)$ -type  
359 nested Bethe ansatz techniques assisted by Lemmas 3.6 and 3.7 leading to the eigenvalue  
360 (3.46) and the corresponding Bethe equations.

361 *The  $\hat{n} = n + 1$  case.* In this case we can not apply Lemma 3.5 directly since this would lead to  
362 the following nested transfer matrix

$$\begin{aligned} \tau(v; \mathbf{u}^{(n)}) &= \text{tr}_a \left( \alpha_a^{(\hat{n})} [M_a^{(N)}]^{(\hat{n})} \{p(v) T_a^{(\hat{n})}(v; \mathbf{u}^{(n)})\}^v \right) \\ &= \text{tr}_a \left( [M_a^{(N)}]^{(n)} \{p(v) [T_a^{(\hat{n})}(v; \mathbf{u}^{(n)})]^{(n)}\}^v \right) + \frac{1}{2} \varepsilon_{\hat{n}} \{p(v) [T_a^{(\hat{n})}(v; \mathbf{u}^{(n)})]_{\hat{n}\hat{n}}\}^v. \end{aligned}$$

363 However, the space  $\bar{L}^{(n-1)}$  is not stable under the action of  $[T_a^{(\hat{n})}(v; \mathbf{u}^{(n)})]_{\hat{n}\hat{n}}$ . This is because  
364  $[T_a^{(\hat{n})}(v; \mathbf{u}^{(n)})]_{\hat{n}\hat{n}}$  has operators  $[\widehat{R}_{\check{a}_i \check{a}}^{(\hat{n}, \hat{n})}(u_i^{(n)} - v)]_{\hat{n}j}$  with  $j \leq n$  that map  $E_{\hat{n}-j+1}^{(\hat{n})} \in \bar{V}_{\check{a}_i}^{\hat{n}}$  to  $E_1^{(\hat{n})}$ .  
365 Therefore, the right hand side of (3.37) would no longer represent a splitting into “wanted”  
366 and “unwanted” terms. A resolution of this issue is to single-out the operator  $s_{\hat{n}\hat{n}}(v)$  from the  
367 very beginning. From (2.11) we know that  $s_{\hat{n}\hat{n}}(\tilde{u}) = s_{\hat{n}\hat{n}}(u)$  giving  $\{p(v) s_{\hat{n}\hat{n}}(v)\}^v = 2s_{\hat{n}\hat{n}}(v)$ .  
368 This allows us to rewrite the transfer matrix as

$$\tau(v) = \text{tr}_a \left( [M_a^{(N)}]^{(n)} \{p(v) [A_a^{(\hat{n})}(v)]^{(n)}\}^v \right) + \varepsilon_{\hat{n}} s_{\hat{n}\hat{n}}(v). \quad (3.52)$$

369 We can now use Lemma 3.5 to exchange  $\{p(v) [A_a^{(\hat{n})}(v)]^{(n)}\}^v$  and  $\mathcal{B}^{(n)}(\mathbf{u}^{(n)})$ , and Lemma 3.8  
370 to compute the action of  $s_{\hat{n}\hat{n}}(v)$  on  $\Psi(\mathbf{u}^{(1\dots n)})$ . This gives an expressions equivalent to (3.51)  
371 except the nested transfer matrix is now given by

$$\tau(v; \mathbf{u}^{(n)}) := \text{tr}_a \left( [M_a^{(N)}]^{(n)} \{p(v) \overline{T_a^{(n)}}(v; \mathbf{u}^{(n)})\}^v \right) + \varepsilon_{\hat{n}} f^-(v, \mathbf{u}^{(n)}) f^+(v, \tilde{\mathbf{u}}^{(n)}) \mu_{\hat{n}}(v).$$

372 Here we invoked Lemma 3.1 to replace  $[T_a^{(n)}(v; \mathbf{u}^{(n)})]^{(n)}$  with  $\overline{T_a^{(n)}}(v; \mathbf{u}^{(n)})$ . The remaining  
373 steps are the same as in the  $\hat{n} = n$  case.  $\square$

374 *Remark 3.10.* Let  $(a_{ij})_{i,j=1}^n$  denote Cartan matrix of type  $A_n$ . Let  $(b_{ij})_{i,j=1}^n$  denote a zero matrix  
 375 when  $\hat{n} = n + 1$  and let  $b_{nn} = 2$ ,  $b_{n-1,n} = b_{n,n-1} = -1$ , and  $b_{ij} = 0$  otherwise, when  $\hat{n} = n$ . Set  
 376  $m_0 := 0$  and  $z_j^{(k)} := u_j^{(k)} - \frac{1}{2}(k - \rho)$ . Then Bethe equations can be written as, for each  $k < n$ ,

$$\prod_{l=k-1}^{k+1} \prod_{i=1}^{m_l} \frac{z_j^{(k)} - z_i^{(l)} + \frac{1}{2}a_{kl}}{z_j^{(k)} - z_i^{(l)} - \frac{1}{2}a_{kl}} \cdot \frac{z_j^{(k)} + z_i^{(l)} + n + \frac{1}{2}b_{kl}}{z_j^{(k)} + z_i^{(l)} + n - \frac{1}{2}b_{kl}} = -\frac{\varepsilon_{k+1}}{\varepsilon_k} \cdot \frac{\mu_{k+1}(u_j^{(k)})}{\mu_k(u_j^{(k)})}, \quad (3.53)$$

$$\frac{z_j^{(n)} + \frac{1}{2}(n+1)}{z_j^{(n)} + \frac{1}{2}(\hat{n}-1)} \prod_{l=n-1}^n \prod_{i=1}^{m_l} \frac{z_j^{(n)} - z_i^{(l)} + \frac{1}{2}a_{nl}}{z_j^{(n)} - z_i^{(l)} - \frac{1}{2}a_{nl}} \prod_{l=\hat{n}-1}^n \prod_{i=1}^{m_l} \frac{z_j^{(n)} + z_i^{(l)} + n + \frac{1}{2}b_{nl}}{z_j^{(n)} + z_i^{(l)} + \hat{n} - \frac{1}{2}b_{nl}} = -\frac{\varepsilon_{\hat{n}}}{\varepsilon_n} \cdot \frac{\mu_{\hat{n}}(\tilde{u}_j^{(n)})}{\mu_n(u_j^{(n)})}. \quad (3.54)$$

### 377 3.7 Trace formula

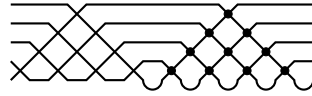
378 Define the “master” creation operator

$$\begin{aligned} \mathcal{B}_N(\mathbf{u}^{(1\dots n)}) := & \prod_{k \leq n} \prod_{j < i} \frac{1}{f+(u_j^{(k)}, u_i^{(k)}) (f+(u_j^{(k)}, \tilde{u}_i^{(k)})) \delta_{\hat{n},n}} \\ & \times \text{tr} \left[ \prod_{(k,i) \succ (l,j)} R_{a_i^k a_j^l}^{(N,N)}(u_i^{(k)} - u_j^{(l)}) \prod_{(k,i)} \left( S_{a_i^k}^{(N)}(u_i^{(k)}) \prod_{(k,i) \succ (l,j)} \widehat{R}_{a_i^k a_j^l}^{(N,N)}(\tilde{u}_i^{(k)} - u_j^{(l)}) \right) \right. \\ & \left. \times (E_{n+1,n}^{(N)})^{\otimes m_n} \otimes \dots \otimes (E_{21}^{(N)})^{\otimes m_1} \right] \end{aligned} \quad (3.55)$$

379 where  $(k, i) \succ (l, j)$  means that  $k > l$  or  $k = l$  and  $i > j$ , and the products over tuples are  
 380 defined in terms of the following rule

$$\prod_{(k,i)} = \overleftarrow{\prod}_{k < n} \overleftarrow{\prod}_{i < m_k}$$

381 In other words, these products are ordered in the reversed lexicographical order. The trace is  
 382 taken over all  $a_i^k$  spaces, including  $a_i^n$ , which are associated with level- $n$  excitations. Note that  
 383  $(k, i)$  is fixed in the third product inside the trace. Diagrammatically, the operator inside the  
 384 trace is of the form



385 where  $\times = R_{a_i^k a_j^l}(u_i^{(k)} - u_j^{(l)})$ ,  $\times = \widehat{R}_{a_i^k a_j^l}(\tilde{u}_i^{(k)} - u_j^{(l)})$ , and  $\cup = S_{a_i^k}(u_i^{(k)})$ .

386 *Example 3.11.* The “master” creation operators of low rank:

$$\begin{aligned} \mathcal{B}_3(u_1^{(1)}) &= s_{12}(u_1^{(1)}), & \mathcal{B}_3(u_1^{(1)}, u_2^{(1)}) &= s_{12}(u_2^{(1)})s_{12}(u_1^{(1)}) + \frac{s_{13}(u_2^{(1)})s_{22}(u_1^{(1)})}{u_1^{(1)} - \tilde{u}_2^{(1)}}, \\ \mathcal{B}_4(u_1^{(1)}, u_1^{(2)}) &= s_{23}(u_1^{(2)})s_{12}(u_1^{(1)}) + \frac{s_{24}(u_1^{(2)})s_{22}(u_1^{(1)})}{u_1^{(1)} - \tilde{u}_1^{(2)}} + \frac{(u_1^{(1)} - \tilde{u}_1^{(2)} + 1)s_{13}(u_1^{(2)})s_{22}(u_1^{(1)})}{(u_1^{(1)} - u_1^{(2)})(u_1^{(1)} - \tilde{u}_1^{(2)}), \\ \mathcal{B}_5(u_1^{(1)}, u_1^{(2)}) &= s_{23}(u_1^{(2)})s_{12}(u_1^{(1)}) + \frac{s_{13}(u_1^{(2)})s_{22}(u_1^{(1)})}{u_1^{(1)} - u_1^{(2)}} + \frac{s_{25}(u_1^{(2)})s_{32}(u_1^{(1)})}{u_1^{(1)} - \tilde{u}_1^{(2)}} \\ &+ \frac{s_{14}(u_1^{(2)})s_{32}(u_1^{(1)})}{(u_1^{(1)} - u_1^{(2)})(u_1^{(1)} - \tilde{u}_1^{(2)})}. \end{aligned}$$



387 **Proposition 3.12.** *The level- $n$  Bethe vector (3.42) can be written as*

$$\Psi(\mathbf{u}^{(1..n)}) = \mathcal{B}_N(\mathbf{u}^{(1..n)}) \cdot \eta. \quad (3.56)$$

388 *Proof.* First, notice that  $R$ -matrices  $R_{a_i^k a_j^k}^{(N,N)}(u_i^{(k)} - u_j^{(k)})$  in (3.55) evaluate to  $f^+(u_j^{(k)} - u_i^{(k)})$  un-  
389 der the trace. This cancels the first overall factor in (3.55). The second overall factor is the  
390 choice of normalisation in (3.26). Next, let  $V_a^{(N)}$  and  $V_b^{(N)}$  denote copies of  $\mathbb{C}^N$ . Then, for any  
391  $\zeta \in (L^{(n)})^0$  and  $E_i^{(N)} \otimes E_j^{(N)} \in V_a^{(N)} \otimes V_b^{(N)}$  with  $1 \leq i, j \leq n$ , we have

$$Q_{ab}^{(N,N)} E_i^{(N)} \otimes E_j^{(N)} = 0$$

392 and

$$Q_{ab}^{(N,N)} S_a^{(N)}(v) \cdot E_i^{(N)} \otimes E_j^{(N)} \otimes \zeta = \sum_k Q_{ab}^{(N,N)} \cdot E_k^{(N)} \otimes E_j^{(N)} \otimes s_{ki}(v) \zeta = 0.$$

393 Thus  $\widehat{R}_{a_i^k a_j^l}^{(N,N)}(u_i^{(k)} - u_j^{(l)})$  with  $1 \leq k, l < n$  act as identity operators in (3.56). This gives an  
394 expression analogous (up to Yang-Baxter moves) to that in Proposition 4.7 of [GMR19]. The  
395  $N = 2n$  case then follows from that proposition. The  $N = 2n+1$  case is proven analogously.  $\square$

## 396 4 Recurrence relations

### 397 4.1 Notation

398 Given any tuple  $\mathbf{u}$  of complex parameters, let  $(\mathbf{u}_I, \mathbf{u}_{II}) \vdash \mathbf{u}$  be a partition of this tuple and let  
399  $\mathbf{u}_{I,II} := \mathbf{u}_I \cup \mathbf{u}_{II} = \mathbf{u}$ . Assume that  $1 \leq k < |\mathbf{u}|$  and set

$$\sum_{|\mathbf{u}_{II}|=k} f(\mathbf{u}_I) := \sum_{i_1 < i_2 < \dots < i_k} f(\mathbf{u} \setminus (u_{i_1}, u_{i_2}, \dots, u_{i_k}))$$

400 for any function or operator  $f$ . We will use a natural generalisation of this notation for any  
401 partition of  $\mathbf{u}$ . For instance, for  $(\mathbf{u}_I, \mathbf{u}_{II}, \mathbf{u}_{III}) \vdash \mathbf{u}$  we have  $\mathbf{u}_{I,II} = \mathbf{u}_I \cup \mathbf{u}_{II}$ ,  $\mathbf{u}_{II,III} = \mathbf{u}_{II} \cup \mathbf{u}_{III}$ , etc.,  
402 and e.g.

$$\sum_{|\mathbf{u}_{III}|=1} \sum_{|\mathbf{u}_{II}|=2} f(\mathbf{u}_{II}) g(\mathbf{u}_I) = \sum_j \sum_{\substack{i_1 < i_2 \\ i_1 \neq j, i_2 \neq j}} f((u_{i_1}, u_{i_2})) g(\mathbf{u} \setminus (u_{i_1}, u_{i_2}, u_j)).$$

403 We extend the notation above to partitions of tuples  $\mathbf{u}^{(1..n)}$  by allowing empty partitions.  
404 The empty partitions will be the ones that are missing from the expressions. For instance, an  
405 expression of the form

$$\sum_{\substack{|\mathbf{u}_{II}^{(r)}|=k \\ i < r \leq n}} f(\mathbf{u}_{II}^{(r)}) g(\mathbf{u}_I^{(1..n)})$$

406 will mean that  $\mathbf{u}_{II}^{(1)} = \dots = \mathbf{u}_{II}^{(i)} = \emptyset$  so that  $\mathbf{u}_I^{(1..n)} = (\mathbf{u}^{(1)}, \dots, \mathbf{u}^{(i)}, \mathbf{u}_I^{(i+1)}, \dots, \mathbf{u}_I^{(n)})$ . We will  
407 also use the notation  $|\mathbf{u}_{III}^{(r)}| = 0$  meaning  $\mathbf{u}_{III}^{(r)} = \emptyset$ .

408 The notation  $|\mathbf{u}_{II,III}^{(r)}| = (k, l)$  will mean that  $|\mathbf{u}_{II}^{(r)}| = k$  and  $|\mathbf{u}_{III}^{(r)}| = l$  and the notation  
409  $|\mathbf{u}_{II}^{(r,s)}| = (k, l)$  will mean that  $|\mathbf{u}_{II}^{(r)}| = k$  and  $|\mathbf{u}_{II}^{(s)}| = l$  so that

$$\sum_{|\mathbf{u}_{II,III}^{(r)}|=(k,l)} = \sum_{|\mathbf{u}_{III}^{(r)}|=l} \sum_{|\mathbf{u}_{II}^{(r)}|=k} \quad \text{and} \quad \sum_{|\mathbf{u}_{II}^{(r,s)}|=(k,l)} = \sum_{|\mathbf{u}_{II}^{(s)}|=l} \sum_{|\mathbf{u}_{II}^{(r)}|=k}.$$

410 A notation of the form  $\mathbf{u}_{II,III}^{(r,s)}$  will not be used.

## 4.2 Recurrence relations

We will combine the composite model method and the known  $Y(\mathfrak{gl}_n)$ -type recurrence relations to obtain recurrence relations for  $Y^\pm(\mathfrak{gl}_N)$ -based Bethe vectors. The composite model method was introduced in [IK84]. For a pedagogical review, see [Sla20]. Recurrence relations for  $Y(\mathfrak{gl}_n)$ -based Bethe vectors were obtained in [HL<sup>+</sup>17b]. We will need the following statement which follows directly from those in [HL<sup>+</sup>17b] recalled in Appendix A.3. Recall the notation (3.9) of rational functions.

**Proposition 4.1.** Consider a  $Y(\mathfrak{gl}_n)$ -based Bethe vector  $\Phi(\mathbf{v}^{(1\dots n-1)} | \mathbf{v}^{(n)})$  in the quantum space

$$V_{a_{m_n}}^{(n)} \otimes \dots \otimes V_{a_1}^{(n)} \otimes L(\boldsymbol{\lambda}) \quad (4.1)$$

with  $V_{a_i}^{(n)} \cong \mathbb{C}^n$ , a finite-dimensional irreducible  $Y(\mathfrak{gl}_n)$ -module  $L(\boldsymbol{\lambda})$ , Bethe roots  $\mathbf{v}^{(1\dots n-1)}$  and inhomogeneities  $\mathbf{v}^{(n)}$  associated with spaces  $V_{a_i}^{(n)}$ . Set

$$\Lambda_k(\mathbf{z}; \mathbf{v}^{(1\dots n-1)}) := f^-(\mathbf{z}, \mathbf{v}^{(k-1)}) f^+(\mathbf{z}, \mathbf{v}^{(k)}) \lambda_k(\mathbf{z}). \quad (4.2)$$

An expansion of  $\Phi(\mathbf{v}^{(1\dots n-1)} | \mathbf{v}^{(n)})$  in the space  $V_{a_{m_n}}^{(n)}$  is given by

$$\Phi(\mathbf{v}^{(1\dots n-1)} | \mathbf{v}^{(n)}) = \sum_{1 \leq i \leq n} \sum_{\substack{|\mathbf{v}_{\text{II}}^{(r)}| = 1 \\ i \leq r < n}} \prod_{i < k \leq n} \frac{\Lambda_k(\mathbf{v}_{\text{II}}^{(k-1)}; \mathbf{v}_I^{(1\dots n)})}{\mathbf{v}_{\text{II}}^{(k-1)} - \mathbf{v}_{\text{II}}^{(k)}} E_i^{(n)} \otimes \Phi(\mathbf{v}_I^{(1\dots n-1)} | \mathbf{v}_I^{(n)}) \quad (4.3)$$

where  $\mathbf{v}_{\text{II}}^{(n)} = \mathbf{v}_{m_n}^{(n)}$  and  $\mathbf{v}_{\text{II}}^{(r)} = \emptyset$  for all  $1 \leq r < i$  so that  $\mathbf{v}_I^{(1\dots n)} = (\mathbf{v}^{(1)}, \dots, \mathbf{v}^{(i-1)}, \mathbf{v}_I^{(i)}, \dots, \mathbf{v}^{(n)})$ .

**Corollary 4.2.** An expansion of Bethe vector  $\Phi(\mathbf{v}^{(1\dots n-1)} | \mathbf{v}^{(n)})$  in the space  $V_{a_{m_n}}^{(n)} \otimes V_{a_{m_{n-1}}}^{(n)}$  is given by

$$\begin{aligned} & \sum_{1 \leq i \leq n} \sum_{\substack{|\mathbf{v}_{\text{II,III}}^{(r)}| = (2,0) \\ i \leq r < n}} \prod_{i < k \leq n} K(\mathbf{v}_{\text{II}}^{(k-1)} | \mathbf{v}_{\text{II,III}}^{(k)}) \Lambda_k(\mathbf{v}_{\text{II}}^{(k-1)}; \mathbf{v}_I^{(1\dots n)}) E_i^{(n)} \otimes E_i^{(n)} \otimes \Phi(\mathbf{v}_I^{(1\dots n-1)} | \mathbf{v}_I^{(n)}) \\ & + \sum_{1 \leq i < j \leq n} \sum_{\substack{|\mathbf{v}_{\text{III}}^{(r)}| = 1 \\ i \leq r < n}} \sum_{\substack{|\mathbf{v}_{\text{II}}^{(s)}| = 1 \\ j \leq s < n}} \prod_{i < k < j} \frac{\Lambda_k(\mathbf{v}_{\text{III}}^{(k-1)}; \mathbf{v}_I^{(1\dots n-1)})}{\mathbf{v}_{\text{III}}^{(k-1)} - \mathbf{v}_{\text{III}}^{(k)}} \cdot \Lambda_j(\mathbf{v}_{\text{III}}^{(j-1)}; \mathbf{v}_I^{(1\dots n-1)}) \\ & \times \prod_{j < k \leq n} \frac{\Lambda_k(\mathbf{v}_{\text{II}}^{(k-1)}; \mathbf{v}_I^{(1\dots n-1)}) \Lambda_k(\mathbf{v}_{\text{III}}^{(k-1)}; \mathbf{v}_{\text{I,II}}^{(1\dots n-1)})}{(\mathbf{v}_{\text{II}}^{(k-1)} - \mathbf{v}_{\text{II}}^{(k)})(\mathbf{v}_{\text{III}}^{(k-1)} - \mathbf{v}_{\text{III}}^{(k)})} \\ & \times \left( \frac{f^+(\mathbf{u}_{\text{III}}^{(j-1)}, \mathbf{v}_{\text{II}}^{(j)})}{\mathbf{v}_{\text{III}}^{(j-1)} - \mathbf{v}_{\text{III}}^{(j)}} E_i^{(n)} \otimes E_j^{(n)} + \frac{1}{\mathbf{v}_{\text{III}}^{(j-1)} - \mathbf{v}_{\text{II}}^{(j)}} E_j^{(n)} \otimes E_i^{(n)} \right) \otimes \Phi(\mathbf{v}_I^{(1\dots n-1)} | \mathbf{v}_I^{(n)}) \quad (4.4) \end{aligned}$$

where  $\mathbf{v}_{\text{III}}^{(n)} = \mathbf{v}_{m_n}^{(n)}$ ,  $\mathbf{v}_{\text{II}}^{(n)} = \mathbf{v}_{m_{n-1}}^{(n)}$  and  $\mathbf{v}_{\text{III}}^{(r)} = \mathbf{v}_{\text{II}}^{(r)} = \emptyset$  for all  $1 \leq r < i$  in the first sum and  $\mathbf{v}_{\text{III}}^{(r)} = \mathbf{v}_{\text{II}}^{(s)} = \emptyset$  for all  $1 \leq r < i$  and  $1 \leq s < j$  in the second sum, and

$$K(\mathbf{u} | \mathbf{v}) := \frac{\prod_{i,j} (u_i - v_j + 1)}{\prod_{i < j} (u_i - u_j)(v_j - v_i)} \det_{i,j} \left( \frac{1}{(u_i - v_j)(u_i - v_j + 1)} \right) \quad (4.5)$$

is the domain wall boundary partition function.

427 *Proof.* Applying (4.3) to  $\Phi(\mathbf{v}^{(1\dots n-1)} | \mathbf{v}^{(n)})$  twice gives

$$\sum_{\substack{1 \leq i, j \leq n \\ i \leq r < n \\ j \leq s < n}} \sum_{\substack{|\mathbf{v}_{\text{III}}^{(r)}|=1 \\ |\mathbf{v}_{\text{II}}^{(s)}|=1}} \prod_{i < k \leq n} \frac{\Lambda_k(\mathbf{v}_{\text{III}}^{(k-1)}; \mathbf{v}_{\text{I,II}}^{(1\dots n)})}{\mathbf{v}_{\text{III}}^{(k-1)} - \mathbf{v}_{\text{III}}^{(k)}} \prod_{j < l \leq n} \frac{\Lambda_l(\mathbf{v}_{\text{II}}^{(l-1)}; \mathbf{v}_{\text{I}}^{(1\dots n)})}{\mathbf{v}_{\text{II}}^{(l-1)} - \mathbf{v}_{\text{II}}^{(l)}} \Phi_{\bar{i}\bar{j}} \quad (4.6)$$

428 where  $\mathbf{v}_{\text{III}}^{(r)} = \mathbf{v}_{\text{II}}^{(s)} = \emptyset$  for all  $1 \leq r < i$  and  $1 \leq s < j$ , and  $\Phi_{ij} := E_i^{(n)} \otimes E_i^{(n)} \otimes \Phi(\mathbf{v}_{\text{I}}^{(1\dots n-1)} | \mathbf{v}_{\text{I}}^{(n)})$ .

429 Cases  $i = j$ . Notice that

$$\Lambda_k(\mathbf{v}_{\text{III}}^{(k-1)}; \mathbf{v}_{\text{I,II}}^{(1\dots n)}) = f^-(\mathbf{v}_{\text{III}}^{(k-1)}, \mathbf{v}_{\text{II}}^{(k-1)}) f^+(\mathbf{v}_{\text{III}}^{(k-1)}, \mathbf{v}_{\text{II}}^{(k)}) \Lambda_k(\mathbf{v}_{\text{III}}^{(k-1)}; \mathbf{v}_{\text{I}}^{(1\dots n)})$$

430 and

$$\frac{f^-(\mathbf{v}_{\text{III}}^{(k-1)}, \mathbf{v}_{\text{II}}^{(k-1)}) f^+(\mathbf{v}_{\text{III}}^{(k-1)}, \mathbf{v}_{\text{II}}^{(k)})}{(\mathbf{v}_{\text{III}}^{(k-1)} - \mathbf{v}_{\text{III}}^{(k)})(\mathbf{v}_{\text{II}}^{(k-1)} - \mathbf{v}_{\text{II}}^{(k)})} + \frac{f^-(\mathbf{v}_{\text{II}}^{(k-1)}, \mathbf{v}_{\text{III}}^{(k-1)}) f^+(\mathbf{v}_{\text{II}}^{(k-1)}, \mathbf{v}_{\text{II}}^{(k)})}{(\mathbf{v}_{\text{II}}^{(k-1)} - \mathbf{v}_{\text{III}}^{(k)})(\mathbf{v}_{\text{III}}^{(k-1)} - \mathbf{v}_{\text{II}}^{(k)})} = K(\mathbf{v}_{\text{II,III}}^{(k-1)} | \mathbf{v}_{\text{II,III}}^{(k)}).$$

431 These identities allow us to rewrite the  $i = j$  cases of (4.6) as

$$\sum_{1 \leq i \leq n} \sum_{\substack{|\mathbf{v}_{\text{II,III}}^{(r)}|=(1,1) \\ i \leq r < n}} \prod_{i < k \leq n} K(\mathbf{v}_{\text{II,III}}^{(k-1)} | \mathbf{v}_{\text{II,III}}^{(k)}) \Lambda_k(\mathbf{v}_{\text{II,III}}^{(k-1)}; \mathbf{v}_{\text{I}}^{(1\dots n)}) \Phi_{\bar{i}\bar{i}}$$

432 giving the first sum in (4.4).

433 Cases  $i < j$ . Since  $\mathbf{v}_{\text{II}}^{(s)} = \emptyset$  for  $s < j$  in (4.6) we have

$$\Lambda_k(\mathbf{v}_{\text{III}}^{(k-1)}; \mathbf{v}_{\text{I,II}}^{(1\dots n)}) = \Lambda_k(\mathbf{v}_{\text{III}}^{(k-1)}; \mathbf{v}_{\text{I}}^{(1\dots n)}) \quad \text{for } k < j$$

434 and

$$\Lambda_j(\mathbf{v}_{\text{III}}^{(j-1)}; \mathbf{v}_{\text{I,II}}^{(1\dots n)}) = f^+(\mathbf{v}_{\text{III}}^{(j-1)}, \mathbf{v}_{\text{II}}^{(j)}) \Lambda_j(\mathbf{v}_{\text{III}}^{(j-1)}; \mathbf{v}_{\text{I}}^{(1\dots n)})$$

435 allowing us to rewrite the  $i < j$  cases as

$$\begin{aligned} & \sum_{1 \leq i < j \leq n} \sum_{\substack{|\mathbf{v}_{\text{III}}^{(r)}|=1 \\ i \leq r < n}} \sum_{\substack{|\mathbf{v}_{\text{II}}^{(s)}|=1 \\ j \leq s < n}} \prod_{i < k < j} \frac{\Lambda_k(\mathbf{v}_{\text{III}}^{(k-1)}; \mathbf{v}_{\text{I}}^{(1\dots n)})}{\mathbf{v}_{\text{III}}^{(k-1)} - \mathbf{v}_{\text{III}}^{(k)}} \cdot \Lambda_j(\mathbf{v}_{\text{III}}^{(j-1)}; \mathbf{v}_{\text{I}}^{(1\dots n)}) \\ & \times \prod_{j < k \leq n} \frac{\Lambda_k(\mathbf{v}_{\text{II}}^{(k-1)}; \mathbf{v}_{\text{I}}^{(1\dots n)}) \Lambda_k(\mathbf{v}_{\text{III}}^{(k-1)}; \mathbf{v}_{\text{I,II}}^{(1\dots n)})}{(\mathbf{v}_{\text{II}}^{(k-1)} - \mathbf{v}_{\text{II}}^{(k)})(\mathbf{v}_{\text{III}}^{(k-1)} - \mathbf{v}_{\text{III}}^{(k)})} \cdot \frac{f^+(\mathbf{v}_{\text{III}}^{(j-1)}, \mathbf{v}_{\text{II}}^{(j)})}{\mathbf{v}_{\text{III}}^{(j-1)} - \mathbf{v}_{\text{III}}^{(j)}} \Phi_{\bar{i}\bar{j}}. \end{aligned} \quad (4.7)$$

436 Cases  $i > j$ . Interchanging indices  $i$  and  $j$  in (4.6) gives

$$\begin{aligned} & \sum_{1 \leq i < j \leq n} \sum_{\substack{|\mathbf{v}_{\text{III}}^{(s)}|=1 \\ j \leq s < n}} \sum_{\substack{|\mathbf{v}_{\text{II}}^{(r)}|=1 \\ i \leq r < n}} \prod_{i < k < j} \frac{\Lambda_k(\mathbf{v}_{\text{II}}^{(k-1)}; \mathbf{v}_{\text{I}}^{(1\dots n)})}{\mathbf{v}_{\text{II}}^{(k-1)} - \mathbf{v}_{\text{II}}^{(k)}} \cdot \Lambda_j(\mathbf{v}_{\text{II}}^{(j-1)}; \mathbf{v}_{\text{I}}^{(1\dots n)}) \\ & \times \prod_{j < k \leq n} \frac{\Lambda_k(\mathbf{v}_{\text{II}}^{(k-1)}; \mathbf{v}_{\text{I}}^{(1\dots n)}) \Lambda_k(\mathbf{v}_{\text{III}}^{(k-1)}; \mathbf{v}_{\text{I,II}}^{(1\dots n)})}{(\mathbf{v}_{\text{II}}^{(k-1)} - \mathbf{v}_{\text{II}}^{(k)})(\mathbf{v}_{\text{III}}^{(k-1)} - \mathbf{v}_{\text{III}}^{(k)})} \cdot \frac{1}{\mathbf{v}_{\text{II}}^{(j-1)} - \mathbf{v}_{\text{II}}^{(j)}} \Phi_{\bar{j}\bar{i}}. \end{aligned} \quad (4.8)$$

437 Since  $i < j$  we can rename  $\mathbf{v}_{\text{II}}^{(r)}$  by  $\mathbf{v}_{\text{III}}^{(r)}$  for  $i \leq r < j$  and combine the result with (4.7). This  
438 gives the second sum in (4.4).  $\square$

439 *Example 4.3.* When  $N = 3$ , expansion (4.4) of  $\Phi(\mathbf{v}^{(1,2)} | \mathbf{v}^{(3)})$  is

$$\begin{aligned}
 & \Phi_{11} + \sum_{|\mathbf{v}_{\text{II}}^{(2)}|=2} K(\mathbf{v}_{\text{II}}^{(2)} | \mathbf{v}_{\text{II,III}}^{(3)}) \Lambda_3(\mathbf{v}_{\text{II}}^{(2)}; \mathbf{v}_{\text{I}}^{(1,2,3)}) \Phi_{22} \\
 & + \sum_{|\mathbf{v}_{\text{II}}^{(1,2)}|=(2,2)} K(\mathbf{v}_{\text{II}}^{(1)} | \mathbf{v}_{\text{II}}^{(2)}) K(\mathbf{v}_{\text{II}}^{(2)} | \mathbf{v}_{\text{II,III}}^{(3)}) \Lambda_2(\mathbf{v}_{\text{II}}^{(1)}; \mathbf{v}_{\text{I}}^{(1,2,3)}) \Lambda_3(\mathbf{v}_{\text{II}}^{(2)}; \mathbf{v}_{\text{I}}^{(1,2,3)}) \Phi_{33} \\
 & + \sum_{|\mathbf{v}_{\text{III}}^{(2)}|=1} \Lambda_3(\mathbf{v}_{\text{III}}^{(2)}; \mathbf{v}_{\text{I}}^{(1,2,3)}) \left( \frac{f^+(\mathbf{v}_{\text{III}}^{(2)}, \mathbf{v}_{\text{II}}^{(3)})}{\mathbf{v}_{\text{III}}^{(2)} - \mathbf{v}_{\text{III}}^{(3)}} \Phi_{21} + \frac{1}{\mathbf{v}_{\text{III}}^{(2)} - \mathbf{v}_{\text{II}}^{(3)}} \Phi_{12} \right) \\
 & + \sum_{|\mathbf{v}_{\text{III}}^{(1,2)}|=(1,1)} \frac{\Lambda_2(\mathbf{v}_{\text{III}}^{(1)}; \mathbf{v}_{\text{I}}^{(1,2,3)})}{\mathbf{v}_{\text{III}}^{(1)} - \mathbf{v}_{\text{III}}^{(2)}} \Lambda_3(\mathbf{v}_{\text{III}}^{(2)}; \mathbf{v}_{\text{I}}^{(1,2,3)}) \left( \frac{f^+(\mathbf{v}_{\text{III}}^{(2)}, \mathbf{v}_{\text{II}}^{(3)})}{\mathbf{v}_{\text{III}}^{(2)} - \mathbf{v}_{\text{III}}^{(3)}} \Phi_{31} + \frac{1}{\mathbf{v}_{\text{III}}^{(2)} - \mathbf{v}_{\text{II}}^{(3)}} \Phi_{13} \right) \\
 & + \sum_{|\mathbf{v}_{\text{III}}^{(1,2)}|=(1,1)} \sum_{|\mathbf{v}_{\text{II}}^{(2)}|=1} \Lambda_2(\mathbf{v}_{\text{III}}^{(1)}; \mathbf{v}_{\text{I}}^{(1,2,3)}) \frac{\Lambda_3(\mathbf{v}_{\text{II}}^{(2)}; \mathbf{v}_{\text{I}}^{(1,2,3)}) \Lambda_3(\mathbf{v}_{\text{III}}^{(2)}; \mathbf{v}_{\text{I,II}}^{(1,2,3)})}{(\mathbf{v}_{\text{II}}^{(2)} - \mathbf{v}_{\text{II}}^{(3)})(\mathbf{v}_{\text{III}}^{(2)} - \mathbf{v}_{\text{III}}^{(3)})} \\
 & \quad \times \left( \frac{f^+(\mathbf{v}_{\text{III}}^{(1)}, \mathbf{v}_{\text{II}}^{(2)})}{\mathbf{v}_{\text{III}}^{(1)} - \mathbf{v}_{\text{III}}^{(2)}} \Phi_{32} + \frac{1}{\mathbf{v}_{\text{III}}^{(1)} - \mathbf{v}_{\text{II}}^{(2)}} \Phi_{23} \right)
 \end{aligned}$$

440 where  $\mathbf{v}_{\text{III}}^{(3)} = \mathbf{v}_{m_3}^{(3)}$ ,  $\mathbf{v}_{\text{II}}^{(3)} = \mathbf{v}_{m_3-1}^{(3)}$ , and  $\mathbf{v}_{\text{III}}^{(2)} = \mathbf{v}_{\text{III}}^{(1)} = \mathbf{v}_{\text{II}}^{(1)} = \emptyset$  in the first sum,  $\mathbf{v}_{\text{III}}^{(2)} = \mathbf{v}_{\text{III}}^{(1)} = \emptyset$  in  
 441 the second sum,  $\mathbf{v}_{\text{III}}^{(1)} = \mathbf{v}_{\text{II}}^{(2)} = \mathbf{v}_{\text{II}}^{(1)} = \emptyset$  in the third sum,  $\mathbf{v}_{\text{II}}^{(2)} = \mathbf{v}_{\text{II}}^{(1)} = \emptyset$  in the fourth sum and  
 442  $\mathbf{v}_{\text{II}}^{(1)} = \emptyset$  in the last sum, and  $\Phi_{ij} = E_i^{(3)} \otimes E_j^{(3)} \otimes \Phi(\mathbf{v}_{\text{I}}^{(1,2)} | \mathbf{v}_{\text{I}}^{(3)})$ .

443 We are ready to state the main results of this section, recurrence relations for twisted  
 444 Yangian based Bethe vectors. The even case follows almost immediately from Corollary 4.2.  
 445 The odd case will require additional steps which are due to the  $E_1^{(n)} = E_2^{(n+1)}$  factors in the  
 446 reference vector  $\eta^m$ .

447 **Proposition 4.4.**  $Y^\pm(\mathfrak{gl}_{2n})$ -based Bethe vectors satisfy the recurrence relation

$$\begin{aligned}
 \Psi(\mathbf{u}^{(1\dots n)}) &= \sum_{1 \leq i \leq n} \sum_{\substack{|\mathbf{u}_{\text{II,III}}^{(r)}|=(2,0) \\ i \leq r < n}} \prod_{i < k \leq n} K(\mathbf{u}_{\text{II}}^{(k-1)} | \mathbf{u}_{\text{II,III}}^{(k)}) \Gamma_k(\mathbf{u}_{\text{II}}^{(k-1)}; \mathbf{u}_{\text{I}}^{(1\dots n)}) s_{i,2n-i+1}(\mathbf{u}_{\text{III}}^{(n)}) \Psi(\mathbf{u}_{\text{I}}^{(1\dots n)}) \\
 & + \sum_{1 \leq i < j \leq n} \sum_{\substack{|\mathbf{u}_{\text{III}}^{(r)}|=1 \\ i \leq r < n}} \sum_{\substack{|\mathbf{u}_{\text{II}}^{(s)}|=1 \\ j \leq s < n}} \prod_{i < k < j} \frac{\Gamma_k(\mathbf{u}_{\text{III}}^{(k-1)}; \mathbf{u}_{\text{I}}^{(1\dots n)})}{\mathbf{u}_{\text{III}}^{(k-1)} - \mathbf{u}_{\text{III}}^{(k)}} \cdot \Gamma_j(\mathbf{u}_{\text{III}}^{(j-1)}; \mathbf{u}_{\text{I}}^{(1\dots n)}) \\
 & \quad \times \prod_{j < k \leq n} \frac{\Gamma_k(\mathbf{u}_{\text{II}}^{(k-1)}; \mathbf{u}_{\text{I}}^{(1\dots n)}) \Gamma_k(\mathbf{u}_{\text{III}}^{(k-1)}; \mathbf{u}_{\text{I,II}}^{(1\dots n)})}{(\mathbf{u}_{\text{II}}^{(k-1)} - \mathbf{u}_{\text{II}}^{(k)})(\mathbf{u}_{\text{III}}^{(k-1)} - \mathbf{u}_{\text{III}}^{(k)})} \\
 & \quad \times \left( \frac{f^+(\mathbf{u}_{\text{III}}^{(j-1)}, \mathbf{u}_{\text{II}}^{(j)})}{\mathbf{u}_{\text{III}}^{(j-1)} - \mathbf{u}_{\text{III}}^{(j)}} s_{i,2n-j+1}(\mathbf{u}_{\text{III}}^{(n)}) + \frac{1}{\mathbf{u}_{\text{III}}^{(j-1)} - \mathbf{u}_{\text{II}}^{(j)}} s_{j,2n-i+1}(\mathbf{u}_{\text{III}}^{(n)}) \right) \Psi(\mathbf{u}_{\text{I}}^{(1\dots n)})
 \end{aligned} \tag{4.9}$$

448 where  $\mathbf{u}_{\text{III}}^{(n)} = \mathbf{u}_j^{(n)}$ ,  $\mathbf{u}_{\text{II}}^{(n)} = \tilde{\mathbf{u}}_j^{(n)}$  and  $\mathbf{u}_{\text{I}}^{(n)} = \mathbf{u}^{(n)} \setminus \mathbf{u}_j^{(n)}$  for any  $1 \leq j \leq m_n$ , and  $\mathbf{u}_{\text{III}}^{(r)} = \mathbf{u}_{\text{II}}^{(r)} = \emptyset$  for  
 449 all  $1 \leq r < i$  in the first sum,  $\mathbf{u}_{\text{III}}^{(r)} = \mathbf{u}_{\text{II}}^{(s)} = \emptyset$  for all  $1 \leq r < i$  and  $1 \leq s < j$  in the second sum,  
 450 and  $\Gamma_k(\mathbf{u}_{\text{III}}^{(k-1)}; \mathbf{u}_{\text{I,II}}^{(1\dots n)})$  when  $k = n$  denotes  $f^+(\mathbf{u}_{\text{III}}^{(n-1)}, \mathbf{u}_{\text{II}}^{(n)}) \Gamma_n(\mathbf{u}_{\text{III}}^{(n-1)}; \mathbf{u}_{\text{I}}^{(1\dots n)})$ .

451 *Example 4.5.* When  $n = 2$ , the recurrence relation (4.9) gives

$$\begin{aligned} \Psi(\mathbf{u}^{(1,2)}) &= s_{23}(\mathbf{u}_{\text{III}}^{(2)})\Psi(\mathbf{u}_I^{(1,2)}) + \sum_{|\mathbf{u}_{\text{II}}^{(1)}|=2} K(\mathbf{u}_{\text{II}}^{(1)} | \mathbf{u}_{\text{II,III}}^{(2)})\Gamma_2(\mathbf{u}_{\text{II}}^{(1)}; \mathbf{u}_I^{(1,2)})s_{14}(\mathbf{u}_{\text{III}}^{(2)})\Psi(\mathbf{u}_I^{(1,2)}) \\ &+ \sum_{|\mathbf{u}_{\text{III}}^{(1)}|=1} \Gamma_2(\mathbf{u}_{\text{III}}^{(1)}; \mathbf{u}_I^{(1,2)}) \left( \frac{f^+(\mathbf{u}_{\text{III}}^{(1)}, \mathbf{u}_{\text{II}}^{(2)})}{\mathbf{u}_{\text{III}}^{(1)} - \mathbf{u}_{\text{III}}^{(2)}} s_{13}(\mathbf{u}_{\text{III}}^{(2)}) + \frac{1}{\mathbf{u}_{\text{III}}^{(1)} - \mathbf{u}_{\text{II}}^{(2)}} s_{24}(\mathbf{u}_{\text{III}}^{(2)}) \right) \Psi(\mathbf{u}_I^{(1,2)}) \end{aligned} \quad (4.10)$$

452 where  $\mathbf{u}_{\text{III}}^{(2)} = u_j^{(2)}$ ,  $\mathbf{u}_{\text{II}}^{(2)} = \tilde{u}_j^{(2)}$  and  $\mathbf{u}_I^{(2)} = \mathbf{u}^{(2)} \setminus u_j^{(2)}$  for any  $1 \leq j \leq m_2$ , and  $\mathbf{u}_{\text{III}}^{(1)} = \emptyset$  in the first  
453 sum and  $\mathbf{u}_{\text{II}}^{(1)} = \emptyset$  in the second sum.

454 *Proof of Proposition 4.4.* By Lemma 3.7, it is sufficient to consider the  $j = m_n$  case. Recall  
455 (3.28), (3.42) and consider a level- $(n-1)$  vector

$$\mathcal{R}^{(\hat{n})}(\mathbf{u}_{m_n}^{(n)}; \mathbf{u}^{(n)} \setminus u_{m_n}^{(n)})\Psi(\mathbf{u}^{(1\dots n-1)} | \mathbf{u}^{(n)}). \quad (4.11)$$

456 With the help of Yang-Baxter equation we can move operator  $\mathcal{R}^{(\hat{n})}(\mathbf{u}_{m_n}^{(n)}; \mathbf{u}^{(n)} \setminus u_{m_n}^{(n)})$  all way to  
457 the reference vector  $\eta^m$ . As a result of this, the level- $(n-1)$  nested monodromy matrix (3.17)  
458 factorises as

$$\widehat{R}_{\hat{a}_{m_n} a}^{(\hat{n}, \hat{n})}(\mathbf{u}_{m_n}^{(n)} - v) \widehat{R}_{\hat{a}_{m_n} a}^{(\hat{n}, \hat{n})}(\tilde{u}_{m_n}^{(n)} - v) T_a^{(\hat{n})}(v; \mathbf{u}^{(n)} \setminus u_{m_n}^{(n)}). \quad (4.12)$$

459 Since  $\mathcal{R}^{(\hat{n})}(\mathbf{u}_{m_n}^{(n)}; \mathbf{u}^{(n)} \setminus u_{m_n}^{(n)}) \cdot \eta^m = \eta^m$  when  $\hat{n} = n$ , we may view vector (4.11) as a  $Y(\mathfrak{gl}_n)$ -  
460 based Bethe vector with monodromy matrix (4.12) and apply expansion (4.4) in the space  
461  $V_{\hat{a}_{m_n}}^{(n)} \otimes V_{\tilde{a}_{m_n}}^{(n)}$ . Recall (3.27), (3.47), (3.48) and act with  $\hat{\beta}_{\hat{a}_{m_n} \tilde{a}_{m_n}}^{(n)}(\mathbf{u}_{m_n}^{(n)}) \mathcal{B}^{(n)}(\mathbf{u}^{(n)} \setminus u_{m_n}^{(n)})$  on the  
462 resulting expression. This immediately gives the wanted result.  $\square$

463 **Proposition 4.6.**  $Y^+(\mathfrak{gl}_{2n+1})$ -based Bethe vectors satisfy the recurrence relation

$$\begin{aligned} \Psi(\mathbf{u}^{(1\dots n)}) &= \sum_{1 \leq i \leq n} \sum_{\substack{|\mathbf{u}_{\text{III}}^{(r)}|=1 \\ i \leq r < n}} \prod_{i < k \leq n} \frac{\Gamma_k(\mathbf{u}_{\text{III}}^{(k-1)}; \mathbf{u}_I^{(1\dots n)})}{\mathbf{u}_{\text{III}}^{(k-1)} - \mathbf{u}_{\text{III}}^{(k)}} s_{i, \hat{n}}(\mathbf{u}_{\text{III}}^{(n)}) \Psi(\mathbf{u}_I^{(1\dots n)}) \\ &+ \sum_{1 \leq i < n} \sum_{\substack{|\mathbf{u}_{\text{II}}^{(r)}|=1 \\ i \leq r \leq n}} \prod_{i < k \leq n} \frac{\Gamma_k(\mathbf{u}_{\text{II}}^{(k-1)}; \mathbf{u}_I^{(1\dots n)})}{\mathbf{u}_{\text{II}}^{(k-1)} - \mathbf{u}_{\text{II}}^{(k)}} \cdot \frac{\Gamma_{\hat{n}}(\mathbf{u}_{\text{II}}^{(n)}; \mathbf{u}_I^{(1\dots n)})}{\mathbf{u}_{\text{II}}^{(n)} - \tilde{\mathbf{u}}_{\text{III}}^{(n)}} \\ &\quad \times \left( \frac{\mathbf{u}_{\text{II}}^{(n-1)} - \mathbf{u}_{\text{II}}^{(n)} + 1}{\mathbf{u}_{\text{II}}^{(n-1)} - \mathbf{u}_{\text{III}}^{(n)}} s_{i, \hat{n}+1}(\mathbf{u}_{\text{III}}^{(n)}) + s_{n, \hat{n}+1}(\mathbf{u}_{\text{III}}^{(n)}) \right) \Psi(\mathbf{u}_I^{(1\dots n)}) \\ &+ \sum_{1 \leq i \leq n} \sum_{\substack{|\mathbf{u}_{\text{II,III}}^{(r)}|=(2,0) \\ i \leq r < n}} \sum_{|\mathbf{u}_{\text{II}}^{(n)}|=1} \prod_{i < k \leq n} \Gamma_k(\mathbf{u}_{\text{II}}^{(k-1)}; \mathbf{u}_I^{(1\dots n)}) K(\mathbf{u}_{\text{II}}^{(k-1)} | \mathbf{u}_{\text{II,III}}^{(k)}) \\ &\quad \times \frac{\Gamma_{\hat{n}}(\mathbf{u}_{\text{II}}^{(n)}; \mathbf{u}_I^{(1\dots n)})}{\mathbf{u}_{\text{II}}^{(n)} - \tilde{\mathbf{u}}_{\text{III}}^{(n)}} s_{i, 2\hat{n}-i}(\mathbf{u}_{\text{III}}^{(n)}) \Psi(\mathbf{u}_I^{(1\dots n)}) \\ &+ \sum_{1 \leq i < j < n} \sum_{\substack{|\mathbf{u}_{\text{III}}^{(r)}|=1 \\ i \leq r < n}} \sum_{\substack{|\mathbf{u}_{\text{II}}^{(s)}|=1 \\ j \leq s \leq n}} \prod_{i < k < j} \frac{\Gamma_k(\mathbf{u}_{\text{III}}^{(k-1)}; \mathbf{u}_I^{(1\dots n)})}{\mathbf{u}_{\text{III}}^{(k-1)} - \mathbf{u}_{\text{III}}^{(k)}} \cdot \Gamma_j(\mathbf{u}_{\text{III}}^{(j-1)}; \mathbf{u}_I^{(1\dots n)}) \\ &\quad \times \prod_{j < k < n} \frac{\Gamma_k(\mathbf{u}_{\text{II}}^{(k-1)}; \mathbf{u}_I^{(1\dots n)}) \Gamma_k(\mathbf{u}_{\text{III}}^{(k-1)}; \mathbf{u}_{\text{I,II}}^{(1\dots n)})}{(\mathbf{u}_{\text{II}}^{(k-1)} - \mathbf{u}_{\text{II}}^{(k)}) (\mathbf{u}_{\text{III}}^{(k-1)} - \mathbf{u}_{\text{III}}^{(k)})} \cdot \Gamma_n(\mathbf{u}_{\text{II,III}}^{(k-1)}; \mathbf{u}_I^{(1\dots n)}) \Gamma_{\hat{n}}(\mathbf{u}_{\text{II}}^{(n)}; \mathbf{u}_I^{(1\dots n)}) \end{aligned}$$

$$\begin{aligned} & \times \left[ \left( \left( \beta_0 + \frac{\beta_2}{2\gamma} \right) \frac{f^+(\mathbf{u}_{\text{III}}^{(j-1)}, \mathbf{u}_{\text{II}}^{(j)})}{\mathbf{u}_{\text{III}}^{(j-1)} - \mathbf{u}_{\text{III}}^{(j)}} + \frac{\beta_1}{2\gamma} \cdot \frac{1}{\mathbf{u}_{\text{III}}^{(j-1)} - \mathbf{u}_{\text{II}}^{(j)}} \right) s_{i,2\hat{n}-j}(\mathbf{u}_{\text{III}}^{(n)}) \right. \\ & \quad \left. + \left( \frac{\beta_1}{2\gamma} \cdot \frac{f^+(\mathbf{u}_{\text{III}}^{(j-1)}, \mathbf{u}_{\text{II}}^{(j)})}{\mathbf{u}_{\text{III}}^{(j-1)} - \mathbf{u}_{\text{III}}^{(j)}} + \left( \beta_0 + \frac{\beta_2}{2\gamma} \right) \frac{1}{\mathbf{u}_{\text{III}}^{(j-1)} - \mathbf{u}_{\text{II}}^{(j)}} \right) s_{j,2\hat{n}-i}(\mathbf{u}_{\text{III}}^{(n)}) \right] \\ & \quad \times \Psi(\mathbf{u}_{\text{I}}^{(1\dots n)}) \end{aligned} \quad (4.13)$$

464 where

$$\begin{aligned} \beta_0 &= \frac{f^-(\mathbf{u}_{\text{III}}^{(n-1)}, \mathbf{u}_{\text{II}}^{(n-1)}) f^+(\mathbf{u}_{\text{III}}^{(n-1)}, \tilde{\mathbf{u}}_{\text{III}}^{(n)})}{(\mathbf{u}_{\text{II}}^{(n-1)} - \tilde{\mathbf{u}}_{\text{III}}^{(n)})(\mathbf{u}_{\text{III}}^{(n-1)} - \mathbf{u}_{\text{III}}^{(n)})(\mathbf{u}_{\text{II}}^{(n)} - \tilde{\mathbf{u}}_{\text{III}}^{(n)})}, \\ \beta_1 &= \frac{\mathbf{u}_{\text{III}}^{(n-1)} - \mathbf{u}_{\text{III}}^{(n)}}{\mathbf{u}_{\text{III}}^{(n-1)} - \mathbf{u}_{\text{II}}^{(n)}} \left( \mathbf{u}_{\text{II}}^{(n-1)} - \tilde{\mathbf{u}}_{\text{III}}^{(n)} + 1 + \frac{\mathbf{u}_{\text{III}}^{(n-1)} - \tilde{\mathbf{u}}_{\text{III}}^{(n)}}{\mathbf{u}_{\text{II}}^{(n-1)} - \mathbf{u}_{\text{II}}^{(n)}} \right), \\ \beta_2 &= f^+(\mathbf{u}_{\text{II}}^{(n-1)}, \mathbf{u}_{\text{II}}^{(n)}) \frac{\mathbf{u}_{\text{II}}^{(n-1)} - \tilde{\mathbf{u}}_{\text{III}}^{(n)}}{\mathbf{u}_{\text{III}}^{(n-1)} - \mathbf{u}_{\text{II}}^{(n)}} + \frac{(\mathbf{u}_{\text{II}}^{(n-1)} - \mathbf{u}_{\text{III}}^{(n)})(\mathbf{u}_{\text{III}}^{(n-1)} - \tilde{\mathbf{u}}_{\text{III}}^{(n)}) + 1}{\mathbf{u}_{\text{II}}^{(n-1)} - \mathbf{u}_{\text{II}}^{(n)}} \end{aligned} \quad (4.14)$$

465 and

$$\gamma = (\mathbf{u}_{\text{II}}^{(n-1)} - \tilde{\mathbf{u}}_{\text{III}}^{(n)})(\mathbf{u}_{\text{III}}^{(n-1)} - \tilde{\mathbf{u}}_{\text{III}}^{(n)})(\mathbf{u}_{\text{II}}^{(n-1)} - \mathbf{u}_{\text{III}}^{(n)})(\mathbf{u}_{\text{III}}^{(n-1)} - \mathbf{u}_{\text{III}}^{(n)}) \quad (4.15)$$

466 and  $\mathbf{u}_{\text{III}}^{(n)} = u_j^{(n)}$  for any  $1 \leq j \leq m_n$ , and  $\mathbf{u}_{\text{III}}^{(r)} = \mathbf{u}_{\text{II}}^{(s)} = \emptyset$  for all  $1 \leq r < i$  and  $1 \leq s \leq n$  in the  
 467 first sum,  $\mathbf{u}_{\text{II}}^{(r)} = \mathbf{u}_{\text{III}}^{(s)} = \emptyset$  for all  $1 \leq r < i$  and  $1 \leq s < n$  in the second sum,  $\mathbf{u}_{\text{II}}^{(r)} = \mathbf{u}_{\text{III}}^{(s)} = \emptyset$  for  
 468 all  $1 \leq r < i$  and  $1 \leq s < n$  in the third sum and  $\mathbf{u}_{\text{III}}^{(r)} = \mathbf{u}_{\text{II}}^{(s)} = \emptyset$  for all  $1 \leq r < i$  and  $1 \leq s < j$   
 469 in the last sum.

470 Example 4.7. When  $n = 1$ , the recurrence relation (4.13) gives

$$\Psi(\mathbf{u}^{(1)}) = s_{12}(\mathbf{u}_{\text{III}}^{(1)}) \Psi(\mathbf{u}_{\text{I}}^{(1)}) + \sum_{|\mathbf{u}_{\text{III}}^{(1)}|=1} \frac{\Gamma_2(\mathbf{u}_{\text{II}}^{(1)}; \mathbf{u}_{\text{I}}^{(1)})}{\mathbf{u}_{\text{II}}^{(1)} - \tilde{\mathbf{u}}_{\text{III}}^{(1)}} s_{13}(\mathbf{u}_{\text{III}}^{(1)}) \Psi(\mathbf{u}_{\text{I}}^{(1)}) \quad (4.16)$$

471 where  $\mathbf{u}_{\text{III}}^{(1)} = u_j^{(1)}$  for any  $1 \leq j \leq m_1$ . When  $n = 2$ , we have

$$\begin{aligned} \Psi(\mathbf{u}^{(1,2)}) &= s_{23}(\mathbf{u}_{\text{III}}^{(2)}) \Psi(\mathbf{u}_{\text{I}}^{(1,2)}) + \sum_{|\mathbf{u}_{\text{III}}^{(1)}|=1} \frac{\Gamma_2(\mathbf{u}_{\text{III}}^{(1)}; \mathbf{u}_{\text{I}}^{(1,2)})}{\mathbf{u}_{\text{III}}^{(1)} - \mathbf{u}_{\text{III}}^{(2)}} s_{13}(\mathbf{u}_{\text{III}}^{(2)}) \Psi(\mathbf{u}_{\text{I}}^{(1,2)}) \\ &+ \sum_{|\mathbf{u}_{\text{II}}^{(1,2)}|=(1,1)} \frac{\Gamma_2(\mathbf{u}_{\text{II}}^{(1)}; \mathbf{u}_{\text{I}}^{(1,2)}) \Gamma_3(\mathbf{u}_{\text{II}}^{(2)}; \mathbf{u}_{\text{I}}^{(1,2)})}{\mathbf{u}_{\text{II}}^{(2)} - \tilde{\mathbf{u}}_{\text{III}}^{(2)}} \\ &\quad \times \left( \frac{f^+(\mathbf{u}_{\text{II}}^{(1)}, \mathbf{u}_{\text{II}}^{(2)})}{\mathbf{u}_{\text{II}}^{(1)} - \mathbf{u}_{\text{III}}^{(2)}} s_{14}(\mathbf{u}_{\text{III}}^{(2)}) + \frac{1}{\mathbf{u}_{\text{II}}^{(1)} - \mathbf{u}_{\text{II}}^{(2)}} s_{25}(\mathbf{u}_{\text{III}}^{(2)}) \right) \Psi(\mathbf{u}_{\text{I}}^{(1,2)}) \\ &+ \sum_{|\mathbf{u}_{\text{II}}^{(1,2)}|=(2,1)} \frac{\Gamma_2(\mathbf{u}_{\text{II}}^{(1)}; \mathbf{u}_{\text{I}}^{(1,2)}) \Gamma_3(\mathbf{u}_{\text{II}}^{(2)}; \mathbf{u}_{\text{I}}^{(1,2)})}{\mathbf{u}_{\text{II}}^{(2)} - \tilde{\mathbf{u}}_{\text{III}}^{(2)}} K(\mathbf{u}_{\text{II}}^{(1)} | \mathbf{u}_{\text{II,III}}^{(2)}) s_{15}(\mathbf{u}_{\text{III}}^{(2)}) \Psi(\mathbf{u}_{\text{I}}^{(1,2)}) \\ &+ \sum_{|\mathbf{u}_{\text{II}}^{(2)}|=1} \frac{\Gamma_3(\mathbf{u}_{\text{II}}^{(2)}; \mathbf{u}_{\text{I}}^{(2)})}{\mathbf{u}_{\text{II}}^{(2)} - \tilde{\mathbf{u}}_{\text{III}}^{(2)}} s_{24}(\mathbf{u}_{\text{III}}^{(2)}) \Psi(\mathbf{u}_{\text{I}}^{(1,2)}) \end{aligned} \quad (4.17)$$

472 where  $\mathbf{u}_{\text{III}}^{(2)} = u_j^{(2)}$  for any  $1 \leq j \leq m_2$ , and  $\mathbf{u}_{\text{II}}^{(1)} = \mathbf{u}_{\text{II}}^{(2)} = \emptyset$  in the first sum,  $\mathbf{u}_{\text{III}}^{(1)} = \emptyset$  in the  
 473 second sum,  $\mathbf{u}_{\text{III}}^{(1)} = \emptyset$  in the third sum and  $\mathbf{u}_{\text{III}}^{(1)} = \mathbf{u}_{\text{II}}^{(1)} = \emptyset$  in the last sum.

474 The technical Lemma below will assist us in proving Proposition 4.6.

475 **Lemma 4.8.** Let  $\Psi_j(\mathbf{u}^{(1\dots n)})$  denote a  $Y^+(\mathfrak{gl}_{2n+1})$ -based Bethe vector with the reference vector  
 476  $\eta_j^m := (E_{12}^{(\hat{n})})_{\hat{a}_j} \eta^m$ . Then

$$\Psi_j(\mathbf{u}^{(1\dots n)}) = \sum_{1 \leq i \leq j} \frac{1}{u_j^{(n)} - u_i^{(n)} + 1} \frac{\Gamma_{\hat{n}}(u_i^{(n)}, \mathbf{u}^{(1\dots n)} \setminus u_i^{(n)})}{\prod_{k>j} f^+(u_k^{(n)}, u_i^{(n)})} \Psi(\mathbf{u}^{(1\dots n)} \setminus u_i^{(n)}). \quad (4.18)$$

477 *Proof.* Recall (3.26) and consider level- $(n-1)$  vector

$$\prod_{j>1}^{\rightarrow} R_{\hat{a}_1 \hat{a}_j}^{(\hat{n}, \hat{n})}(\tilde{u}_1^{(n)} - u_j^{(n)}) \Psi_1(\mathbf{u}^{(1\dots n-1)} | \mathbf{u}^{(n)}). \quad (4.19)$$

478 With the help of Yang-Baxter equation we can move the product of  $R$ -matrices all way to the  
 479 reference vector  $\eta_1^m$ . As a result of this, the level- $(n-1)$  nested monodromy matrix (3.17)  
 480 takes the form

$$\prod_{i>1}^{\leftarrow} \widehat{R}_{\hat{a}_i \hat{a}}^{(\hat{n}, \hat{n})}(u_i^{(n)} - v) \prod_{i>1}^{\leftarrow} \widehat{R}_{\hat{a}_i \hat{a}}^{(\hat{n}, \hat{n})}(\tilde{u}_i^{(n)} - v) \widehat{R}_{\hat{a}_1 \hat{a}}^{(\hat{n}, \hat{n})}(u_1^{(n)} - v) \widehat{R}_{\hat{a}_1 \hat{a}}^{(\hat{n}, \hat{n})}(\tilde{u}_1^{(n)} - v) A_{\hat{a}}^{(\hat{n})}(v). \quad (4.20)$$

481 In the space  $L^{(n-1)'}$ , it is equivalent to  $T_a^{(n)'}(v; \mathbf{u}^{(n)} \setminus u_1^{(n)})$ . Next, recall (3.41) and note that

$$\prod_{j>1}^{\rightarrow} R_{\hat{a}_1 \hat{a}_j}^{(\hat{n}, \hat{n})}(\tilde{u}_1^{(n)} - u_j^{(n)}) \cdot \eta_1^m = f^+(u_1^{(n)}, \tilde{\mathbf{u}}^{(n)} \setminus \tilde{u}_1^{(n)}) \eta_1^m. \quad (4.21)$$

482 Hence, vector (4.19) can be expanded in the space  $V_{\hat{a}_1}^{(\hat{n})} \otimes V_{\hat{a}_1}^{(\hat{n})}$  as

$$f^+(u_1^{(n)}, \tilde{\mathbf{u}}^{(n)} \setminus \tilde{u}_1^{(n)}) \cdot E_1^{(\hat{n})} \otimes E_1^{(\hat{n})} \otimes \Psi_1(\mathbf{u}^{(1\dots n-1)} | \mathbf{u}^{(n)} \setminus u_1^{(n)}). \quad (4.22)$$

483 From (3.27) note that  $\check{b}_{\hat{a}_1 \hat{a}_1}^{(n)}(v) \cdot E_1^{(\hat{n})} \otimes E_1^{(\hat{n})} = s_{\hat{n}\hat{n}}(v)$ . Defining relations of  $Y^+(\mathfrak{gl}_{2n+1})$  imply  
 484 that

$$s_{\hat{n}\hat{n}}(u_1^{(n)}) \prod_{i<n}^{\leftarrow} \mathcal{B}^{(i)}(\mathbf{u}^{(i)}; \mathbf{u}^{(i+1\dots n)} \setminus u_1^{(n)}) = \prod_{i<n}^{\leftarrow} \mathcal{B}^{(i)}(\mathbf{u}^{(i)}; \mathbf{u}^{(i+1\dots n)} \setminus u_1^{(n)}) s_{\hat{n}\hat{n}}(u_1^{(n)}) + UWT$$

485 where  $UWT$  denotes “unwanted” terms, all of which act by 0 on  $\eta_1^m$ . We have thus shown  
 486 that

$$\begin{aligned} \Psi_1(\mathbf{u}^{(1\dots n)}) &= \mathcal{B}^{(n)}(\mathbf{u}^{(n)} \setminus u_1^{(n)}) \check{b}_{\hat{a}_1 \hat{a}_1}^{(n)}(u_1^{(n)}) \prod_{j>1}^{\rightarrow} R_{\hat{a}_1 \hat{a}_j}^{(\hat{n}, \hat{n})}(\tilde{u}_1^{(n)} - u_j^{(n)}) \Psi_1(\mathbf{u}^{(1\dots n-1)} | \mathbf{u}^{(n)}) \\ &= \mu_{\hat{n}}(v) f^+(u_1^{(n)}, \tilde{\mathbf{u}}^{(n)} \setminus \tilde{u}_1^{(n)}) \Psi(\mathbf{u}^{(1\dots n)} \setminus u_1^{(n)}). \end{aligned} \quad (4.23)$$

487 This gives the  $j = 1$  case of the claim. Then, using Yang-Baxter equation, Lemma 3.4, and the  
 488 identity

$$\eta_{j+1}^m = f^+(u_j^{(n)}, u_{j+1}^{(n)}) \check{R}_{\hat{a}_{j+1} \hat{a}_j}^{(\hat{n}, \hat{n})}(u_{j+1}^{(n)} - u_j^{(n)}) \check{R}_{\hat{a}_{j+1} \hat{a}_j}^{(\hat{n}, \hat{n})}(u_j^{(n)} - u_{j+1}^{(n)}) \cdot \eta_j^m + \frac{1}{u_{j+1}^{(n)} - u_j^{(n)}} \eta_j^m$$

489 we find

$$\Psi_{j+1}(\mathbf{u}^{(1\dots n)}) = f^+(u_j^{(n)}, u_{j+1}^{(n)}) \Psi_j(\mathbf{u}^{(1\dots n)} \setminus u_{j+1}^{(n)}) + \frac{1}{u_{j+1}^{(n)} - u_j^{(n)}} \Psi_j(\mathbf{u}^{(1\dots n)}). \quad (4.24)$$

490 A simple induction on  $j$  together with Lemma 3.7 gives the wanted result.  $\square$

491 *Proof of Proposition 4.6.* The main idea of the proof is similar to that of Proposition 4.4.  
 492 We start from the level- $(n-1)$  vector (4.11) and move operator  $\mathcal{R}^{(\hat{n})}(u_{m_n}^{(n)}; \mathbf{u}^{(n)} \setminus u_{m_n}^{(n)})$  all way to  
 493 the reference vector  $\eta^m$ . In the odd case  $E_1^{(\hat{n})} = E_2^{(n+1)}$  giving (recall (3.29))

$$\mathcal{R}^{(\hat{n})}(u_{m_n}^{(n)}; \mathbf{u}^{(n)} \setminus u_{m_n}^{(n)}) \cdot \eta^m = \eta^m + \sum_{j < m_n} \frac{\prod_{j < k < m_n} f^+(u_k^{(n)}, \tilde{u}_{m_n}^{(n)})}{u_j^{(n)} - \tilde{u}_{m_n}^{(n)}} P_{\hat{a}_j \hat{a}_{m_n}}^{(\hat{n}, \hat{n})} \eta^m. \quad (4.25)$$

494 Hence, in the odd case we can rewrite (4.11) as

$$\dot{\Psi}_{2,1}(\mathbf{u}^{(1\dots n-1)} | \mathbf{u}^{(n)}) + \sum_{j < m_n} \frac{\prod_{j < k < m_n} f^+(u_k^{(n)}, \tilde{u}_{m_n}^{(n)})}{u_j^{(n)} - \tilde{u}_{m_n}^{(n)}} \dot{\Psi}_{2,2;j}(\mathbf{u}^{(1\dots n-1)} | \mathbf{u}^{(n)}) \quad (4.26)$$

495 where  $\dot{\Psi}_{k,l}$  and  $\dot{\Psi}_{k,l;j}$  denote level- $(n-1)$  Bethe vectors based on the transfer matrix (4.12) and  
 496 reference vectors  $(E_{k,2}^{(\hat{n})})_{\hat{a}_{m_n}} (E_{l,1}^{(\hat{n})})_{\hat{a}_{m_n}} \eta^m$  and  $(E_{k,2}^{(\hat{n})})_{\hat{a}_{m_n}} (E_{l,1}^{(\hat{n})})_{\hat{a}_{m_n}} (E_{1,2}^{(\hat{n})})_{\hat{a}_j} \eta^m$ , respectively.

497 Consider the second term in (4.26). Acting with  $\mathcal{B}^{(n)}(\mathbf{u}^{(n)} \setminus u_{m_n}^{(n)})$  and applying Lemma 4.8  
 498 gives

$$\begin{aligned} \sum_{i \leq j < m_n} \prod_{j < k < m_n} \frac{f^+(u_k^{(n)}, \tilde{u}_{m_n}^{(n)})}{f^+(u_k^{(n)}, u_i^{(n)})} \cdot \frac{\Gamma_{\hat{n}}(u_i^{(n)}; \mathbf{u}^{(1\dots n)} \setminus (u_i^{(n)}, u_{m_n}^{(n)}))}{(u_j^{(n)} - u_i^{(n)} + 1)(u_j^{(n)} - \tilde{u}_{m_n}^{(n)})} \\ \times \mathcal{B}^{(n)}(\mathbf{u}^{(n)} \setminus (u_i^{(n)}, u_{m_n}^{(n)})) \dot{\Psi}_{2,2}(\mathbf{u}^{(1\dots n-1)} | \mathbf{u}^{(n)} \setminus u_i^{(n)}). \end{aligned} \quad (4.27)$$

499 Using the identity

$$\frac{1}{u_i^{(n)} - \tilde{u}_{m_n}^{(n)}} = \sum_{i \leq j < m_n} \frac{1}{(u_j^{(n)} - u_i^{(n)} + 1)(u_j^{(n)} - \tilde{u}_{m_n}^{(n)})} \prod_{j < k < m_n} \frac{f^+(u_k^{(n)}, \tilde{u}_{m_n}^{(n)})}{f^+(u_k^{(n)}, u_i^{(n)})} \quad (4.28)$$

500 which follows by a descending induction on  $i$ , expression (4.27) becomes

$$\sum_{i < m_n} \frac{\Gamma_{\hat{n}}(u_i^{(n)}; \mathbf{u}^{(n)} \setminus (u_i^{(n)}, u_{m_n}^{(n)}))}{u_i^{(n)} - \tilde{u}_{m_n}^{(n)}} \mathcal{B}^{(n)}(\mathbf{u}^{(n)} \setminus (u_i^{(n)}, u_{m_n}^{(n)})) \dot{\Psi}_{2,2}(\mathbf{u}^{(1\dots n-1)} | \mathbf{u}^{(n)} \setminus u_i^{(n)}). \quad (4.29)$$

501 Thus, acting with  $\theta_{\hat{a}_{m_n} \hat{a}_{m_n}}^{(n)}(u_{m_n}^{(n)}) \mathcal{B}^{(n)}(\mathbf{u}^{(n)} \setminus u_{m_n}^{(n)})$  on (4.26) we obtain

$$\begin{aligned} \Psi(\mathbf{u}^{(1\dots n)}) = \theta_{\hat{a}_{m_n} \hat{a}_{m_n}}^{(n)}(u_{m_n}^{(n)}) \left( \mathcal{B}^{(n)}(\mathbf{u}^{(n)} \setminus u_{m_n}^{(n)}) \dot{\Psi}_{2,1}(\mathbf{u}^{(1\dots n-1)} | \mathbf{u}^{(n)}) \right. \\ \left. + \sum_{i < m_n} \frac{\Gamma_{\hat{n}}(u_i^{(n)}; \mathbf{u}^{(1\dots n)} \setminus (u_i^{(n)}, u_{m_n}^{(n)}))}{u_i^{(n)} - \tilde{u}_{m_n}^{(n)}} \right. \\ \left. \times \mathcal{B}^{(n)}(\mathbf{u}^{(n)} \setminus (u_i^{(n)}, u_{m_n}^{(n)})) \dot{\Psi}_{2,2}(\mathbf{u}^{(1\dots n-1)} | \mathbf{u}^{(n)} \setminus u_i^{(n)}) \right). \end{aligned} \quad (4.30)$$

502 We will view vectors  $\dot{\Psi}_{2,1}$  and  $\dot{\Psi}_{2,2}$  as  $Y(\mathfrak{gl}_n)$ -based Bethe vectors and apply  $Y(\mathfrak{gl}_n)$ -based re-  
 503 currence relations.

504 First, consider vector  $\dot{\Psi}_{2,2}$ . Its reference vector is annihilated by the  $(j, i)$ -th entries of the  
 505 monodromy matrix (4.12) satisfying the condition  $i < j$ . Hence, we may use (4.4) to obtain



506 an expansion in the space  $V_{\hat{a}_{m_n}}^{(\hat{n})} \otimes V_{\hat{a}_{m_n}}^{(\hat{n})}$ . Taking  $\mathbf{u}_{\text{III}}^{(n)} = u_{m_n}^{(n)}$ , the second term inside the brackets  
 507 of (4.30) becomes (we have singled out the  $i < j = n$  terms for further convenience)

$$\sum_{1 \leq i \leq n} \sum_{\substack{|\mathbf{u}_{\text{II,III}}^{(r)}|=(2,0) \\ i \leq r < n}} \sum_{|\mathbf{u}_{\text{II}}^{(n)}|=1} \prod_{i < k \leq n} K(\mathbf{u}_{\text{II}}^{(k-1)} | \mathbf{u}_{\text{II,III}}^{(k)}) \Gamma_k(\mathbf{u}_{\text{II}}^{(k-1)}; \mathbf{u}_{\text{I}}^{(1\dots n)}) \\ \times \frac{\Gamma_{\hat{n}}(\mathbf{u}_{\text{II}}^{(n)}; \mathbf{u}_{\text{I}}^{(1\dots n)})}{\mathbf{u}_{\text{II}}^{(n)} - \tilde{\mathbf{u}}_{\text{III}}^{(n)}} E_{\bar{i}}^{(\hat{n})} \otimes E_{\bar{i}}^{(\hat{n})} \otimes \Psi(\mathbf{u}_{\text{I}}^{(1\dots n)}) \quad (4.31)$$

$$+ \sum_{1 \leq i < n} \sum_{\substack{|\mathbf{u}_{\text{II}}^{(r)}|=1 \\ i \leq r \leq n}} \prod_{i < k < n} \frac{\Gamma_k(\mathbf{u}_{\text{II}}^{(k-1)}; \mathbf{u}_{\text{I}}^{(1\dots n)})}{\mathbf{u}_{\text{II}}^{(k-1)} - \mathbf{u}_{\text{II}}^{(k)}} \cdot \frac{\Gamma_n(\mathbf{u}_{\text{II}}^{(n-1)}; \mathbf{u}_{\text{I}}^{(1\dots n)}) \Gamma_{\hat{n}}(\mathbf{u}_{\text{II}}^{(n)}; \mathbf{u}_{\text{I}}^{(1\dots n)})}{\mathbf{u}_{\text{II}}^{(n)} - \tilde{\mathbf{u}}_{\text{III}}^{(n)}} \\ \times \left( \frac{f^+(\mathbf{u}_{\text{II}}^{(n-1)}, \tilde{\mathbf{u}}_{\text{III}}^{(n)})}{\mathbf{u}_{\text{II}}^{(n-1)} - \mathbf{u}_{\text{III}}^{(n)}} E_{\bar{i}}^{(\hat{n})} \otimes E_{\bar{2}}^{(\hat{n})} + \frac{1}{\mathbf{u}_{\text{II}}^{(n-1)} - \tilde{\mathbf{u}}_{\text{III}}^{(n)}} E_{\bar{2}}^{(\hat{n})} \otimes E_{\bar{i}}^{(\hat{n})} \right) \otimes \Psi(\mathbf{u}_{\text{I}}^{(1\dots n)}) \quad (4.32)$$

$$+ \sum_{1 \leq i < j < n} \sum_{\substack{|\mathbf{u}_{\text{III}}^{(r)}|=1 \\ i \leq r < n}} \sum_{\substack{|\mathbf{u}_{\text{II}}^{(s)}|=1 \\ j \leq s \leq n}} \prod_{i < k < j} \frac{\Gamma_k(\mathbf{u}_{\text{III}}^{(k-1)}; \mathbf{u}_{\text{I}}^{(1\dots n)})}{\mathbf{u}_{\text{III}}^{(k-1)} - \mathbf{u}_{\text{III}}^{(k)}} \cdot \Gamma_j(\mathbf{u}_{\text{III}}^{(j-1)}; \mathbf{u}_{\text{I}}^{(1\dots n)}) \\ \times \prod_{j < k < n} \frac{\Gamma_k(\mathbf{u}_{\text{II}}^{(k-1)}; \mathbf{u}_{\text{I}}^{(1\dots n)}) \Gamma_k(\mathbf{u}_{\text{III}}^{(k-1)}; \mathbf{u}_{\text{I,II}}^{(1\dots n)})}{(\mathbf{u}_{\text{II}}^{(k-1)} - \mathbf{u}_{\text{II}}^{(k)})(\mathbf{u}_{\text{III}}^{(k-1)} - \mathbf{u}_{\text{III}}^{(k)})} \cdot \frac{\Gamma_n(\mathbf{u}_{\text{II,III}}^{(n-1)}; \mathbf{u}_{\text{I}}^{(1\dots n)}) \Gamma_{\hat{n}}(\mathbf{u}_{\text{II}}^{(n)}; \mathbf{u}_{\text{I}}^{(1\dots n)})}{\mathbf{u}_{\text{II}}^{(n)} - \tilde{\mathbf{u}}_{\text{III}}^{(n)}} \\ \times \frac{f^-(\mathbf{u}_{\text{III}}^{(n-1)}, \mathbf{u}_{\text{II}}^{(n-1)}) f^+(\mathbf{u}_{\text{III}}^{(n-1)}, \tilde{\mathbf{u}}_{\text{III}}^{(n)})}{(\mathbf{u}_{\text{II}}^{(n-1)} - \tilde{\mathbf{u}}_{\text{III}}^{(n)})(\mathbf{u}_{\text{III}}^{(n-1)} - \mathbf{u}_{\text{III}}^{(n)})} \\ \times \left( \frac{f^+(\mathbf{u}_{\text{III}}^{(j-1)}, \mathbf{u}_{\text{II}}^{(j)})}{\mathbf{u}_{\text{III}}^{(j-1)} - \mathbf{u}_{\text{III}}^{(j)}} E_{\bar{i}}^{(\hat{n})} \otimes E_{\bar{j}}^{(\hat{n})} + \frac{1}{\mathbf{u}_{\text{III}}^{(j-1)} - \mathbf{u}_{\text{II}}^{(j)}} E_{\bar{j}}^{(\hat{n})} \otimes E_{\bar{i}}^{(\hat{n})} \right) \otimes \Psi(\mathbf{u}_{\text{I}}^{(1\dots n)}). \quad (4.33)$$

508 Next, consider vector  $\hat{\Psi}_{2,1}$ . This time we can not apply expansion (4.4). Instead, we will  
 509 use the composite model approach to obtain the wanted expansion. Set  $L^{\text{II}} := V_{\hat{a}_{m_n}}^{(\hat{n})} \otimes V_{\hat{a}_{m_n}}^{(\hat{n})}$  and  
 510  $L^{\text{I}} := W_{\hat{a} \setminus \hat{a}_{m_n}}^{(\hat{n})} \otimes W_{\hat{a} \setminus \hat{a}_{m_n}}^{(\hat{n})} \otimes (L^{(n)})^0$  so that  $L^{(n-1)} \cong L^{\text{II}} \otimes L^{\text{I}}$ . Recall (3.31) and set

$$\omega_{a_i^{n-1}, k}^{\text{II}}(v) := \sum_{j < n} (E_j^{(n-1)})_{a_i^{n-1}}^* \otimes [R_{\hat{a}_{m_n} a}^{(\hat{n}, \hat{n})}(v - u_{m_n}^{(n)}) R_{\hat{a}_{m_n} a}^{(\hat{n}, \hat{n})}(v - \tilde{u}_{m_n}^{(n)})]_{n-j, k}, \\ \hat{\theta}_k^{\text{I}}(v) := [T_a^{(k+1)}(u_i^{(k)}; \mathbf{u}^{(n)} \setminus u_{m_n}^{(n)})]_{k, n}.$$

511 The cases when  $k = n$ ,  $\hat{n}$  will be denoted by

$$\hat{\theta}_{a_i^{n-1}}^{\text{II}}(v) := \omega_{a_i^{n-1}, n}^{\text{II}}(v), \quad p_{a_i^{n-1}}^{\text{II}}(v) := \omega_{a_i^{n-1}, \hat{n}}^{\text{II}}(v), \quad d^{\text{I}}(v) := \hat{\theta}_n^{\text{I}}(v), \quad c^{\text{I}}(v) := \hat{\theta}_{\hat{n}}^{\text{I}}(v)$$

512 so that

$$\hat{\theta}_{a_i^{n-1}}^{(n-1)}(v; \mathbf{u}^{(n)}) = \sum_{k < n} \omega_{a_i^{n-1}, k}^{\text{II}}(v) \hat{\theta}_k^{\text{I}}(v) + \hat{\theta}_{a_i^{n-1}}^{\text{II}}(v) d^{\text{I}}(v) + p_{a_i^{n-1}}^{\text{II}}(v) c^{\text{I}}(v).$$

513 This notation is reminiscent of the Bethe ansatz notation commonly used in the composite  
 514 model approach only  $p_{a_i^{n-1}}^{\text{II}}$  is an additional creation operator specific to the case at hand.

515 Consider the II-labelled operators. Their action on the reference state  $E_2^{(\hat{n})} \otimes E_1^{(\hat{n})} \in L^{\text{II}}$  is given

516 by

$$\begin{aligned}
\alpha_{a_i^{n-1}, j}^{\parallel}(\nu) \cdot E_2^{(\hat{n})} \otimes E_1^{(\hat{n})} &= (E_{n-j}^{(n-1)})_{a_i^{n-1}}^* \cdot E_2^{(\hat{n})} \otimes E_1^{(\hat{n})}, \\
\theta_{a_i^{n-1}}^{\parallel}(\nu) \cdot E_2^{(\hat{n})} \otimes E_1^{(\hat{n})} &= \frac{1}{\nu - u_{m_n}^{(n)}} \sum_{j < n} (E_j^{(n-1)})_{a_i^{n-1}}^* \cdot E_{j+2}^{(\hat{n})} \otimes E_1^{(\hat{n})}, \\
p_{a_i^{n-1}}^{\parallel}(\nu) \cdot E_2^{(\hat{n})} \otimes E_1^{(\hat{n})} &= \frac{1}{\nu - \tilde{u}_{m_n}^{(n)}} \sum_{j < n} (E_j^{(n-1)})_{a_i^{n-1}}^* \left( \frac{1}{\nu - u_{m_n}^{(n)}} E_{j+2}^{(\hat{n})} \otimes E_2^{(\hat{n})} + E_2^{(\hat{n})} \otimes E_{j+2}^{(\hat{n})} \right), \\
p_{a_i^{n-1}}^{\parallel}(w) \theta_{a_i^{n-1}}^{\parallel}(\nu) \cdot E_2^{(\hat{n})} \otimes E_1^{(\hat{n})} &= \frac{1}{(w - \tilde{u}_{m_n}^{(n)})(\nu - u_{m_n}^{(n)})} \sum_{j, k < n} (E_j^{(n-1)})_{a_i^{n-1}}^* (E_k^{(n-1)})_{a_i^{n-1}}^* \\
&\quad \times \left( \frac{1}{w - u_{m_n}^{(n)}} E_{j+2}^{(\hat{n})} \otimes E_{k+2}^{(\hat{n})} + E_{k+2}^{(\hat{n})} \otimes E_{j+2}^{(\hat{n})} \right).
\end{aligned}$$

517 The products  $\theta_{a_j^{n-1}}^{\parallel}(\nu) \theta_{a_i^{n-1}}^{\parallel}(u)$ ,  $p_{a_j^{n-1}}^{\parallel}(\nu) p_{a_i^{n-1}}^{\parallel}(u)$ , and  $p_{a_k^{n-1}}^{\parallel}(w) p_{a_j^{n-1}}^{\parallel}(\nu) \theta_{a_i^{n-1}}^{\parallel}(u)$  act by zero on  
518  $E_2^{(\hat{n})} \otimes E_1^{(\hat{n})}$ . The homogeneous ( $aa$  and  $bb$ ,  $pp$ ) exchange relations of the  $\parallel$ -labelled operators  
519 are analogous to (3.32) and (3.33), respectively. The mixed ( $ab$ ,  $ap$ ,  $bp$ ) exchange relations  
520 have the form

$$\alpha_{a_j^{n-1}}^{\parallel}(\nu) \theta_{a_i^{n-1}}^{\parallel}(u) = \theta_{a_i^{n-1}}^{\parallel}(u) \alpha_{a_j^{n-1}}^{\parallel}(\nu) R_{a_i^{n-1}, a_j^{n-1}}^{(n-1, n-1)}(u - \nu) + \frac{1}{u - \nu} \theta_{a_i^{n-1}}^{\parallel}(\nu) \alpha_{a_j^{n-1}}^{\parallel}(u) P_{a_j^{n-1}, a_i^{n-1}}^{(n-1, n-1)}.$$

521 Consider the  $\perp$ -labelled operators. The  $dc$ ,  $cb$ ,  $db$  exchange relations have the form

$$d^{\perp}(\nu) c^{\perp}(u) = f^{-}(\nu, u) c^{\perp}(u) d^{\perp}(\nu) + \frac{1}{\nu - u} c^{\perp}(\nu) d^{\perp}(u).$$

522 The standard Bethe ansatz arguments then imply

$$\begin{aligned}
&\overleftarrow{\prod}_i \theta_{a_i^{n-1}}^{(n-1)}(u_i^{(n-1)}; \mathbf{u}^{(n)}) \cdot E_2^{(\hat{n})} \otimes E_1^{(\hat{n})} \otimes \Psi^{(n-2)}(\mathbf{u}^{(1\dots n-2)} | \mathbf{u}^{(n-1, n)} \setminus u_{m_n}^{(n)}) \\
&= \left[ E_2^{(\hat{n})} \otimes E_1^{(\hat{n})} \otimes \overleftarrow{\prod}_i \theta_{a_i^{n-1}}^{\perp}(u_i^{(n-1)}) \right. \tag{4.34}
\end{aligned}$$

$$\begin{aligned}
&+ \sum_j \frac{f^{-}(u_j^{(n-1)}, \mathbf{u}^{(n-1)} \setminus u_j^{(n-1)})}{u_j^{(n-1)} - u_{m_n}^{(n)}} \sum_{k < n} E_{k+2}^{(\hat{n})} \otimes E_1^{(\hat{n})} \otimes (E_k^{(n-1)})_{a_j^{n-1}}^* \\
&\quad \times \overleftarrow{\prod}_{i \neq j} \theta_{a_i^{n-1}}^{\perp}(u_i^{(n-1)}) d^{\perp}(u_j^{(n-1)}) \tag{4.35}
\end{aligned}$$

$$\begin{aligned}
&+ \sum_j \frac{f^{-}(u_j^{(n-1)}, \mathbf{u}^{(n-1)} \setminus u_j^{(n-1)})}{u_j^{(n-1)} - \tilde{u}_{m_n}^{(n)}} \sum_{k < n} \left( \frac{1}{u_j^{(n-1)} - u_{m_n}^{(n)}} E_{k+2}^{(\hat{n})} \otimes E_2^{(\hat{n})} + E_2^{(\hat{n})} \otimes E_{k+2}^{(\hat{n})} \right) \\
&\quad \otimes (E_k^{(n-1)})_{a_j^{n-1}}^* \overleftarrow{\prod}_{i \neq j} \theta_{a_i^{n-1}}^{\perp}(u_i^{(n-1)}) c^{\perp}(u_j^{(n-1)}) \tag{4.36}
\end{aligned}$$

$$+ \sum_{j < j'} f^{-}((u_j^{(n-1)}, u_{j'}^{(n-1)}), \mathbf{u}^{(n-1)} \setminus (u_j^{(n-1)}, u_{j'}^{(n-1)}))$$

523

$$\begin{aligned}
& \times \sum_{k,l < n} \left( \frac{1}{\gamma} \left( \alpha_{11} E_{k+2}^{(\hat{n})} \otimes E_{l+2}^{(\hat{n})} + \alpha_{12} E_{l+2}^{(\hat{n})} \otimes E_{k+2}^{(\hat{n})} \right) \otimes (E_k^{(n-1)})_{a_j^{n-1}}^* (E_l^{(n-1)})_{a_{j'}^{n-1}}^* \right. \\
& \quad \times \prod_{i \neq j, j'}^{\leftarrow} \theta_{a_i^{n-1}}^1(u_i^{(n-1)}) c^l(u_{j'}^{(n-1)}) d^l(u_j^{(n-1)}) \\
& \quad + \frac{1}{\gamma} \left( \alpha_{21} E_{k+2}^{(\hat{n})} \otimes E_{l+2}^{(\hat{n})} + \alpha_{22} E_{l+2}^{(\hat{n})} \otimes E_{k+2}^{(\hat{n})} \right) \otimes (E_k^{(n-1)})_{a_j^{n-1}}^* (E_l^{(n-1)})_{a_{j'}^{n-1}}^* \\
& \quad \times \prod_{i \neq j, j'}^{\leftarrow} \theta_{a_i^{n-1}}^1(u_i^{(n-1)}) c^l(u_j^{(n-1)}) d^l(u_{j'}^{(n-1)}) \left. \right) \Big] \quad (4.37) \\
& \times \Psi(\mathbf{u}^{(1\dots n-2)} | \mathbf{u}^{(n-1,n)} \setminus u_{m_n}^{(n)})
\end{aligned}$$

524 where

$$\begin{aligned}
\alpha_{11} &:= (u_{j'}^{(n-1)} - u_{m_n}^{(n)})(u_j^{(n-1)} - \tilde{u}_{m_n}^{(n)}) - (u_{j'}^{(n-1)} - u_{m_n}^{(n)}) / (u_j^{(n-1)} - u_{j'}^{(n-1)}), \\
\alpha_{12} &:= u_j^{(n-1)} - \tilde{u}_{m_n}^{(n)} - ((u_j^{(n-1)} - u_{m_n}^{(n)})(u_{j'}^{(n-1)} - \tilde{u}_{m_n}^{(n)}) + 1) / (u_j^{(n-1)} - u_{j'}^{(n-1)}), \\
\alpha_{21} &:= f^+(u_j^{(n-1)}, u_{j'}^{(n-1)})(u_{j'}^{(n-1)} - u_{m_n}^{(n)}), \\
\alpha_{22} &:= f^+(u_j^{(n-1)}, u_{j'}^{(n-1)})((u_j^{(n-1)} - u_{m_n}^{(n)})(u_{j'}^{(n-1)} - \tilde{u}_{m_n}^{(n)}) + 1), \\
\gamma &:= (u_j^{(n-1)} - u_{m_n}^{(n)})(u_j^{(n-1)} - \tilde{u}_{m_n}^{(n)})(u_{j'}^{(n-1)} - u_{m_n}^{(n)})(u_{j'}^{(n-1)} - \tilde{u}_{m_n}^{(n)}).
\end{aligned} \quad (4.38)$$

525 We will consider the terms (4.34–4.37) individually.

526 First, consider the term (4.34). Acting with  $\theta_{\hat{a}_{m_n} \check{a}_{m_n}}^{(n)}(u_{m_n}^{(n)}) \mathcal{B}^{(n)}(\mathbf{u}^{(n)} \setminus u_{m_n}^{(n)})$  gives the  $i = n$   
527 case of the first term on the right hand side of (4.13).

528 Next, consider the term (4.35). The operator  $d^l(u_j^{(n-1)})$  acts on  $\Psi(\mathbf{u}^{(1\dots n-2)} | \mathbf{u}^{(n-1,n)} \setminus u_{m_n}^{(n)})$   
529 via multiplication by  $f^+(u_j^{(n-1)}, \mathbf{u}^{(n)} \setminus u_{m_n}^{(n)}) \mu_n(u_j^{(n-1)})$  giving

$$\begin{aligned}
& \sum_j \frac{\Gamma_n(u_j^{(n-1)}; \mathbf{u}^{(1\dots n)} \setminus (u_j^{(n-1)}, u_{m_n}^{(n)}))}{u_j^{(n-1)} - u_{m_n}^{(n)}} \sum_{k < n} E_{k+2}^{(\hat{n})} \otimes E_1^{(\hat{n})} \otimes (E_k^{(n-1)})_{a_j^{n-1}}^* \\
& \quad \times \Psi(\mathbf{u}^{(1\dots n-1)} \setminus u_j^{(n-1)} | u_j^{(n-1)}, \mathbf{u}^{(n)} \setminus u_{m_n}^{(n)}). \quad (4.39)
\end{aligned}$$

530 Using (4.3), we expand  $\Psi(\mathbf{u}^{(1\dots n-1)} \setminus u_j^{(n-1)} | u_j^{(n-1)}, \mathbf{u}^{(n)} \setminus u_{m_n}^{(n)})$  in the space  $V_{a_j^{n-1}}^{(n-1)}$ :

$$\sum_{i < n} \sum_{\substack{|\mathbf{u}_{\text{III}}^{(r)}|=1 \\ i \leq r < n-1}} \prod_{i < k < n} \frac{\Gamma_k(\mathbf{u}_{\text{III}}^{(k-1)}; \mathbf{u}_{\text{I}}^{(1\dots n)})}{\mathbf{u}_{\text{III}}^{(k-1)} - \mathbf{u}_{\text{III}}^{(k)}} E_{n-i}^{(n-1)} \otimes \Psi(\mathbf{u}_{\text{I}}^{(1\dots n-1)} | \mathbf{u}_{\text{I}}^{(n)}) \quad (4.40)$$

531 where  $\mathbf{u}_{\text{III}}^{(n-1)} := u_j^{(n-1)}$  and  $\mathbf{u}_{\text{I}}^{(n)} := \mathbf{u}^{(n)} \setminus u_{m_n}^{(n)}$ . Substituting (4.40) into (4.39) yields

$$\sum_{i < n} \sum_{\substack{|\mathbf{u}_{\text{III}}^{(r)}|=1 \\ i \leq r < n}} \prod_{i < k \leq n} \frac{\Gamma_k(\mathbf{u}_{\text{III}}^{(k-1)}; \mathbf{u}_{\text{I}}^{(1\dots n)})}{\mathbf{u}_{\text{III}}^{(k-1)} - \mathbf{u}_{\text{III}}^{(k)}} E_{\hat{n}-i+1}^{(\hat{n})} \otimes E_1^{(\hat{n})} \otimes \Psi(\mathbf{u}_{\text{I}}^{(1\dots n-1)} | \mathbf{u}_{\text{I}}^{(n)}). \quad (4.41)$$

532 Acting with  $\theta_{\hat{a}_{m_n} \check{a}_{m_n}}^{(n)}(u_{m_n}^{(n)}) \mathcal{B}^{(n)}(\mathbf{u}^{(n)} \setminus u_{m_n}^{(n)})$  gives the  $i < n$  cases of the first term on the right  
533 hand side of (4.13).

534 We are now ready to consider the term (4.36). Let  $\eta^l$  denote the restriction of  $\eta^m$  to the  
 535 space  $L^l$ . Set  $\eta_i^l := (E_{12}^{(\hat{n})})_{\tilde{a}_i} \cdot \eta^l$ . Using the explicit form of  $c^l(u_j^{(n-1)})$  we find

$$c^l(u_j^{(n-1)}) \cdot \eta^l = \sum_{l < m_n} \frac{\prod_{k < l} f^+(u_j^{(n-1)}, u_k^{(n)})}{u_j^{(n-1)} - u_l^{(n)}} \mu_n(u_j^{(n-1)}) \eta_l^l \quad (4.42)$$

536 giving

$$\begin{aligned} & \sum_j \frac{f^-(u_j^{(n-1)}, \mathbf{u}^{(n-1)} \setminus u_j^{(n-1)})}{u_j^{(n-1)} - \tilde{u}_{m_n}^{(n)}} \sum_{k < n} \left( \frac{1}{u_j^{(n-1)} - u_{m_n}^{(n)}} E_{k+2}^{(\hat{n})} \otimes E_2^{(\hat{n})} + E_2^{(\hat{n})} \otimes E_{k+2}^{(\hat{n})} \right) \otimes (E_k^{(n-1)})_{a_j^{n-1}}^* \\ & \times \sum_{l < m_n} \frac{\prod_{k < l} f^+(u_j^{(n-1)}, u_k^{(n)})}{u_j^{(n-1)} - u_l^{(n)}} \mu_n(u_j^{(n-1)}) \Psi_l(\mathbf{u}^{(1\dots n-1)} \setminus u_j^{(n-1)} | u_j^{(n-1)}, \mathbf{u}^{(n)} \setminus u_{m_n}^{(n)}). \end{aligned} \quad (4.43)$$

537 Acting with  $\mathcal{B}^{(n)}(\mathbf{u}^{(n)} \setminus u_{m_n}^{(n)})$  and applying Lemma 4.8 to the second line of (4.43) gives

$$\begin{aligned} & \sum_{i \leq l < m_n} \frac{\prod_{k < l} f^+(u_j^{(n-1)}, u_k^{(n)})}{\prod_{l < k < m_n} f^+(u_k^{(n)}, u_i^{(n)})} \cdot \frac{\Gamma_{\hat{n}}(u_i^{(n)}, \mathbf{u}^{(1\dots n)} \setminus (u_i^{(n)}, u_{m_n}^{(n)}))}{(u_l^{(n)} - u_i^{(n)} + 1)(u_j^{(n-1)} - u_l^{(n)})} \\ & \times \mu_n(u_j^{(n-1)}) \Psi(\mathbf{u}^{(1\dots n)} \setminus (u_j^{(n-1)}, u_i^{(n)}, u_{m_n}^{(n)}) | u_j^{(n-1)}). \end{aligned} \quad (4.44)$$

538 Using the identity

$$\frac{f^+(u_j^{(n-1)}, \mathbf{u}^{(n)} \setminus (u_i^{(n)}, u_{m_n}^{(n)}))}{u_j^{(n-1)} - u_i^{(n)}} = \sum_{i \leq l < m_n} \frac{\prod_{k < l} f^+(u_j^{(n-1)}, u_k^{(n)})}{\prod_{l < k < m_n} f^+(u_k^{(n)}, u_i^{(n)})} \cdot \frac{1}{(u_l^{(n)} - u_i^{(n)} + 1)(u_j^{(n-1)} - u_l^{(n)})}$$

539 which follows by a descending induction on  $i$ , expression (4.44) becomes

$$\begin{aligned} & \sum_{i < m_n} \frac{f^+(u_j^{(n-1)}, \mathbf{u}^{(n)} \setminus (u_i^{(n)}, u_{m_n}^{(n)}))}{u_j^{(n-1)} - u_i^{(n)}} \Gamma_{\hat{n}}(u_i^{(n)}, \mathbf{u}^{(1\dots n)} \setminus (u_i^{(n)}, u_{m_n}^{(n)})) \\ & \times \mu_n(u_j^{(n-1)}) \Psi(\mathbf{u}^{(1\dots n)} \setminus (u_j^{(n-1)}, u_i^{(n)}, u_{m_n}^{(n)}) | u_j^{(n-1)}). \end{aligned} \quad (4.45)$$

540 Therefore, action of  $\mathcal{B}^{(n)}(\mathbf{u}^{(n)} \setminus u_{m_n}^{(n)})$  on (4.43) gives

$$\begin{aligned} & \sum_j \sum_{i < m_n} \frac{\Gamma_n(u_j^{(n-1)}; \mathbf{u}^{(1\dots n)} \setminus (u_j^{(n-1)}, u_i^{(n)}, u_{m_n}^{(n)})) \Gamma_{\hat{n}}(u_i^{(n)}; \mathbf{u}^{(1\dots n)} \setminus (u_i^{(n)}, u_{m_n}^{(n)}))}{(u_j^{(n-1)} - u_i^{(n)})(u_j^{(n-1)} - \tilde{u}_{m_n}^{(n)})} \\ & \times \sum_{k < n} \left( \frac{1}{u_j^{(n-1)} - u_{m_n}^{(n)}} E_{k+2}^{(\hat{n})} \otimes E_2^{(\hat{n})} + E_2^{(\hat{n})} \otimes E_{k+2}^{(\hat{n})} \right) \otimes (E_k^{(n-1)})_{a_j^{n-1}}^* \\ & \times \Psi(\mathbf{u}^{(1\dots n)} \setminus (u_j^{(n-1)}, u_i^{(n)}, u_{m_n}^{(n)}) | u_j^{(n-1)}). \end{aligned} \quad (4.46)$$

541 Finally, we expand  $\Psi(\mathbf{u}^{(1\dots n)} \setminus (u_j^{(n-1)}, u_i^{(n)}, u_{m_n}^{(n)}) | u_j^{(n-1)})$  in the space  $V_{a_j^{n-1}}^{(n-1)}$  analogously to (4.40).

542 This gives

$$\begin{aligned} & \sum_{i < n} \sum_{\substack{|u_{\text{II}}^{(r)}|=1 \\ i \leq r \leq n}} \prod_{i < k < n} \frac{\Gamma_k(\mathbf{u}_{\text{II}}^{(k-1)}; \mathbf{u}_{\text{I}}^{(1\dots n)})}{\mathbf{u}_{\text{II}}^{(k-1)} - \mathbf{u}_{\text{II}}^{(k)}} \cdot \frac{\Gamma_n(\mathbf{u}_{\text{II}}^{(n-1)}; \mathbf{u}_{\text{I}}^{(1\dots n)}) \Gamma_{\hat{n}}(\mathbf{u}_{\text{II}}^{(n)}; \mathbf{u}_{\text{I}}^{(1\dots n)})}{(\mathbf{u}_{\text{II}}^{(n-1)} - \mathbf{u}_{\text{II}}^{(n)})(\mathbf{u}_{\text{II}}^{(n-1)} - \tilde{\mathbf{u}}_{\text{III}}^{(n)})} \\ & \times \left( \frac{1}{\mathbf{u}_{\text{II}}^{(n-1)} - \mathbf{u}_{\text{III}}^{(n)}} E_i^{(\hat{n})} \otimes E_2^{(\hat{n})} + E_2^{(\hat{n})} \otimes E_i^{(\hat{n})} \right) \otimes \Psi(\mathbf{u}_{\text{I}}^{(1\dots n)}). \end{aligned} \quad (4.47)$$

543 Combining (4.47) with (4.32) and acting with  $\ell_{\tilde{a}_{m_n} \tilde{a}_{m_n}}^{(n)}(u_{m_n}^{(n)})$  gives the second term on the right  
 544 hand side of (4.13).

545 It remains to consider the term (4.37). Using the same arguments as above, and renaming  
 546  $j \rightarrow p, j' \rightarrow p'$ , we obtain

$$\begin{aligned} & \sum_{i < m_n} \sum_{p < p'} \Gamma_n((u_p^{(n-1)}, u_{p'}^{(n-1)}); \mathbf{u}^{(1\dots n)} \setminus (u_p^{(n-1)}, u_{p'}^{(n-1)}, u_i^{(n)}, u_{m_n}^{(n)})) \Gamma_{\hat{n}}(u_i^{(n)}; \mathbf{u}^{(1\dots n)} \setminus (u_i^{(n)}, u_{m_n}^{(n)})) \\ & \times \sum_{k, l < n} \frac{1}{\gamma} \left( \beta_1 E_{k+2}^{(\hat{n})} \otimes E_{l+2}^{(\hat{n})} + \beta_2 E_{l+2}^{(\hat{n})} \otimes E_{k+2}^{(\hat{n})} \right) \otimes (E_k^{(n-1)})_{a_{p'}^{n-1}}^* (E_l^{(n-1)})_{a_{p'}^{n-1}}^* \\ & \times \Psi(\mathbf{u}^{(1\dots n)} \setminus (u_p^{(n-1)}, u_{p'}^{(n-1)}, u_i^{(n)}, u_{m_n}^{(n)}) | u_p^{(n-1)}, u_{p'}^{(n-1)}) \end{aligned} \quad (4.48)$$

547 where

$$\begin{aligned} \beta_1 & := \frac{f^+(u_p^{(n-1)}, u_i^{(n)})}{u_{p'}^{(n-1)} - u_i^{(n)}} \alpha_{11} + \frac{f^+(u_{p'}^{(n-1)}, u_i^{(n)})}{u_p^{(n-1)} - u_i^{(n)}} \alpha_{21} \\ & = \frac{u_{p'}^{(n-1)} - u_{m_n}^{(n)}}{u_{p'}^{(n-1)} - u_i^{(n)}} \left( u_p^{(n-1)} - \tilde{u}_{m_n}^{(n)} + 1 + \frac{u_{p'}^{(n-1)} - \tilde{u}_{m_n}^{(n)}}{u_p^{(n-1)} - u_i^{(n)}} \right), \\ \beta_2 & := \frac{f^+(u_p^{(n-1)}, u_i^{(n)})}{u_{p'}^{(n-1)} - u_i^{(n)}} \alpha_{12} + \frac{f^+(u_{p'}^{(n-1)}, u_i^{(n)})}{u_p^{(n-1)} - u_i^{(n)}} \alpha_{22} \\ & = f^+(u_p^{(n-1)}, u_i^{(n)}) \frac{u_p^{(n-1)} - \tilde{u}_{m_n}^{(n)}}{u_{p'}^{(n-1)} - u_i^{(n)}} + \frac{(u_p^{(n-1)} - u_{m_n}^{(n)})(u_{p'}^{(n-1)} - \tilde{u}_{m_n}^{(n)}) + 1}{u_p^{(n-1)} - u_i^{(n)}}. \end{aligned} \quad (4.49)$$

548 Note that

$$\beta_1 + \beta_2 = \frac{\gamma}{u_i^{(n)} - \tilde{u}_{m_n}^{(n)}} \left( K(u_p^{(n-1)}, u_{p'}^{(n-1)} | u_i^{(n)}, u_{m_n}^{(n)}) - K(u_p^{(n-1)}, u_{p'}^{(n-1)} | \tilde{u}_{m_n}^{(n)}, u_{m_n}^{(n)}) \right). \quad (4.50)$$

549 We can now use (4.4) to expand vector

$$\Psi(\mathbf{u}^{(1\dots n)} \setminus (u_p^{(n-1)}, u_{p'}^{(n-1)}, u_i^{(n)}, u_{m_n}^{(n)}) | u_p^{(n-1)}, u_{p'}^{(n-1)})$$

550 in the space  $V_{a_{p'}^{n-1}}^{(n-1)} \otimes V_{a_p^{n-1}}^{(n-1)}$ :

$$\sum_{1 \leq i < n} \sum_{|\mathbf{u}_{\text{III}}^{(r)}|=(2,0)} \prod_{i < k < n} \Gamma_k(\mathbf{u}_{\text{II}}^{(k-1)}; \mathbf{u}_I^{(1\dots n)}) K(\mathbf{u}_{\text{II}}^{(k-1)} | \mathbf{u}_{\text{II,III}}^{(k)}) E_{n-i}^{(n-1)} \otimes E_{n-i}^{(n-1)} \otimes \Psi(\mathbf{u}_I^{(1\dots n)}) \quad (4.51)$$

$$\begin{aligned} & + \sum_{1 \leq i < j < n} \sum_{|\mathbf{u}_{\text{III}}^{(r)}|=1} \sum_{|\mathbf{u}_{\text{II}}^{(s)}|=1} \prod_{i < k < j} \frac{\Gamma_k(\mathbf{u}_{\text{III}}^{(k-1)}; \mathbf{u}_I^{(1\dots n)})}{\mathbf{u}_{\text{III}}^{(k-1)} - \mathbf{u}_{\text{III}}^{(k)}} \cdot \Gamma_j(\mathbf{u}_{\text{III}}^{(j-1)}; \mathbf{u}_I^{(1\dots n)}) \\ & \times \prod_{j < k < n} \frac{\Gamma_k(\mathbf{u}_{\text{II}}^{(k-1)}; \mathbf{u}_I^{(1\dots n)}) \Gamma_k(\mathbf{u}_{\text{III}}^{(k-1)}; \mathbf{u}_{\text{I,II}}^{(1\dots n)})}{(\mathbf{u}_{\text{II}}^{(k-1)} - \mathbf{u}_{\text{II}}^{(k)})(\mathbf{u}_{\text{III}}^{(k-1)} - \mathbf{u}_{\text{III}}^{(k)})} \\ & \times \left( \frac{f^+(\mathbf{u}_{\text{III}}^{(j-1)}, \mathbf{u}_{\text{II}}^{(j)})}{\mathbf{u}_{\text{III}}^{(j-1)} - \mathbf{u}_{\text{III}}^{(j)}} E_{n-i}^{(n-1)} \otimes E_{n-j}^{(n-1)} + \frac{1}{\mathbf{u}_{\text{III}}^{(j-1)} - \mathbf{u}_{\text{II}}^{(j)}} E_{n-j}^{(n-1)} \otimes E_{n-i}^{(n-1)} \right) \otimes \Psi(\mathbf{u}_I^{(1\dots n)}) \end{aligned} \quad (4.52)$$

551 where  $\mathbf{u}_{\text{II}}^{(n-1)} := u_p^{(n-1)}$ ,  $\mathbf{u}_{\text{III}}^{(n-1)} := u_{p'}^{(n-1)}$  and  $\mathbf{u}_I^{(n)} := \mathbf{u}^{(n)} \setminus (u_i^{(n)}, u_{m_n}^{(n)})$ .

552 Substituting the term (4.51) into (4.48) and applying (4.50) gives

$$\begin{aligned} & \sum_{\substack{1 \leq i < n \\ i \leq r < n}} \sum_{|u_{\text{II}}^{(r)}|=2} \sum_{|u_{\text{II}}^{(n)}|=1} \prod_{i < k \leq n} \Gamma_k(\mathbf{u}_{\text{II}}^{(k-1)}; \mathbf{u}_{\text{I}}^{(1\dots n)}) \prod_{i < k < n} K(\mathbf{u}_{\text{II}}^{(k-1)} | \mathbf{u}_{\text{II}}^{(k)}) \\ & \times \frac{\Gamma_{\hat{n}}(\mathbf{u}_{\text{II}}^{(n)}; \mathbf{u}_{\text{I}}^{(1\dots n)})}{\mathbf{u}_{\text{II}}^{(n)} - \tilde{\mathbf{u}}_{\text{III}}^{(n)}} \left( K(\mathbf{u}_{\text{II}}^{(n-1)} | \mathbf{u}_{\text{II,III}}^{(n)}) - K(\mathbf{u}_{\text{II}}^{(n-1)} | \tilde{\mathbf{u}}_{\text{III}}^{(n)}, \mathbf{u}_{\text{III}}^{(n)}) \right) E_{\tilde{\tau}}^{(\hat{n})} \otimes E_{\tilde{\tau}}^{(\hat{n})} \otimes \Phi(\mathbf{u}_{\text{I}}^{(1\dots n)}). \end{aligned} \quad (4.53)$$

553 Upon combining (4.53) with (4.31) and acting with  $\theta_{\hat{a}_{m_n} \check{a}_{m_n}}^{(n)}(\mathbf{u}_{m_n}^{(n)})$  gives the third term on the  
554 right hand side of (4.13).

555 Finally, substituting (4.52) into (4.48) and exploiting symmetry of Bethe vectors gives

$$\begin{aligned} & \sum_{\substack{1 \leq i < j < n \\ i \leq r < n}} \sum_{|u_{\text{III}}^{(r)}|=1} \sum_{|u_{\text{III}}^{(s)}|=1} \prod_{i < k < j} \frac{\Gamma_k(\mathbf{u}_{\text{III}}^{(k-1)}; \mathbf{u}_{\text{I}}^{(1\dots n)})}{\mathbf{u}_{\text{III}}^{(k-1)} - \mathbf{u}_{\text{III}}^{(k)}} \cdot \Gamma_j(\mathbf{u}_{\text{III}}^{(j-1)}; \mathbf{u}_{\text{I}}^{(1\dots n)}) \\ & \times \prod_{j < k < n} \frac{\Gamma_k(\mathbf{u}_{\text{II}}^{(k-1)}; \mathbf{u}_{\text{I}}^{(1\dots n)}) \Gamma_k(\mathbf{u}_{\text{III}}^{(k-1)}; \mathbf{u}_{\text{I,II}}^{(1\dots n)})}{(\mathbf{u}_{\text{II}}^{(k-1)} - \mathbf{u}_{\text{II}}^{(k)})(\mathbf{u}_{\text{III}}^{(k-1)} - \mathbf{u}_{\text{III}}^{(k)})} \cdot \Gamma_n(\mathbf{u}_{\text{II,III}}^{(n-1)}; \mathbf{u}_{\text{I}}^{(1\dots n)}) \Gamma_{\hat{n}}(\mathbf{u}_{\text{II}}^{(n)}; \mathbf{u}_{\text{I}}^{(1\dots n)}) \\ & \times \frac{1}{2\gamma} \left[ \left( \left( \beta_2 \frac{f^+(\mathbf{u}_{\text{III}}^{(j-1)}, \mathbf{u}_{\text{II}}^{(j)})}{\mathbf{u}_{\text{III}}^{(j-1)} - \mathbf{u}_{\text{III}}^{(j)}} + \beta_1 \frac{1}{\mathbf{u}_{\text{III}}^{(j-1)} - \mathbf{u}_{\text{II}}^{(j)}} \right) E_{\tilde{\tau}}^{(\hat{n})} \otimes E_{\tilde{\tau}}^{(\hat{n})} \right. \right. \\ & \left. \left. + \left( \beta_1 \frac{f^+(\mathbf{u}_{\text{III}}^{(j-1)}, \mathbf{u}_{\text{II}}^{(j)})}{\mathbf{u}_{\text{III}}^{(j-1)} - \mathbf{u}_{\text{III}}^{(j)}} + \beta_2 \frac{1}{\mathbf{u}_{\text{III}}^{(j-1)} - \mathbf{u}_{\text{II}}^{(j)}} \right) E_{\tilde{\tau}}^{(\hat{n})} \otimes E_{\tilde{\tau}}^{(\hat{n})} \right] \otimes \Psi(\mathbf{u}_{\text{I}}^{(1\dots n)}). \end{aligned} \quad (4.54)$$

556 Combining (4.54) with (4.33) and acting with  $\theta_{\hat{a}_{m_n} \check{a}_{m_n}}^{(n)}(\mathbf{u}_{m_n}^{(n)})$  gives the last term on the right  
557 hand side of (4.13).  $\square$

### 558 4.3 Proof of Lemma 3.8

559 The idea of the proof is to construct a certain Bethe vector and evaluate this vector in two  
560 different ways. Equating the resulting expressions will yield the claim of the Lemma.

561 We begin by rewriting the wanted relation in a more convenient way. From (2.20) and  
562 (3.27) we find that

$$\left\{ \frac{p(v)}{u_i^{(n)} - v} \theta_{\hat{a}_{m_n} \check{a}_{m_n}}^{(n)}(v) \right\}^v = \theta_{\hat{a}_{m_n} \check{a}_{m_n}}^{(n)}(v) \left( \frac{f^+(u_i^{(n)}, \tilde{v})}{u_i^{(n)} - v} + \frac{1}{u_i^{(n)} - \tilde{v}} P_{\hat{a}_{m_n} \check{a}_{m_n}}^{(\hat{n}, \hat{n})} \right). \quad (4.55)$$

563 Repeating the steps used in deriving (4.30) and applying (4.55) we rewrite (3.43) as

$$\begin{aligned} & s_{\hat{n}\hat{n}}(v) \Psi(\mathbf{u}^{(1\dots n)}) = \Gamma_{\hat{n}}(v, \mathbf{u}^{(1\dots n)}) \Psi(\mathbf{u}^{(1\dots n)}) \\ & - \sum_i \theta_{\hat{a}_{m_n} \check{a}_{m_n}}^{(n)}(v) \left( \frac{f^+(u_i^{(n)}, \tilde{v})}{u_i^{(n)} - v} + \frac{1}{u_i^{(n)} - \tilde{v}} P_{\hat{a}_{m_n} \check{a}_{m_n}}^{(\hat{n}, \hat{n})} \right) \\ & \times \Gamma_{\hat{n}}(u_i^{(n)}, \mathbf{u}^{(1\dots n)} \setminus u_i^{(n)}) \mathcal{B}^{(n)}(\mathbf{u}_{\sigma_i}^{(n)} \setminus u_i^{(n)}) \dot{\Psi}_{2,1}(\mathbf{u}^{(1\dots n-1)} | \mathbf{u}_{\sigma_i}^{(n)}) \\ & - \sum_{i \neq i'} \theta_{\hat{a}_{m_n-1} \check{a}_{m_n-1}}^{(n)}(v) \left( \frac{f^+(u_i^{(n)}, \tilde{v})}{u_i^{(n)} - v} + \frac{1}{u_i^{(n)} - \tilde{v}} P_{\hat{a}_{m_n-1} \check{a}_{m_n-1}}^{(\hat{n}, \hat{n})} \right) \\ & \times \Gamma_{\hat{n}}((u_i^{(n)}, u_{i'}^{(n)}); \mathbf{u}^{(1\dots n)} \setminus (u_i^{(n)}, u_{i'}^{(n)})) \frac{f^-(u_i^{(n)}, u_{i'}^{(n)}) f^+(u_i^{(n)}, \tilde{u}_{i'}^{(n)})}{u_{i'}^{(n)} - \tilde{u}_{i'}^{(n)}} \\ & \times \mathcal{B}^{(n)}(\mathbf{u}^{(n)} \setminus (u_i^{(n)}, u_{i'}^{(n)})) \dot{\Psi}_{2,2}(\mathbf{u}^{(1\dots n-1)} | \mathbf{u}_{\sigma_i}^{(n)} \setminus u_{i'}^{(n)}). \end{aligned} \quad (4.56)$$

564 Let  $\Psi_{m_n+1}(\mathbf{u}^{(1\dots n)} \cup \nu)$  denote a Bethe vector with  $m_n+1$  level- $n$  excitations and the reference  
 565 vector  $\eta_{m_n+1}^m := (E_{12}^{(\hat{n})})_{\hat{a}_{m_n+1}} \eta^m$ ; here  $\nu$  denotes the  $(m_n+1)$ -st level- $n$  Bethe root. Applying  
 566 (4.18) and (4.30) to this Bethe vector we obtain

$$\begin{aligned} \Psi_{m_n+1}(\mathbf{u}^{(1\dots n)} \cup \nu) &= \Gamma_{\hat{n}}(\nu, \mathbf{u}^{(1\dots n)}) \Psi(\mathbf{u}^{(1\dots n)}) \\ &\quad - \sum_i \frac{f^+(u_i^{(n)}, \tilde{\nu})}{u_i^{(n)} - \nu} \Gamma_{\hat{n}}(u_i^{(n)}; \mathbf{u}^{(1\dots n)} \setminus u_i^{(n)}) \\ &\quad \times \ell_{\hat{a}_{m_n} \hat{a}_{m_n}}^{(n)}(\nu) \mathcal{B}^{(n)}(\mathbf{u}^{(n)} \setminus u_i^{(n)}) \dot{\Psi}_{2,1}(\mathbf{u}^{(1\dots n-1)} | \mathbf{u}^{(n)} \setminus u_i^{(n)} \cup \nu) \\ &\quad - \sum_{i' \neq i} \Gamma_{\hat{n}}((u_i^{(n)}, u_{i'}^{(n)}); \mathbf{u}^{(1\dots n)} \setminus (u_i^{(n)}, u_{i'}^{(n)})) K(u_i^{(n)}, u_{i'}^{(n)} | \nu, \tilde{\nu}) f^+(u_i^{(n)}, u_{i'}^{(n)}) \\ &\quad \times \ell_{\hat{a}_{m_n-1} \hat{a}_{m_n-1}}^{(n)}(\nu) \mathcal{B}^{(n)}(\mathbf{u}^{(n)} \setminus (u_i^{(n)}, u_{i'}^{(n)})) \\ &\quad \times \dot{\Psi}_{2,2}(\mathbf{u}^{(1\dots n-1)} | \mathbf{u}^{(n)} \setminus (u_i^{(n)}, u_{i'}^{(n)}) \cup \nu). \end{aligned} \quad (4.57)$$

567 Next, recall (4.25) and note that  $P_{\hat{a}_i \hat{a}_{m_n}}^{(\hat{n}, \hat{n})} \eta_{m_n}^m = P_{\hat{a}_{m_n} \hat{a}_{m_n}}^{(\hat{n}, \hat{n})} \eta_i^m$  giving

$$\mathcal{R}^{(\hat{n})}(u_{m_n}^{(n)}; \mathbf{u}^{(n)} \setminus u_{m_n}^{(n)}) \cdot \eta_{m_n}^m = \eta_{m_n}^m + \sum_{i < m_n} \frac{\prod_{i < k < m_n} f^+(u_k^{(n)}, \tilde{u}_{m_n}^{(n)})}{u_i^{(n)} - \tilde{u}_{m_n}^{(n)}} P_{\hat{a}_{m_n} \hat{a}_{m_n}}^{(\hat{n}, \hat{n})} \eta_i^m. \quad (4.58)$$

568 This yields an analogue of (4.30) for  $\Psi_{m_n+1}(\mathbf{u}^{(1\dots n)} \cup \nu)$ :

$$\begin{aligned} \Psi_{m_n+1}(\mathbf{u}^{(1\dots n)} \cup \nu) &= \ell_{\hat{a}_{m_n+1} \hat{a}_{m_n+1}}^{(n)}(\nu) \mathcal{B}^{(n)}(\mathbf{u}^{(n)}) \dot{\Psi}_{1,1}(\mathbf{u}^{(1\dots n-1)} | \mathbf{u}^{(n)} \cup \nu) \\ &\quad + \sum_i \frac{\Gamma_{\hat{n}}(u_i^{(n)}; \mathbf{u}^{(1\dots n)} \setminus u_i^{(n)})}{u_i^{(n)} - \tilde{\nu}} \\ &\quad \times \ell_{\hat{a}_{m_n} \hat{a}_{m_n}}^{(n)}(\nu) \mathcal{B}^{(n)}(\mathbf{u}^{(n)} \setminus u_i^{(n)}) \dot{\Psi}_{1,2}(\mathbf{u}^{(1\dots n-1)} | \mathbf{u}^{(n)} \setminus u_i^{(n)} \cup \nu). \end{aligned} \quad (4.59)$$

569 The next step is to evaluate products of creation operators  $\mathcal{B}^{(n)}$  and the dotted Bethe vectors  
 570  $\dot{\Psi}$ . This is done by applying the same techniques used in the proof of Proposition 4.6. Hence,  
 571 we will skip the technical details and state the final expressions only.

572 Evaluating the named products in (4.57) and (4.59) gives

$$\begin{aligned} &\mathcal{B}^{(n)}(\mathbf{u}^{(n)} \setminus u_i^{(n)}) \dot{\Psi}_{2,1}(\mathbf{u}^{(1\dots n-1)} | \mathbf{u}^{(n)} \setminus u_i^{(n)} \cup \nu) \\ &= E_2^{(\hat{n})} \otimes E_1^{(\hat{n})} \otimes \Psi(\mathbf{u}^{(1\dots n)} \setminus u_i^{(n)}) \\ &\quad + \sum_j \frac{\Gamma_n(u_j^{(n-1)}; \mathbf{u}^{(1\dots n)} \setminus (u_j^{(n-1)}, u_i^{(n)}))}{u_j^{(n-1)} - \nu} \\ &\quad \times \sum_{1 \leq k < n} E_{k+2}^{(\hat{n})} \otimes E_1^{(\hat{n})} \otimes (E_k^{(n-1)})_{\hat{a}_j^{n-1}}^* \Psi(\mathbf{u}^{(1\dots n)} \setminus (u_j^{(n-1)}, u_i^{(n)} | u_j^{(n-1)})) \\ &\quad + \sum_j \sum_{i' \neq i} \frac{\Gamma_n(u_j^{(n-1)}; \mathbf{u}^{(1\dots n)} \setminus (u_j^{(n-1)}, u_i^{(n)}, u_{i'}^{(n)})) \Gamma_{\hat{n}}(u_{i'}^{(n)}; \mathbf{u}^{(1\dots n)} \setminus (u_i^{(n)}, u_{i'}^{(n)}))}{(u_j^{(n-1)} - \tilde{\nu})(u_j^{(n-1)} - u_{i'}^{(n)})} \\ &\quad \times \sum_{1 \leq k < n} \left( \frac{1}{u_j^{(n-1)} - \nu} E_{k+2}^{(\hat{n})} \otimes E_2^{(\hat{n})} + E_2^{(\hat{n})} \otimes E_{k+2}^{(\hat{n})} \right) \otimes (E_k^{(n-1)})_{\hat{a}_j^{n-1}}^* \\ &\quad \times \Psi(\mathbf{u}^{(1\dots n)} \setminus (u_j^{(n-1)}, u_i^{(n)}, u_{i'}^{(n)} | u_j^{(n-1)})) \end{aligned}$$

$$\begin{aligned}
 & + \sum_{j < j'} \sum_{i' \neq i} \Gamma_n((u_j^{(n-1)}, u_{j'}^{(n-1)}); \mathbf{u}^{(1\dots n)} \setminus (u_j^{(n-1)}, u_{j'}^{(n-1)}, u_i^{(n)}, u_{i'}^{(n)})) \Gamma_{\hat{n}}(u_{i'}^{(n)}; \mathbf{u}^{(1\dots n)} \setminus (u_i^{(n)}, u_{i'}^{(n)})) \\
 & \times \sum_{1 \leq k, l < n} \frac{1}{\gamma} \left( \beta_1^{(21)} E_{k+2}^{(\hat{n})} \otimes E_{l+2}^{(\hat{n})} + \beta_2^{(21)} E_{l+2}^{(\hat{n})} \otimes E_{k+2}^{(\hat{n})} \right) \otimes (E_k^{(n-1)})_{a_j^{n-1}}^* (E_l^{(n-1)})_{a_{j'}^{n-1}}^* \\
 & \times \Psi(\mathbf{u}^{(1\dots n)} \setminus (u_j^{(n-1)}, u_{j'}^{(n-1)}, u_i^{(n)}, u_{i'}^{(n)}) | u_j^{(n-1)}, u_{j'}^{(n-1)}) \quad (4.60)
 \end{aligned}$$

573 and

$$\begin{aligned}
 & \mathcal{B}^{(n)}(\mathbf{u}^{(n)} \setminus u_i^{(n)}) \dot{\Psi}_{1,2}(\mathbf{u}^{(1\dots n-1)} | \mathbf{u}^{(n)} \setminus u_i^{(n)} \cup v) \\
 & = E_1^{(\hat{n})} \otimes E_2^{(\hat{n})} \otimes \Psi(\mathbf{u}^{(1\dots n)} \setminus u_i^{(n)}) \\
 & + \sum_j \frac{\Gamma_n(u_j^{(n-1)}; \mathbf{u}^{(1\dots n)} \setminus (u_j^{(n-1)}, u_i^{(n)}))}{u_j^{(n-1)} - \tilde{v}} \\
 & \times \sum_{1 \leq k < n} \left( \frac{1}{u_j^{(n-1)} - v} E_{k+2}^{(\hat{n})} \otimes E_1^{(\hat{n})} + E_1^{(\hat{n})} \otimes E_{k+2}^{(\hat{n})} \right) \otimes (E_k^{(n-1)})_{a_j^{n-1}}^* \\
 & \times \Psi(\mathbf{u}^{(1\dots n)} \setminus (u_j^{(n-1)}, u_i^{(n)}) | u_j^{(n-1)}) \\
 & + \sum_j \sum_{i' \neq i} \frac{\Gamma_n(u_j^{(n-1)}; \mathbf{u}^{(1\dots n)} \setminus (u_j^{(n-1)}, u_i^{(n)}, u_{i'}^{(n)})) \Gamma_{\hat{n}}(u_{i'}^{(n)}; \mathbf{u}^{(1\dots n)} \setminus (u_i^{(n)}, u_{i'}^{(n)}))}{(u_j^{(n-1)} - v)(u_j^{(n-1)} - u_{i'}^{(n)})} \\
 & \times \sum_{1 \leq k < n} E_{k+2}^{(\hat{n})} \otimes E_2^{(\hat{n})} \otimes (E_k^{(n-1)})_{a_j^{n-1}}^* \Psi(\mathbf{u}^{(1\dots n)} \setminus (u_j^{(n-1)}, u_i^{(n)}, u_{i'}^{(n)}) | u_j^{(n-1)}) \\
 & + \sum_{j < j'} \sum_{i' \neq i} \Gamma_n((u_j^{(n-1)}, u_{j'}^{(n-1)}); \mathbf{u}^{(1\dots n)} \setminus (u_j^{(n-1)}, u_{j'}^{(n-1)}, u_i^{(n)}, u_{i'}^{(n)})) \Gamma_{\hat{n}}(u_{i'}^{(n)}; \mathbf{u}^{(1\dots n)} \setminus (u_i^{(n)}, u_{i'}^{(n)})) \\
 & \times \sum_{1 \leq k, l < n} \frac{1}{\gamma} \left( \beta_1^{(12)} E_{k+2}^{(\hat{n})} \otimes E_{l+2}^{(\hat{n})} + \beta_{12}^{(12)} E_{l+2}^{(\hat{n})} \otimes E_{k+2}^{(\hat{n})} \right) \otimes (E_k^{(n-1)})_{a_j^{n-1}}^* (E_l^{(n-1)})_{a_{j'}^{n-1}}^* \\
 & \times \Psi(\mathbf{u}^{(1\dots n)} \setminus (u_j^{(n-1)}, u_{j'}^{(n-1)}, u_i^{(n)}, u_{i'}^{(n)}) | u_j^{(n-1)}, u_{j'}^{(n-1)}) \quad (4.61)
 \end{aligned}$$

574 and

$$\begin{aligned}
 & \mathcal{B}^{(n)}(\mathbf{u}^{(n)}) \dot{\Psi}_{1,1}(\mathbf{u}^{(1\dots n-1)} | \mathbf{u}^{(n)} \cup v) \\
 & = E_1^{(\hat{n})} \otimes E_1^{(\hat{n})} \otimes \Psi(\mathbf{u}^{(1\dots n)}) \\
 & + \sum_j \sum_i \frac{\Gamma_n(u_j^{(n-1)}; \mathbf{u}^{(1\dots n)} \setminus (u_j^{(n-1)}, u_i^{(n)})) \Gamma_{\hat{n}}(u_i^{(n)}; \mathbf{u}^{(1\dots n)} \setminus u_i^{(n)})}{u_j^{(n-1)} - u_i^{(n)}} \\
 & \times \sum_{1 \leq k < n} \left( \frac{f^+(u_j^{(n-1)}, \tilde{v})}{u_j^{(n-1)} - v} E_{k+2}^{(\hat{n})} \otimes E_1^{(\hat{n})} + \frac{1}{u_j^{(n-1)} - \tilde{v}} E_1^{(\hat{n})} \otimes E_{k+2}^{(\hat{n})} \right) \otimes (E_k^{(n-1)})_{a_j^{n-1}}^* \\
 & \times \Psi(\mathbf{u}^{(1\dots n)} \setminus (u_j^{(n-1)}, u_i^{(n)}) | u_j^{(n-1)}) \\
 & + \sum_{j < j'} \sum_{i < i'} \Gamma_n((u_j^{(n-1)}, u_{j'}^{(n-1)}); \mathbf{u}^{(1\dots n)} \setminus (u_j^{(n-1)}, u_{j'}^{(n-1)}, u_i^{(n)}, u_{i'}^{(n)})) \\
 & \times \Gamma_{\hat{n}}((u_i^{(n)}, u_{i'}^{(n)}); \mathbf{u}^{(1\dots n)} \setminus (u_i^{(n)}, u_{i'}^{(n)})) K(u_j^{(n-1)}, u_{j'}^{(n-1)} | u_i^{(n)}, u_{i'}^{(n)}) f^+(u_i^{(n)}, \tilde{u}_{i'}^{(n)}) \\
 & \times \sum_{1 \leq k, l < n} \left( \beta_1^{(11)} E_{k+2}^{(\hat{n})} \otimes E_{l+2}^{(\hat{n})} + \beta_2^{(11)} E_{l+2}^{(\hat{n})} \otimes E_{k+2}^{(\hat{n})} \right) \otimes (E_k^{(n-1)})_{a_j^{n-1}}^* (E_l^{(n-1)})_{a_{j'}^{n-1}}^* \\
 & \times \Psi(\mathbf{u}^{(1\dots n)} \setminus (u_j^{(n-1)}, u_{j'}^{(n-1)}, u_i^{(n)}, u_{i'}^{(n)}) | u_j^{(n-1)}, u_{j'}^{(n-1)}) \quad (4.62)
 \end{aligned}$$



575 and

$$\begin{aligned}
& \mathcal{B}^{(n)}(\mathbf{u}^{(n)} \setminus (u_i^{(n)}, u_{i'}^{(n)})) \dot{\Psi}_{2,2}(\mathbf{u}^{(1\dots n-1)} \mid \mathbf{u}^{(n)} \setminus (u_i^{(n)}, u_{i'}^{(n)}) \cup v) \\
&= E_2^{(\hat{n})} \otimes E_2^{(\hat{n})} \otimes \Psi(\mathbf{u}^{(1\dots n)} \setminus (u_i^{(n)}, u_{i'}^{(n)})) \\
&+ \sum_j \Gamma_n(u_j^{(n-1)}; \mathbf{u}^{(1\dots n)} \setminus (u_j^{(n-1)}, u_i^{(n)}, u_{i'}^{(n)})) \\
&\quad \times \sum_{1 \leq k < n} \left( \frac{f^+(u_j^{(n-1)}, \tilde{v})}{u_j^{(n-1)} - v} E_{k+2}^{(\hat{n})} \otimes E_2^{(\hat{n})} + \frac{1}{u_j^{(n-1)} - \tilde{v}} E_2^{(\hat{n})} \otimes E_{k+2}^{(\hat{n})} \right) \otimes (E_k^{(n-1)})_{a_j^{n-1}}^* \\
&\quad \quad \quad \times \Psi(\mathbf{u}^{(1\dots n)} \setminus (u_j^{(n-1)}, u_i^{(n)}, u_{i'}^{(n)}) \mid u_j^{(n-1)}) \\
&+ \sum_{j < j'} \Gamma_n((u_j^{(n-1)}, u_{j'}^{(n-1)}); \mathbf{u}^{(1\dots n)} \setminus (u_j^{(n-1)}, u_{j'}^{(n-1)}, u_i^{(n)}, u_{i'}^{(n)})) \\
&\quad \times \sum_{1 \leq k, l < n} \left( \beta_1^{(11)} E_{k+2}^{(\hat{n})} \otimes E_{l+2}^{(\hat{n})} + \beta_2^{(11)} E_{l+2}^{(\hat{n})} \otimes E_{k+2}^{(\hat{n})} \right) \otimes (E_k^{(n-1)})_{a_j^{n-1}}^* (E_l^{(n-1)})_{a_{j'}^{n-1}}^* \\
&\quad \quad \quad \times \Psi(\mathbf{u}^{(1\dots n)} \setminus (u_j^{(n-1)}, u_{j'}^{(n-1)}, u_i^{(n)}, u_{i'}^{(n)}) \mid u_j^{(n-1)}, u_{j'}^{(n-1)}, \mathbf{u}^{(n)}) \quad (4.63)
\end{aligned}$$

576 where  $\beta_1^{(21)}$ ,  $\beta_2^{(21)}$  and  $\gamma$  are given by (4.49) and (4.38) except  $u_{m_n}^{(n)}$  should be replaced by  $v$ ,  
577 and

$$\begin{aligned}
\beta_1^{(12)} &:= \frac{u_{j'}^{(n-1)} - v}{u_j^{(n-1)} - u_{i'}^{(n)}} \left( f^+(u_j^{(n-1)}, u_{i'}^{(n)}) + \frac{(u_{j'}^{(n-1)} - u_{i'}^{(n)})(u_j^{(n-1)} - \tilde{v})}{u_j^{(n-1)} - u_{i'}^{(n)}} \right), \\
\beta_2^{(12)} &:= \frac{u_j^{(n-1)} - \tilde{v}}{u_j^{(n-1)} - u_{i'}^{(n)}} f^+(u_{j'}^{(n-1)}, u_{i'}^{(n)}) f^+(u_j^{(n-1)}, u_{j'}^{(n-1)}) \\
&\quad + \frac{u_{j'}^{(n-1)} - \tilde{v}}{u_j^{(n-1)} - u_{i'}^{(n)}} f^+(u_j^{(n-1)}, u_{i'}^{(n)}) \left( u_j^{(n-1)} - v - \frac{1}{u_j^{(n-1)} - u_{j'}^{(n-1)}} \right), \\
\beta_1^{(11)} &:= \frac{f^+(u_j^{(n-1)}, \tilde{v})}{(u_j^{(n-1)} - v)(u_{j'}^{(n-1)} - \tilde{v})}, \quad \beta_2^{(11)} := \frac{1}{u_{j'}^{(n-1)} - v} \left( \beta_1^{(11)} + \frac{1}{u_j^{(n-1)} - \tilde{v}} \right). \quad (4.64)
\end{aligned}$$

578 Adapting (4.60) and (4.63) to the relevant products in (4.56) allows us to rewrite the latter  
579 as

$$\begin{aligned}
& \Gamma_{\hat{n}}(u_i^{(n)}, \mathbf{u}^{(1\dots n)} \setminus u_i^{(n)}) E_2^{(\hat{n})} \otimes E_1^{(\hat{n})} \otimes \Psi(\mathbf{u}^{(1\dots n)} \setminus u_i^{(n)}) \\
&+ \sum_j \frac{\Gamma_{\hat{n}}(u_i^{(n)}, \mathbf{u}^{(1\dots n)} \setminus u_i^{(n)}) \Gamma_n(u_j^{(n-1)}; \mathbf{u}^{(1\dots n)} \setminus (u_j^{(n-1)}, u_i^{(n)}))}{u_j^{(n-1)} - u_i^{(n)}} \\
&\quad \times \sum_{1 \leq k < n} E_{k+2}^{(\hat{n})} \otimes E_1^{(\hat{n})} \otimes (E_k^{(n-1)})_{a_j^{n-1}}^* \Psi(\mathbf{u}^{(1\dots n)} \setminus (u_j^{(n-1)}, u_i^{(n)}) \mid u_j^{(n-1)}) \\
&+ \sum_{i' \neq i} \Gamma_{\hat{n}}((u_i^{(n)}, u_{i'}^{(n)}); \mathbf{u}^{(1\dots n)} \setminus (u_i^{(n)}, u_{i'}^{(n)})) \frac{f^-(u_i^{(n)}, u_{i'}^{(n)}) f^+(u_i^{(n)}, \tilde{u}_{i'}^{(n)})}{u_{i'}^{(n)} - \tilde{u}_i^{(n)}} \\
&\quad \times \left( E_2^{(\hat{n})} \otimes E_2^{(\hat{n})} \otimes \Psi(\mathbf{u}^{(1\dots n)} \setminus (u_i^{(n)}, u_{i'}^{(n)})) + A \right) \quad (4.65)
\end{aligned}$$

580 where

$$\begin{aligned}
A := & \sum_j \Gamma_n(u_j^{(n-1)}; \mathbf{u}^{(1\dots n)} \setminus (u_j^{(n-1)}, u_i^{(n)}, u_{i'}^{(n)})) \\
& \times \sum_{1 \leq k < n} \left( \frac{f^+(u_j^{(n-1)}, \tilde{u}_{i'}^{(n)})}{u_j^{(n-1)} - u_i^{(n)}} E_{k+2}^{(\hat{n})} \otimes E_2^{(\hat{n})} + \frac{1}{u_j^{(n-1)} - u_{i'}^{(n)}} E_2^{(\hat{n})} \otimes E_{k+2}^{(\hat{n})} \right) \otimes (E_k^{(n-1)})_{a_j}^* \\
& \times \Psi^{(n-1)}(\mathbf{u}^{(1\dots n)} \setminus (u_j^{(n-1)}, u_i^{(n)}, u_{i'}^{(n)}) | u_j^{(n-1)}) \\
& + \sum_{j < j'} \frac{\Gamma_n((u_j^{(n-1)}, u_{j'}^{(n-1)}); \mathbf{u}^{(1\dots n)} \setminus (u_j^{(n-1)}, u_{j'}^{(n-1)}, u_i^{(n)}, u_{i'}^{(n)}))}{(u_j^{(n-1)} - u_i^{(n)})(u_{j'}^{(n-1)} - u_{i'}^{(n)})} \\
& \times \sum_{1 \leq k, l < n} \left( f^+(u_j^{(n-1)}, u_i^{(n)}) E_{k+2}^{(\hat{n})} \otimes E_{l+2}^{(\hat{n})} + \theta E_{l+2}^{(\hat{n})} \otimes E_{k+2}^{(\hat{n})} \right) \otimes (E_k^{(n-1)})_{a_j}^* (E_l^{(n-1)})_{a_{j'}}^* \\
& \times \Psi(\mathbf{u}^{(1\dots n)} \setminus (u_j^{(n-1)}, u_{j'}^{(n-1)}, u_i^{(n)}, u_{i'}^{(n)}) | u_j^{(n-1)}, u_{j'}^{(n-1)})
\end{aligned}$$

581 and

$$\theta := \frac{(u_j^{(n-1)} - u_i^{(n)})(u_{j'}^{(n-1)} - u_{i'}^{(n)}) + u_j^{(n-1)} - u_{i'}^{(n)} + 1}{(u_j^{(n-1)} - u_{i'}^{(n)})(u_{j'}^{(n-1)} - u_i^{(n)})}.$$

582 The final step is to substitute (4.60)–(4.63) into the difference of (4.59) and (4.57), and (4.65)  
583 into (4.56), and equate the resulting expressions.

## 584 5 Conclusions

585 This paper is a continuation of [GMR19], where twisted Yangian based models, known as one-  
586 dimensional “soliton non-preserving” open spin chains, were studied by means of algebraic  
587 Bethe ansatz. The present paper extends the results of [GMR19] to the odd case, when the  
588 bulk symmetry is  $\mathfrak{gl}_{2n+1}$  and the boundary symmetry is  $\mathfrak{so}_{2n+1}$ . Theorem 3.9 states that Bethe  
589 vectors, defined by formula (3.42), are eigenvectors of the transfer matrix, defined by formula  
590 (3.44), provided Bethe equations (3.53) and (3.54) hold. It is important to note that Bethe  
591 equations for  $Y^\pm(\mathfrak{gl}_N)$ -based models were first considered in [Doi00, AA+05]. However, the  
592 completeness of solutions of such Bethe equations is still an open question. Investigation of  
593 higher-order transfer matrices and  $Q$ -operators might help to shed more light on this problem.

594 In Proposition 3.12 we presented a more symmetric form of the trace formula for Bethe  
595 vectors than the one found in [GMR19]. This formula can be used to obtain Bethe vectors when  
596 the number of excitations is not large since the complexity of the “master” creation operator  
597 grows rapidly when the total excitation number increases. This is a well-known issue of trace  
598 formulas for both closed and open spin chains. Low rank examples of the “master” creation  
599 operator are given in Example 3.11.

600 We also obtained recurrence relations for twisted Yangian based Bethe vectors. They are  
601 given in Propositions 4.4 and 4.6 for even and odd cases, respectively. Repeated application  
602 of these relations allow us to express  $Y^\pm(\mathfrak{gl}_N)$ -based Bethe vectors in terms of  $Y(\mathfrak{gl}_n)$ -based  
603 Bethe vectors obeying recurrence relations found in [HL+17b] and recalled in Appendix A.3.  
604 The recurrence relations found in this paper provide elegant expressions when the rank is  
605 small, see Examples 4.5 and 4.7. The  $n = 2$  even case in Example 4.5 may help investigating  
606 the open fishchain studied in [GJP21]. However, recurrence relations become rather complex  
607 when the rank is not small, especially in the odd case. This raises a natural question, if there  
608 exists an alternative (simpler) method of constructing Bethe vectors for open spin chains. For  
609 closed spin chains the current (“Drinfeld New”) presentation of Yangians and quantum loop

610 algebras [Dri88] has played a significant role in obtaining not only recurrence relations, but  
 611 also action relations, scalar products and norms of Bethe vectors, see [HL<sup>+</sup>17a,HL<sup>+</sup>17b,HL<sup>+</sup>18a,  
 612 HL<sup>+</sup>18b,HL<sup>+</sup>20]. Thus, it is natural to expect that a current presentation of twisted Yangians  
 613 could pave a fruitful path for open spin chains analysis.

614 A current presentation of twisted Yangian  $Y^+(\mathfrak{gl}_N)$  was recently obtained in [LWZ23].  
 615 (The rank 2 case was considered earlier in [Brw16].) However, in [LWZ23] a different, the  
 616 so-called non-split, presentation of twisted Yangian is considered (see Chapter 2 in [Mol07]),  
 617 which is based on the Chevalley involution of  $\mathfrak{gl}_N$  and is not compatible (at least in a natural  
 618 way) with the Bethe vacuum state. Nonetheless, we believe that the presentation obtained  
 619 in [LWZ23] may have applications in open spin chain analysis and deserves attention. For  
 620 example, integrable overlaps for twisted boundary states are constructed using the non-split  
 621 presentation of twisted Yangians [Gom24].

622 Overall, the approach presented in this paper does open a door to an exploration of scalar  
 623 products and norms of Bethe vectors for twisted Yangian based models. However, developing  
 624 Bethe ansatz techniques in the current presentation of twisted Yangians might open a broader  
 625 path to open spin chain analysis. An alternative path could be a development of separation of  
 626 variable techniques along the lines of e.g. [GLMS17,RV21].

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 630 recurrence relations and Paul Ryan for helpful discussions on applications of twisted Yangian  
 631 based models.

## 632 A Appendix

### 633 A.1 Weight grading of $Y^\pm(\mathfrak{gl}_N)$

634 Define an  $n$ -tuple  $\omega_i \in \mathbb{Z}^n$  by  $(\omega_i)_j := \delta_{ij}$  and recall the notation  $\bar{j} = N - j + 1$ . Then define  
 635 weights of the elements  $s_{ij}[r]$  using the following rule

$$\text{wt}(s_{ij}[r]) := \sum_{i \leq k < j} \omega_k + \sum_{\bar{j} \leq k < \hat{n}} \omega_k \quad \text{when } i < j, \quad i + j \leq N + 1 \quad (\text{A.1})$$

636 and require

$$\text{wt}(s_{\bar{j}\bar{i}}[r]) = \text{wt}(s_{ij}[r]), \quad \text{wt}(s_{ji}[r]) = -\text{wt}(s_{ij}[r]) \quad (\text{A.2})$$

637 for all  $1 \leq i, j \leq N$ . Note that  $\text{wt}(s_{ii}[r]) = (0, \dots, 0) \in \mathbb{Z}^n$ . Extending linearly on all monomials  
 638 this defines a weight grading on  $Y^\pm(\mathfrak{gl}_N)$ .

639 The recurrence relations (4.9) and (4.13) are compatible with this grading. The master  
 640 creation operator (3.55) has the weight

$$\omega := \text{wt}(\mathcal{B}_N(\mathbf{u}^{(1\dots n)})) = \begin{cases} (m_1, \dots, m_{n-1}, m_n) & \text{when } \hat{n} = n, \\ (m_1, \dots, m_{n-1}, 2m_n) & \text{when } \hat{n} = n + 1 \end{cases} \quad (\text{A.3})$$

641 which we assign to the corresponding Bethe vector. Then (4.9) and (4.13) can be schematically  
 642 written as

$$\Psi^\omega = \sum_{\omega' \in W} s_{\omega'} \Psi^{\omega - \omega'} \quad (\text{A.4})$$

643 where  $W$  is the set of weights of  $s_{i,n+j}[r]$  with  $1 \leq i \leq n$  and  $1 \leq j \leq \hat{n}$ , the  $s_{\omega'}$  is a generating  
 644 series of  $Y^\pm(\mathfrak{gl}_N)$  of weight  $\omega'$ , and all scalar factors and spectral parameter dependencies are  
 645 omitted, as in (1.1) and (1.2).

## 646 A.2 Commutativity of transfer matrices

647 **Lemma A.1.** *Transfer matrices  $\tau(u)$  defined by (3.44) form a commuting family of operators.*

648 *Proof.* We follow arguments in the Proof of Theorem 2.4 in [V115]. In this proof, we will write  
649  $S_a(u)$  instead of  $S_a^{(N)}(u)$  and  $R_{ab}(u)$  instead of  $R_{ab}^{(N,N)}(u)$ . Then

$$\tau(u)\tau(v) = \text{tr}_a M_a^{t_a}(u) S_a^{t_a}(u) \text{tr}_b M_b(v) S_b(v) = \text{tr}_{ab} M_a^{t_a}(u) M_b(v) S_a^{t_a}(u) S_b(v) \quad (\text{A.5})$$

650 where  $t_a$  denotes the usual matrix transposition in the space labelled  $a$ . Upon inserting a  
651 resolution of identity in terms of  $\widehat{R}$ -matrices and using properties of matrix transposition and  
652 the trace (see Appendix A in [V115]) we rewrite the right hand side of (A.5) as

$$\begin{aligned} & \text{tr}_{ab} M_a^{t_a}(u) M_b(v) (\widehat{R}_{ab}^{t_a}(\tilde{v}-u))^{-1} \widehat{R}_{ab}^{t_a}(\tilde{v}-u) S_a^{t_a}(u) S_b(v) \\ &= \text{tr}_{ab} \left( M_a^{t_a}(u) ((\widehat{R}_{ab}^{t_a}(\tilde{v}-u))^{-1})^{t_b} M_b^{t_b}(v) \right)^{t_b} \left( S_a(u) \widehat{R}_{ab}(\tilde{v}-u) S_b(v) \right)^{t_a} \\ &= \text{tr}_{ab} \left( M_a^{t_a}(u) ((\widehat{R}_{ab}^{t_a}(\tilde{v}-u))^{-1})^{t_b} M_b^{t_b}(v) \right)^{t_b t_a} S_a(u) \widehat{R}_{ab}(\tilde{v}-u) S_b(v). \end{aligned} \quad (\text{A.6})$$

653 We insert a resolution identity in terms of  $R$ -matrices and use properties of matrix transposition  
654 and the trace once again. This gives

$$\begin{aligned} & \text{tr}_{ab} \left( M_a^{t_a}(u) ((\widehat{R}_{ab}^{t_a}(\tilde{v}-u))^{-1})^{t_b} M_b^{t_b}(v) \right)^{t_b t_a} \\ & \quad \times (R_{ab}(u-v))^{-1} R_{ab}(u-v) S_a(u) \widehat{R}_{ab}(\tilde{v}-u) S_b(v) \\ &= \text{tr}_{ab} \left( ((R_{ab}(u-v))^{-1})^{t_a t_b} M_a^{t_a}(u) ((\widehat{R}_{ab}^{t_a}(\tilde{v}-u))^{-1})^{t_b} M_b^{t_b}(v) \right)^{t_b t_a} \\ & \quad \times R_{ab}(u-v) S_a(u) \widehat{R}_{ab}(\tilde{v}-u) S_b(v). \end{aligned} \quad (\text{A.7})$$

655 The  $R$ -matrix (2.6) satisfies

$$((R_{ab}(u))^{-1})^{t_a t_b} = r(u) R_{ab}(-u), \quad ((\widehat{R}_{ab}^{t_a}(u))^{-1})^{t_b} = r(u) \widehat{R}_{ab}(-u) \quad (\text{A.8})$$

656 where  $r(u) := u^2/(u^2 - 1)$ . Relations (A.8) and the dual twisted reflection equation (3.45)  
657 imply

$$\begin{aligned} & \left( ((R_{ab}(u-v))^{-1})^{t_a t_b} M_a^{t_a}(u) ((\widehat{R}_{ab}^{t_a}(\tilde{v}-u))^{-1})^{t_b} M_b^{t_b}(v) \right)^{t_b t_a} \\ &= r(u-v) r(\tilde{v}-u) \left( R_{ab}(v-u) M_a^{t_a}(u) \widehat{R}_{ab}(u-\tilde{v}) M_b^{t_b}(v) \right)^{t_b t_a} \\ &= r(u-v) r(\tilde{v}-u) \left( M_b^{t_b}(v) \widehat{R}_{ab}(u-\tilde{v}) M_a^{t_a}(u) R_{ab}(v-u) \right)^{t_b t_a} \\ &= \left( M_b^{t_b}(v) ((\widehat{R}_{ab}^{t_a}(\tilde{v}-u))^{-1})^{t_b} M_a^{t_a}(u) ((R_{ab}(u-v))^{-1})^{t_a t_b} \right)^{t_b t_a}. \end{aligned} \quad (\text{A.9})$$

658 Applying (A.9) to the right hand side of (A.7) gives

$$\begin{aligned} & \text{tr}_{ab} \left( M_b^{t_b}(v) ((R_{ab}^{t_a}(\tilde{v}-u))^{-1})^{t_b} M_a^{t_a}(u) ((R_{ab}(u-v))^{-1})^{t_a t_b} \right)^{t_b t_a} \\ & \quad \times S_b(v) \widehat{R}_{ab}(\tilde{v}-u) S_a(u) R_{ab}(u-v). \end{aligned} \quad (\text{A.10})$$

659 It remains to repeat similar steps as above in reversed order and use cyclicity of the trace.

660 The (A.10) then becomes

$$\begin{aligned} & \text{tr}_{ab} (R_{ab}(u-v))^{-1} \left( M_b^{t_b}(v) ((R_{ab}^{t_a}(\tilde{v}-u))^{-1})^{t_b} M_a^{t_a}(u) \right)^{t_b t_a} \\ & \quad \times S_b(v) \widehat{R}_{ab}(\tilde{v}-u) S_a(u) R_{ab}(u-v) \\ &= \text{tr}_{ab} \left( M_b^{t_b}(v) ((R_{ab}^{t_a}(\tilde{v}-u))^{-1})^{t_b} M_a^{t_a}(u) \right)^{t_b t_a} S_b(v) \widehat{R}_{ab}(\tilde{v}-u) S_a(u) \\ &= \text{tr}_{ab} \left( M_b^{t_b}(v) ((R_{ab}^{t_a}(\tilde{v}-u))^{-1})^{t_b} M_a^{t_a}(u) \right)^{t_b} \left( S_b(v) \widehat{R}_{ab}(\tilde{v}-u) S_a(u) \right)^{t_a} \\ &= \text{tr}_{ab} (R_{ab}^{t_a}(\tilde{v}-u))^{-1} M_b^{t_b}(v) M_a^{t_a}(u) S_b(v) S_a(u)^{t_a} \widehat{R}_{ab}^{t_a}(\tilde{v}-u) \\ &= \text{tr}_b M_b^{t_b}(v) S_b(v) \text{tr}_a M_a^{t_a}(u) S_a(u)^{t_a} \widehat{R}_{ab}^{t_a}(\tilde{v}-u) = \tau(v) \tau(u) \end{aligned} \quad (\text{A.11})$$

661 as required.  $\square$

### 662 A.3 A recurrence relation for $Y(\mathfrak{gl}_n)$ -based models

663 The Proposition below is a restatement of Proposition 4.2 in [HL<sup>+</sup>17b] in terms of notation  
664 introduced in Section 4.1 and Proposition 4.1. Recall (4.2):

$$\Lambda_k(z; \mathbf{v}^{(1\dots n-1)}) = f^-(z, \mathbf{v}^{(k-1)}) f^+(z, \mathbf{v}^{(k)}) \lambda_k(z).$$

665 Let  $t_{ij}(z)$  denote the standard generating series of  $Y(\mathfrak{gl}_n)$ .

666 **Proposition A.2.**  $Y(\mathfrak{gl}_n)$ -based Bethe vectors satisfy the recurrence relation

$$\Phi(\mathbf{v}^{(1\dots n-1)}) = \sum_{1 \leq i < n} \sum_{\substack{|\mathbf{v}_{\text{II}}^{(r)}|=1 \\ i \leq r < n-1}} \prod_{i < k < n} \frac{\Lambda_k(\mathbf{v}_{\text{II}}^{(k-1)}; \mathbf{v}_{\text{I}}^{(1\dots n-1)})}{\mathbf{v}_{\text{II}}^{(k-1)} - \mathbf{v}_{\text{II}}^{(k)}} t_{in}(\mathbf{v}_{\text{II}}^{(n-1)}) \Phi(\mathbf{v}_{\text{I}}^{(1\dots n-1)}) \quad (\text{A.12})$$

667 where  $\mathbf{v}_{\text{II}}^{(n-1)} = \mathbf{v}_j^{(n-1)}$  for any  $1 \leq j \leq m_{n-1}$  and  $\mathbf{v}_{\text{II}}^{(s)} = \emptyset$  for all  $1 \leq s < i$  so that

$$\mathbf{v}_{\text{I}}^{(1\dots n-1)} = (\mathbf{v}^{(1)}, \dots, \mathbf{v}^{(i-1)}, \mathbf{v}_{\text{I}}^{(i)}, \dots, \mathbf{v}_{\text{I}}^{(n-1)}).$$

668 *Example A.3.* When  $n = 4$ , the recurrence relation (A.12) gives

$$\begin{aligned} \Phi(\mathbf{v}^{(1,2,3)}) &= t_{34}(\mathbf{v}_{\text{II}}^{(3)}) \Phi((\mathbf{v}^{(1)}, \mathbf{v}^{(2)}, \mathbf{v}_{\text{I}}^{(3)})) \\ &+ \sum_{|\mathbf{v}_{\text{II}}^{(2)}|=1} t_{24}(\mathbf{v}_{\text{II}}^{(3)}) \Phi((\mathbf{v}^{(1)}, \mathbf{v}_{\text{I}}^{(2)}, \mathbf{v}_{\text{I}}^{(3)})) \frac{f^-(\mathbf{v}_{\text{II}}^{(2)}, \mathbf{v}_{\text{I}}^{(2)}) f^+(\mathbf{v}_{\text{II}}^{(2)}, \mathbf{v}_{\text{I}}^{(3)}) \lambda_3(\mathbf{v}_{\text{II}}^{(2)})}{\mathbf{v}_{\text{II}}^{(2)} - \mathbf{v}_{\text{II}}^{(3)}} \\ &+ \sum_{\substack{|\mathbf{v}_{\text{II}}^{(r)}|=1 \\ r=1,2}} t_{14}(\mathbf{v}_{\text{II}}^{(3)}) \Phi((\mathbf{v}_{\text{I}}^{(1)}, \mathbf{v}_{\text{I}}^{(2)}, \mathbf{v}_{\text{I}}^{(3)})) \frac{f^-(\mathbf{v}_{\text{II}}^{(1)}, \mathbf{v}_{\text{I}}^{(1)}) f^+(\mathbf{v}_{\text{II}}^{(1)}, \mathbf{v}_{\text{I}}^{(2)}) \lambda_2(\mathbf{v}_{\text{II}}^{(1)})}{\mathbf{v}_{\text{II}}^{(1)} - \mathbf{v}_{\text{II}}^{(2)}} \\ &\quad \times \frac{f^-(\mathbf{v}_{\text{II}}^{(2)}, \mathbf{v}_{\text{I}}^{(2)}) f^+(\mathbf{v}_{\text{II}}^{(2)}, \mathbf{v}_{\text{I}}^{(3)}) \lambda_3(\mathbf{v}_{\text{II}}^{(2)})}{\mathbf{v}_{\text{II}}^{(2)} - \mathbf{v}_{\text{II}}^{(3)}} \end{aligned} \quad (\text{A.13})$$

669 where  $\mathbf{v}_{\text{II}}^{(3)} = \mathbf{v}_j^{(3)}$  for any  $1 \leq j \leq m_3$ , and

## 670 References

- 671 [ADK15] J. Avan, A. Doikou, N. Karaiskos, *The  $\mathfrak{sl}(N)$  twisted Yangian: bulk-boundary*  
672 *scattering and defects*, J. Stat. Mech. P05024 (2015), doi:[10.1088/1742-](https://doi.org/10.1088/1742-5468/2015/05/P05024)  
673 [5468/2015/05/P05024](https://doi.org/10.1088/1742-5468/2015/05/P05024), arXiv:[1412.6480](https://arxiv.org/abs/1412.6480).
- 674 [AA<sup>+</sup>05] D. Arnaudon, J. Avan, N. Crampé, A. Doikou, L. Frappat, E. Ragoucy, *General bound-*  
675 *ary conditions for the  $\mathfrak{sl}(N)$  and  $\mathfrak{sl}(M|N)$  open spin chains*, J. Stat. Mech. P08005  
676 (2004), doi:[10.1088/1742-5468/2004/08/P08005](https://doi.org/10.1088/1742-5468/2004/08/P08005), arXiv:[math-ph/0406021](https://arxiv.org/abs/math-ph/0406021).
- 677 [AC<sup>+</sup>06a] D. Arnaudon, N. Crampe, A. Doikou, L. Frappat, E. Ragoucy, *Analytical Bethe Ansatz*  
678 *for open spin chains with soliton non preserving boundary conditions*, Int. J. Mod. Phys.  
679 **A 21**, 1537 (2006), doi:[10.1142/S0217751X06029077](https://doi.org/10.1142/S0217751X06029077), arXiv:[math-ph/0503014](https://arxiv.org/abs/math-ph/0503014).
- 680 [AC<sup>+</sup>06b] D. Arnaudon, N. Crampe, A. Doikou, L. Frappat, E. Ragoucy, *Spectrum and Bethe*  
681 *ansatz equations for the  $U_q(\mathfrak{gl}(N))$  closed and open spin chains in any representation*,  
682 *Ann. H. Poincaré* vol. 7, 1217 (2006), doi:[10.1007/s00023-006-0280-x](https://doi.org/10.1007/s00023-006-0280-x), arXiv:[math-](https://arxiv.org/abs/math-ph/0512037)  
683 [ph/0512037](https://arxiv.org/abs/math-ph/0512037).

- 684 [Brw16] J. S. Brown, *A Drinfeld presentation for the twisted Yangian  $Y_3^+$* , arXiv:1601.05701.
- 685 [BR08] S. Belliard and E. Ragoucy, *Nested Bethe ansatz for ‘all’ closed spin chains*, J. Phys. A  
686 **41**, 295202 (2008), doi:10.1088/1751-8113/41/29/295202, arXiv:0804.2822.
- 687 [dL<sup>+</sup>19] M. de Leeuw, T. Gombor, C. Kristjansen, G. Linardopoulos, B. Pozsgay,  
688 *Spin chain overlaps and the twisted Yangian*, JHEP **2020**, 176 (2020),  
689 doi:10.1007/JHEP01(2020)176, arXiv:1912.09338.
- 690 [Doi00] A. Doikou, *Quantum spin chain with “soliton non-preserving” boundary conditions*,  
691 J. Phys. A **33**, 8797–8808 (2000), doi:10.1088/0305-4470/33/48/315, arXiv:hep-  
692 th/0006197.
- 693 [Dri88] V. Drinfeld, *A new realization of Yangians and quantized affine algebras*, Soviet Math.  
694 Dokl. **36** (1988), 212–216.
- 695 [DVK87] H.J. De Vega and M. Karowski, *Exact Bethe ansatz solution of  $O(2N)$  symmetric theo-*  
696 *ries*, Nuc. Phys. B **280**, 225–254 (1987), doi:10.1016/0550-3213(87)90146-5.
- 697 [GJP21] N. Gromov, J. Julius, N. Primi, *Open Fishchain in  $N=4$  Supersymmetric Yang-Mills The-*  
698 *ory*, JHEP **2021**, 127 (2021), doi:10.1007/JHEP07(2021)127, arXiv:2101.01232.
- 699 [Ger24] A. J. Gerrard, *On the nested algebraic Bethe ansatz for spin chains with simple  $\mathfrak{g}$ -*  
700 *symmetry*, arXiv:2405.20177.
- 701 [GLMS17] N. Gromov, F. Levkovich-Maslyuk, G. Sizov, *New Construction of Eigenstates and*  
702 *Separation of Variables for  $SU(N)$  Quantum Spin Chains*, JHEP **2017**, 111 (2017),  
703 doi:10.1007/JHEP09(2017)111, arXiv:1610.08032.
- 704 [GMR19] A. Gerrard, N. MacKay, V. Regelskis, *Nested algebraic Bethe ansatz for open spin*  
705 *chains with even twisted Yangian symmetry*, Ann. Henri Poincare **20**, 339–392 (2019),  
706 doi:10.1007/s00023-018-0731-1, arXiv:1710.08409.
- 707 [Gom24] T. Gombor, *Exact overlaps for all integrable two-site boundary states of  $\mathfrak{gl}(N)$*   
708 *symmetric spin chains*, JHEP **2024**, 194 (2024), doi:10.1007/JHEP05(2024)194,  
709 arXiv:2311.04870.
- 710 [GP16] T. Gombor, L. Palla, *Algebraic Bethe Ansatz for  $O(2N)$  sigma models with integrable*  
711 *diagonal boundaries*, JHEP **2016**, 158 (2016), doi:10.1007/JHEP02(2016)158,  
712 arXiv:1511.03107.
- 713 [GR20a] A. Gerrard, V. Regelskis, *Nested algebraic Bethe ansatz for orthogonal*  
714 *and symplectic open spin chains*, Nuc. Phys. B **952**, 114909 (2020),  
715 doi:10.1016/j.nuclphysb.2019.114909, arXiv:1909.12123.
- 716 [GR20b] A. Gerrard, V. Regelskis, *Nested algebraic Bethe ansatz for deformed or-*  
717 *thogonal and symplectic spin chains*, Nuc. Phys. B **956**, 115021 (2020),  
718 doi:10.1016/j.nuclphysb.2020.115021, arXiv:1912.11497.
- 719 [HL<sup>+</sup>17a] A. Hutsalyuk, A. Liashyk, S. Z. Pakuliak, E. Ragoucy, N. A. Slavnov, *Current presenta-*  
720 *tion for the double super-Yangian  $DY(\mathfrak{gl}(m|n))$  and Bethe vectors*, Russ. Math. Survey  
721 **72**, 33–99 (2017), doi:10.1070/RM9754, arXiv:1611.09620.
- 722 [HL<sup>+</sup>17b] A. Hutsalyuk, A. Liashyk, S. Z. Pakuliak, E. Ragoucy, N. A. Slavnov, *Scalar products*  
723 *of Bethe vectors in the models with  $\mathfrak{gl}(m|n)$  symmetry*, Nucl. Phys. B **923**, 277–311  
724 (2017), doi:10.1016/j.nuclphysb.2017.07.020, arXiv:1704.08173.

- 725 [HL<sup>+</sup>18a] A. Hutsalyuk, A. Liashyk, S. Z. Pakuliak, E. Ragoucy, N. A. Slavnov, *Norm of Bethe*  
726 *vectors in models with  $\mathfrak{gl}(m|n)$  symmetry*, Nucl. Phys. B **926**, 256–278 (2018),  
727 doi:[10.1016/j.nuclphysb.2017.07.020](https://doi.org/10.1016/j.nuclphysb.2017.07.020), arXiv:[1705.09219](https://arxiv.org/abs/1705.09219).
- 728 [HL<sup>+</sup>18b] A. Hutsalyuk, A. Liashyk, S. Z. Pakuliak, E. Ragoucy, N. A. Slavnov, *Scalar products*  
729 *and norm of Bethe vectors for integrable models based on  $U_q(\hat{\mathfrak{gl}}_n)$* , SciPost Phys. **4**  
730 (2018) 006, doi:[10.21468/SciPostPhys.4.1.006](https://doi.org/10.21468/SciPostPhys.4.1.006), arXiv:[1711.03867](https://arxiv.org/abs/1711.03867).
- 731 [HL<sup>+</sup>20] A. Hutsalyuk, A. Liashyk, S. Z. Pakuliak, E. Ragoucy, N. A. Slavnov, *Actions of*  
732 *the monodromy matrix elements onto  $\mathfrak{gl}(m|n)$ -invariant Bethe vectors*, J. Stat. Mech.  
733 **2009**, 093104 (2020), doi:[10.1088/1742-5468/abacb2](https://doi.org/10.1088/1742-5468/abacb2), arXiv:[2005.09249](https://arxiv.org/abs/2005.09249).
- 734 [IK84] A. G. Izergin, V. E. Korepin, *The quantum inverse scattering method approach to corre-*  
735 *lation functions*, Comm. Math. Phys. **94**, 67–92 (1984), doi:[10.1007/BF01212350](https://doi.org/10.1007/BF01212350).
- 736 [LWZ23] K. Lu, W. Wang, W. Zhang, *A Drinfeld type presentation of twisted Yangians*,  
737 arXiv:[2308.12254](https://arxiv.org/abs/2308.12254).
- 738 [Mol07] A. Molev, *Yangians and Classical Lie Algebra*, Mathematical Surveys and Monographs  
739 **143**, American Mathematical Society, Providence, RI, 2007, xviii+400 pp.
- 740 [Ols92] G. Olshanskii, *Twisted Yangians and infinite-dimensional classical Lie algebras*, Quan-  
741 *tum groups* (Leningrad, 1990), 104–119, Lecture Notes in Math. **1510**, Springer,  
742 Berlin 1992, doi:[10.1007/BFb0101183](https://doi.org/10.1007/BFb0101183).
- 743 [Reg22] *Algebraic Bethe Ansatz for Spinor R-matrices*, SciPost Phys. **12**, 067 (2022),  
744 doi:[SciPostPhys.12.2.067](https://doi.org/10.21468/SciPostPhys.12.2.067), arXiv:[2108.07580](https://arxiv.org/abs/2108.07580).
- 745 [Rsh85] N. Yu. Reshetikhin, *Integrable Models of Quantum One-dimensional Magnets*  
746 *With  $O(N)$  and  $SP(2k)$  Symmetry*, Theor. Math. Phys. **63**, 555–569 (1985),  
747 doi:[10.1007/BF01017501](https://doi.org/10.1007/BF01017501) (Teor. Mat. Fiz. **63**, no. 3, 347–366 (1985), [link](#)).
- 748 [Rsh91] N. Yu. Reshetikhin, *Algebraic Bethe Ansatz for  $SO(N)$ -invariant Transfer Matrices*, J.  
749 Sov. Math. **54**, 940–951 (1991), doi:[10.1007/BF01101125](https://doi.org/10.1007/BF01101125).
- 750 [RV21] P. Ryan, D. Volin, *Separation of variables for rational  $\mathfrak{gl}(n)$  spin chains in any compact*  
751 *representation, via fusion, embedding morphism and Backlund flow*, Comm. Math.  
752 Phys. **383**, 311–343 (2021), doi:[10.1007/s00220-021-03990-7](https://doi.org/10.1007/s00220-021-03990-7), arXiv:[2002.12341](https://arxiv.org/abs/2002.12341).
- 753 [Sla20] N. A. Slavnov, *Introduction to the nested algebraic Bethe ansatz*, SciPost Phys. Lect.  
754 Notes **19** (2020), doi:[10.21468/SciPostPhysLectNotes.19](https://doi.org/10.21468/SciPostPhysLectNotes.19), arXiv:[1911.12811](https://arxiv.org/abs/1911.12811).
- 755 [Vl15] B. Vlaar, *Boundary transfer matrices and boundary quantum KZ equations*, J. Math.  
756 Phys. **56**, 071705 (2015), doi:[10.1063/1.4927305](https://doi.org/10.1063/1.4927305), arXiv:[1408.3364](https://arxiv.org/abs/1408.3364).